

A look at turbulent impurity and energy transport in stellarators using linear gyrokinetic simulations



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- Neoclassically optimized stellarators: relatively large turbulent energy transport
- \Rightarrow Want to minimize turbulent energy transport.
 - Quasi-isodynamic stellarators: small neoclassical transport at high collisionality
- \implies Turbulent impurity transport likely needed to avoid impurity accumulation.

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 - Compare the ITG turbulence and resulting energy & impurity fluxes between different quasi-symmetric stellarators



Recently published large database of quasi-symmetric stellarators

Theory and definitions

Electrostatic gyrokinetic equation solved with Stella code M. Barnes 2019, JoCP 391 p 365

$$\frac{\partial g_{s}}{\partial t} + v_{\parallel} \vec{b} \cdot \nabla z \left(\frac{\partial g_{s}}{\partial z} + \frac{Z_{s} e}{T_{s}} \frac{\partial \langle \phi \rangle_{\vec{R}}}{\partial z} F_{s} \right) - \frac{\mu_{s}}{m_{s}} \vec{b} \cdot \nabla B \frac{\partial g_{s}}{\partial v_{\parallel}} + \vec{v}_{Ms} \cdot \left(\nabla_{\perp} g_{s} + \frac{Z_{s} e}{T_{s}} \nabla_{\perp} \langle \phi \rangle_{\vec{R}} F_{s} \right) + \langle \vec{v}_{E} \rangle_{\vec{R}} \cdot \nabla_{\perp} g_{s} + \langle \vec{v}_{E} \rangle_{\vec{R}} \cdot \nabla_{E} F_{s} = 0,$$
(1)

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$$\begin{aligned} \frac{\partial g_s}{\partial t} + v_{\parallel} \vec{b} \cdot \nabla z \left(\frac{\partial g_s}{\partial z} + \frac{Z_s e}{T_s} \frac{\partial \langle \phi \rangle_{\vec{R}}}{\partial z} F_s \right) - \frac{\mu_s}{m_s} \vec{b} \cdot \nabla B \frac{\partial g_s}{\partial v_{\parallel}} \\ + \vec{v}_{Ms} \cdot \left(\nabla_{\perp} g_s + \frac{Z_s e}{T_s} \nabla_{\perp} \langle \phi \rangle_{\vec{R}} F_s \right) + \langle \vec{v}_E \rangle_{\vec{R}} \cdot \nabla_{\perp} g_s + \langle \vec{v}_E \rangle_{\vec{R}} \cdot \nabla_E F_s = 0, \end{aligned}$$
Expanded $g = g_{k_x, k_y} \exp(ik_x x + ik_y y)$ and linearize. $(k_x = 0 \text{ here})$
(Raw) fluxes

$$\Gamma_{s,k_y} = -\frac{1}{n_s} k_y \mathcal{I} \int d^3 v g_s \mathcal{J}_0 \phi^*, \qquad (2)$$
$$Q_{s,k_y} = -k_y \mathcal{I} \int d^3 v \frac{m v^2}{2} g_s \mathcal{J}_0 \phi^*. \qquad (3)$$

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Mixing-length estimate of fluxes

$$\Gamma_{k_{y}}^{\text{QL}} = \frac{\Gamma_{k_{y}}\gamma_{k_{y}}}{\langle |\phi_{k_{y}}|^{2}\rangle k_{y}^{2}}$$
(4)

Introduction

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- Magnetic configurations specified by n_{fp} and boundary Fourier coefs.
- Gradients set to promote ITG $\frac{d \ln T_i}{d\psi_t} = -5,$ $\frac{d \ln T_e}{d\psi_t} = \frac{d \ln n}{d\psi_t} = -0.1$
- Trace Ar+16 (Z = 16, A = 40) with gradients: $\frac{d \ln T_z}{d\psi_t} = -2.75,$ $\frac{d \ln n_z}{d\psi_t} = -0.42$







1D scan in a boundary Fourier mode



Comparison of different configurations

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1D scan in a boundary Fourier mode



Comparison of different configurations





- We have investigated a mixing-length estimate for impurity and energy fluxes from linear gyrokinetic simulations
- Impurity and energy flux are found to generally be correlated.
 However, we have found some configurations with twice the impurity flux per energy flux.
- Optimization on one field-line is probably not sufficient.
- Nonlinear verification of the optimization is ongoing.



