

## Motivation

- Forces  $\propto B^2$ . High *B* limited by support structure.
- Superconductor quench limits depend on local **B**.
- Need to be able to dissipate stored energy  $W = \frac{1}{2}LI^2$ .
- Coil shapes can probably be optimized for these quantities.

Field and force on coil 2 due to current in coil 1 can be computed quickly: 1D filament models are ok.

eld:  $\mathbf{B}(\mathbf{r}) = \frac{\mu_0 I}{4\pi} \int d\tilde{\phi} \frac{d\tilde{\mathbf{r}}}{|\mathbf{r} - \tilde{\mathbf{r}}|^3} \times (\mathbf{r} - \tilde{\mathbf{r}})$ Tricky part is the self-field: singularity in **Biot-Savart Law:** 

Simply skipping the singular point in a 1D filament calculation gives a non-converging result with significant error:





Also apparent from analytic formulas for a circular coil:



 $\frac{d\mathbf{F}}{d\ell} = \frac{\mu_0 I^2}{4\pi R} \left[ \ln\left(\frac{8R}{a}\right) - \frac{3}{4} \right] \mathbf{e}_R$ 

Diverges if minor radius  $a \rightarrow 0$ 

Can we avoid high-dimensional integrals or PDE solve? Field: 3D integral

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int d^3 \tilde{r} \frac{\mathbf{J}(\tilde{\mathbf{r}}) \times (\mathbf{r} - \tilde{\mathbf{r}})}{|\mathbf{r} - \tilde{\mathbf{r}}|^3}$$

Force per unit length: 5D integral

$$\frac{d\mathbf{F}}{d\ell} = \frac{\mu_0}{4\pi} \frac{d\phi}{d\ell} \int dx \int dy \int d^3 \tilde{r} \sqrt{g} \frac{\mathbf{J}(\mathbf{r}) \times [\mathbf{J}(\tilde{\mathbf{r}}) \times (\mathbf{r} - \tilde{\mathbf{r}})]}{|\mathbf{r} - \tilde{\mathbf{r}}|^3}$$

Self-inductance & stored energy: 6D integral

 $L = \frac{\mu_0}{4\pi I^2} \int d^3r \int d^3\tilde{r} \frac{\mathbf{J}(\mathbf{r}) \cdot \mathbf{J}(\tilde{\mathbf{r}})}{1}$ 

# **Reduced models of self-force, stored energy, & critical current for** design optimization of electromagnets

# **Derivation of reduced models**

• Parameterize the conductor volume:



 $\mathbf{r}(\boldsymbol{\rho},\boldsymbol{\theta},\boldsymbol{\phi}) = \mathbf{r}_{c}(\boldsymbol{\phi}) + \boldsymbol{\rho}a\cos\theta\mathbf{n}(\boldsymbol{\phi}) + \boldsymbol{\rho}a\sin\theta\mathbf{b}(\boldsymbol{\phi})$ 

- Expansion parameter:  $a / R \ll 1$ , where  $R \sim$  scales of curve centerline, and  $b \sim a$ .
- Introduce intermediate scale d, with  $a \ll d \ll R$ . Split integrals into "near part" + "far part".
- Far part defined by  $|\mathbf{r} \tilde{\mathbf{r}}| > d$ . Finite cross-section can be neglected.
- Near part defined by  $|r \tilde{r}| < d$ . Coil centerline can be Taylor-expanded, so integrals can be done explicitly.
- Identify a 1D integral that has the same near part and far part as the above "high fidelity" calculation for  $a / R \ll 1$ .



$$\begin{aligned} \mathbf{Self-field} & \mathbf{B} = \mathbf{B}_{reg} + \mathbf{B}_0 + \mathbf{B}_{\kappa} + \mathbf{B}_b, \\ \mathbf{B}_{reg}(\phi) &= \frac{\mu_0 I}{4\pi} \int_0^{2\pi} d\tilde{\phi} \frac{\tilde{\mathbf{r}}'_c \times (\mathbf{r}_c - \tilde{\mathbf{r}}_c)}{\left(|\mathbf{r}_c - \tilde{\mathbf{r}}_c|^2 + \delta ab\right)^{3/2}} & \mathbf{B}_0 &= \frac{\mu_0 I}{4\pi ab} \sum_{s_u, s_v} s_u s_v \left[G\left(b(v - s_v), a(u - s_u)\right)\mathbf{q} - G\left(a(u - s_u), b(v - s_v)\right)\right] \mathbf{p} \\ \delta &= \exp\left(-\frac{25}{6} + k\right) & G(x, y) = y \tan^{-1} \frac{x}{y} + \frac{x}{2} \ln\left(1 + \frac{y^2}{x^2}\right) \\ \mathbf{B}_{\kappa} &= \frac{4b}{3a} \tan^{-1} \frac{a}{b} + \frac{4a}{3b} \tan^{-1} \frac{b}{a} + \frac{b^2}{6a^2} \ln \frac{b}{a} + \frac{a^2}{6b^2} \ln \frac{a}{b} \\ &- \frac{a^4 - 6a^2b^2 + b^4}{6a^2b^2} \ln\left(\frac{a}{b} + \frac{b}{a}\right) & K(U, V) = -2UV\left(\kappa_1 \mathbf{q} - \kappa_2 \mathbf{p}\right) \ln\left(\frac{aU^2}{b} + \frac{bV^2}{a}\right) \\ &+ \left(\kappa_2 \mathbf{q} - \kappa_1 \mathbf{p}\right) \left(\frac{aU^2}{b} + \frac{bV^2}{a}\right) \ln\left(\frac{aU^2}{b} + \frac{bV^2}{a}\right) \\ &+ \frac{4aU^2\kappa_2 \mathbf{p}}{b} \tan^{-1} \frac{bV}{aU} - \frac{4bV^2\kappa_1 \mathbf{q}}{a} \tan^{-1} \frac{aU}{bV}, \end{aligned}$$

$$\begin{aligned} \mathbf{\hat{b}elf-field} & \mathbf{B} = \mathbf{B}_{reg} + \mathbf{B}_0 + \mathbf{B}_\kappa + \mathbf{B}_b, \\ \mathbf{B}_{reg}(\phi) &= \frac{\mu_0 I}{4\pi} \int_0^{2\pi} d\tilde{\phi} \frac{\tilde{\mathbf{r}}'_c \times (\mathbf{r}_c - \tilde{\mathbf{r}}_c)}{\left(|\mathbf{r}_c - \tilde{\mathbf{r}}_c|^2 + \delta ab\right)^{3/2}} \\ \delta &= \exp\left(-\frac{25}{6} + k\right) \\ k &= \frac{4b}{3a} \tan^{-1} \frac{a}{b} + \frac{4a}{3b} \tan^{-1} \frac{b}{a} + \frac{b^2}{6a^2} \ln \frac{b}{a} + \frac{a^2}{6b^2} \ln \frac{a}{b} \\ &- \frac{a^4 - 6a^2b^2 + b^4}{6a^2b^2} \ln \left(\frac{a}{b} + \frac{b}{a}\right) \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} \mathbf{b}_b = \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} (4 + 2\ln 2 + \ln \delta) \end{aligned} \\ \mathbf{B}_b &= \frac{\mu_0 I \kappa \mathbf{b}}{8\pi} \mathbf{b}_b = \frac{\mu_0 I \kappa \mathbf{b}}$$

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$$\mathbf{B}_{b} = \frac{\mu_{0} I \kappa \mathbf{b}}{8\pi} \left(4 + 2\ln 2 + \ln \delta\right)$$

### High fidelity |B| [Tesla] Reduced model |B| [Tesla] High fidelity $B_z$ [Tesla]







High

fidelity

model

$$\mathbf{r}(u,v,\phi) = \mathbf{r}_{c}(\phi) + \frac{au}{2}\mathbf{p}(\phi) + \frac{bv}{2}\mathbf{q}(\phi)$$

1 →	$\mathbf{B} = \mathbf{B}_{near} + \mathbf{B}_{far}$	
	Far contribution: $ \boldsymbol{r} - \tilde{\boldsymbol{r}}  > d$	
	Near contribution: $ \boldsymbol{r} - \tilde{\boldsymbol{r}}  < d$	

Reduced model  $B_z$  [Tesla]



estimate the global critical current as

$$I_c = \min_{\phi} r_{-s}$$

### References Supported by DE-FG02-93ER54 the Simons Foundation (560651



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	M Landreman, S Hurwitz, & T Antonsen, arXiv:2310.12087