

Efficient Calculation of the Self-Inductance, Self-Force, and Internal Magnetic Field for Thin Electromagnetic Coils

Siena Hurwitz¹, Matt Landreman¹, Thomas M. Antonsen¹

¹University of Maryland, College Park



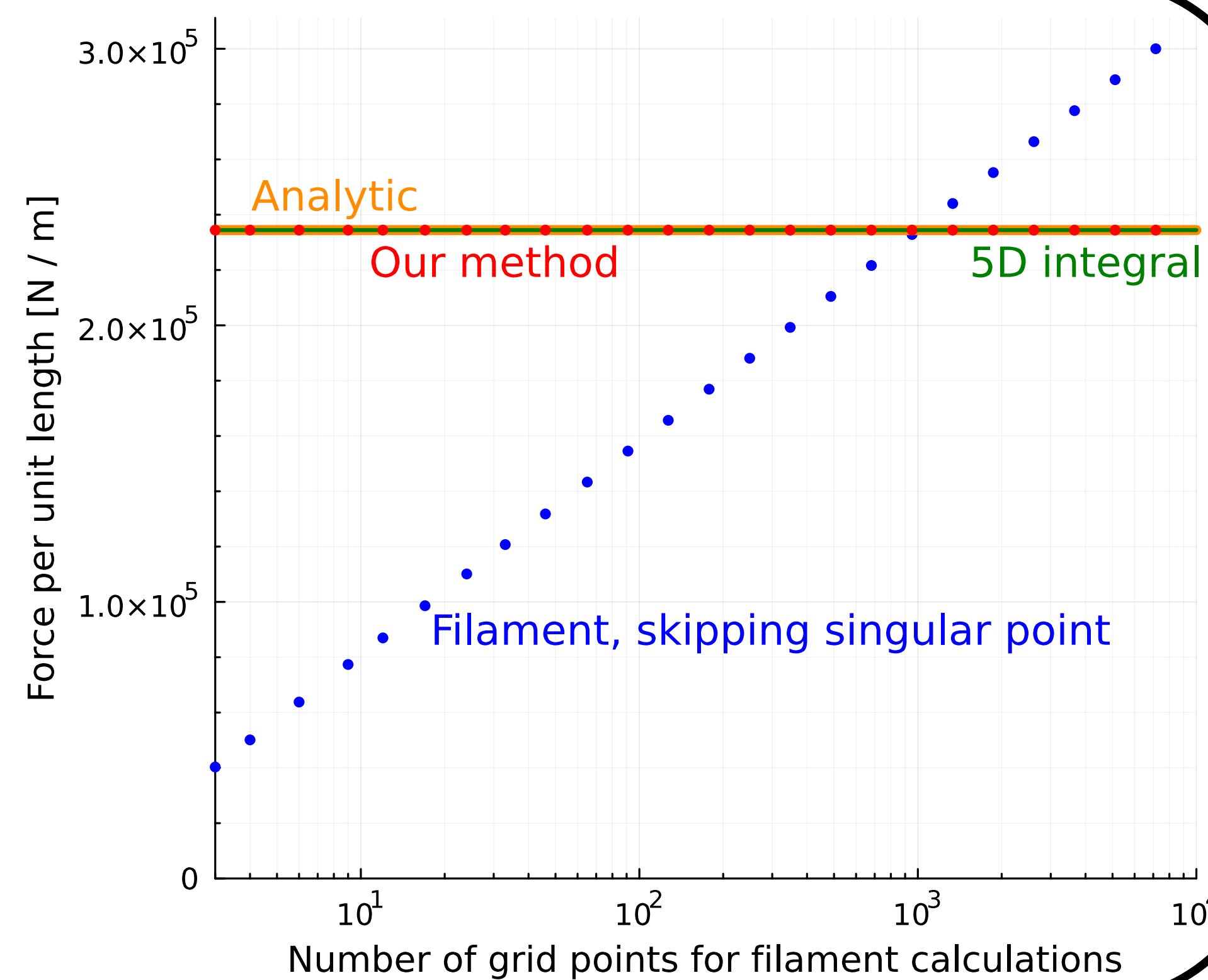
Motivations

- Quantities of high-field coils such as the self-inductance, self-force, and internal magnetic field may be important in coil design, and potential applications include optimization of stellarator magnets for low stresses and high critical current.
- The finite-thickness formulae are inefficient to compute due to high dimensionality of the integrals and a singularity at $\tilde{r} = r$, e.g.

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int d^3\tilde{r} \frac{\tilde{\mathbf{J}} \times \Delta\mathbf{r}}{|\Delta\mathbf{r}|^3}, \quad \Delta\mathbf{r} \equiv \mathbf{r} - \tilde{r}_c$$

- One cannot take the zero-thickness limit as the results diverge logarithmically. For a circular wire loop with radii R and a ,

$$\frac{dF_o}{dl} = \frac{\mu_0 I^2}{4\pi R} \left[\ln \frac{8R}{a} - \frac{3}{4} \right] \hat{\mathbf{r}}$$



Coils, Coordinate System, and Exact Model

We examine closed coils with a circular cross-section of minor radius, a , and uniform current, I . The coils are described by a center-line, $\mathbf{r}_c(\phi)$, where $\phi \in [0, 2\pi)$. The Frenet-Serret unit vectors can then be found for \mathbf{r}_c :

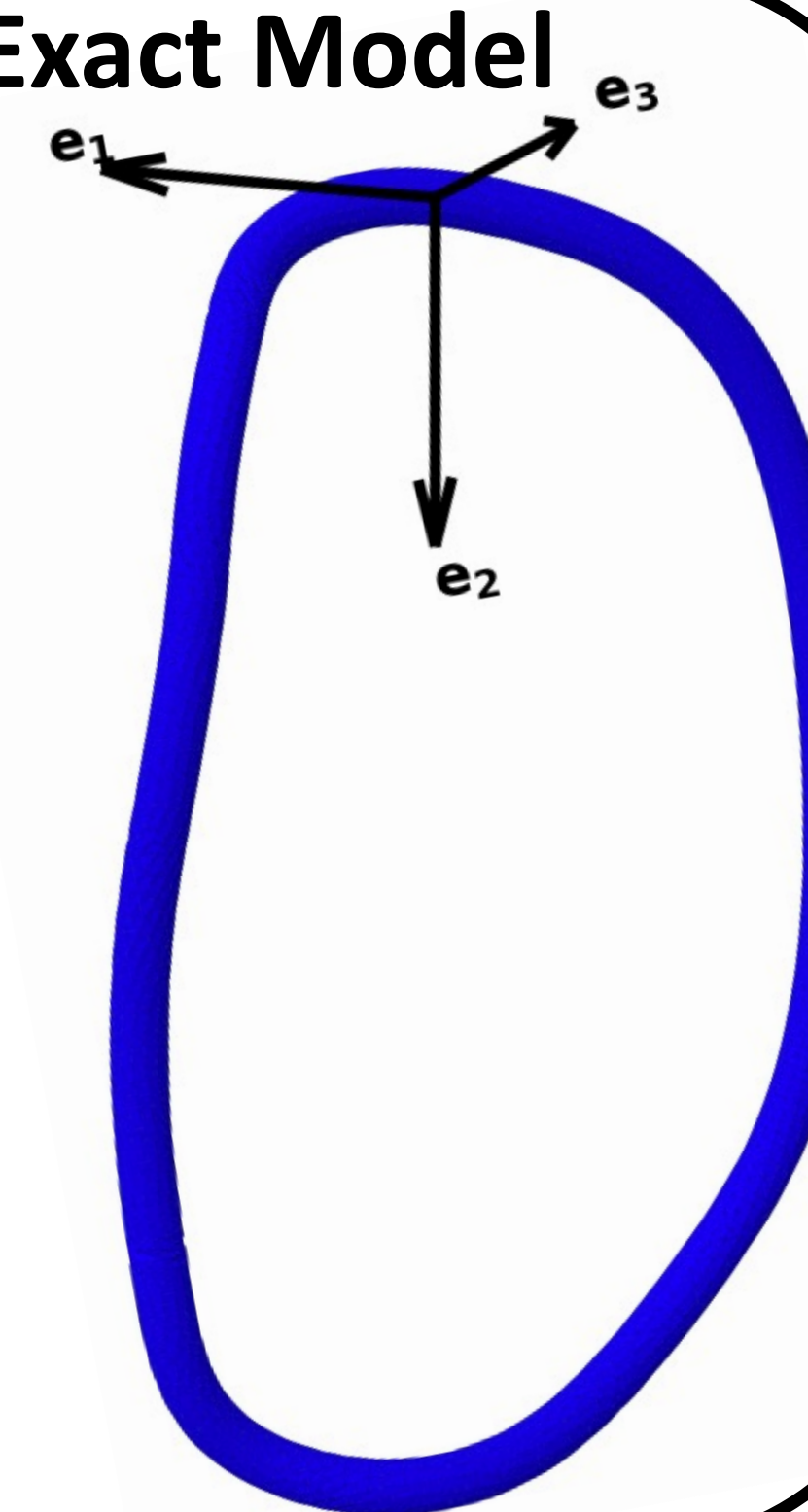
$$\mathbf{e}_1 = \mathbf{r}'_c / |\mathbf{r}'_c|, \mathbf{e}_2 = \mathbf{e}'_1 / |\mathbf{e}'_1|, \mathbf{e}_3 = \mathbf{e}_1 \times \mathbf{e}_2$$

We form the coordinate system, $\mathbf{r}(s, \theta, \phi) = \mathbf{r}_c + s \cos \theta \mathbf{e}_2 + s \sin \theta \mathbf{e}_3$, where $\theta \in [0, 2\pi)$ and $s \geq 0$. This gives:

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0 I}{4\pi a^2} \int_0^{2\pi} d\tilde{\phi} \int_0^{2\pi} d\tilde{\theta} \int_0^a ds \tilde{s} (1 - \tilde{\kappa}\tilde{s}) \frac{\tilde{\mathbf{r}}'_c}{|\Delta\mathbf{r}|}$$

$$L = \frac{1}{\pi a^2 I} \int_0^{2\pi} d\phi \int_0^{2\pi} d\theta \int_0^a ds s (1 - \kappa s) |\mathbf{r}'_c| (\mathbf{e}_1 \cdot \mathbf{A})$$

$$\frac{d\mathbf{F}}{dl}(\phi) = \frac{I}{\pi a^2} \int_0^{2\pi} d\theta \int_0^a ds s (1 - \kappa s) (\mathbf{e}_1 \times \mathbf{B})$$



Derivation Overview

- Assume $a/\mathcal{L} \ll 1$, where \mathcal{L} is any length scale of \mathbf{r}_c (e.g., $|\mathbf{r}'_c|, \kappa^{-1}, \tau^{-1}$, etc).
- Partition the integral over $\tilde{\phi}$ in $\mathbf{A}(\mathbf{r})$ into two sections about ϕ_0 , where ϕ_0 satisfies $a/\mathcal{L} \ll \phi_0 \ll 1$.
 - The "near region" ($|\phi - \phi_0| < \phi_0$) satisfies $|\Delta\phi| \ll 1$.
 - The "far region" ($|\phi - \phi_0| > \phi_0$) satisfies $a/|\Delta\mathbf{r}| \ll 1$.
- Solve for $\mathbf{A}(\mathbf{r})$ and simplify under these assumptions.
- Using this new formula, obtain desired formulae by
 - positing a regularized form of the magnetic field, \mathbf{B}_{reg}
 - calculating $\mathbf{B} = \nabla \times \mathbf{A}$, comparing to \mathbf{B}_{reg} , and substituting in \mathbf{B}_{reg} where appropriate
 - applying previous two methods to the formulae for L and $d\mathbf{F}/dl$

Self-Magnetic Field

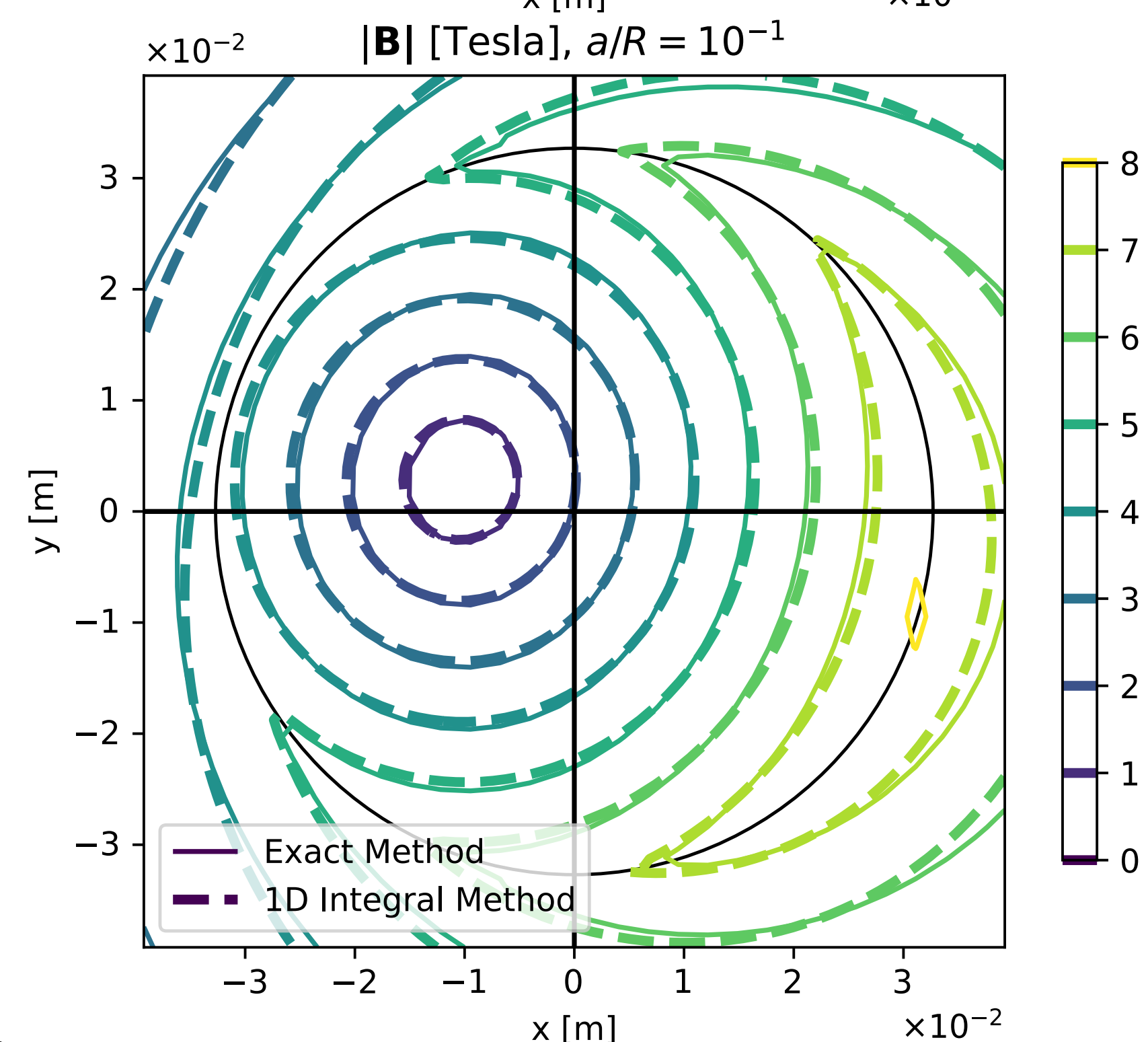
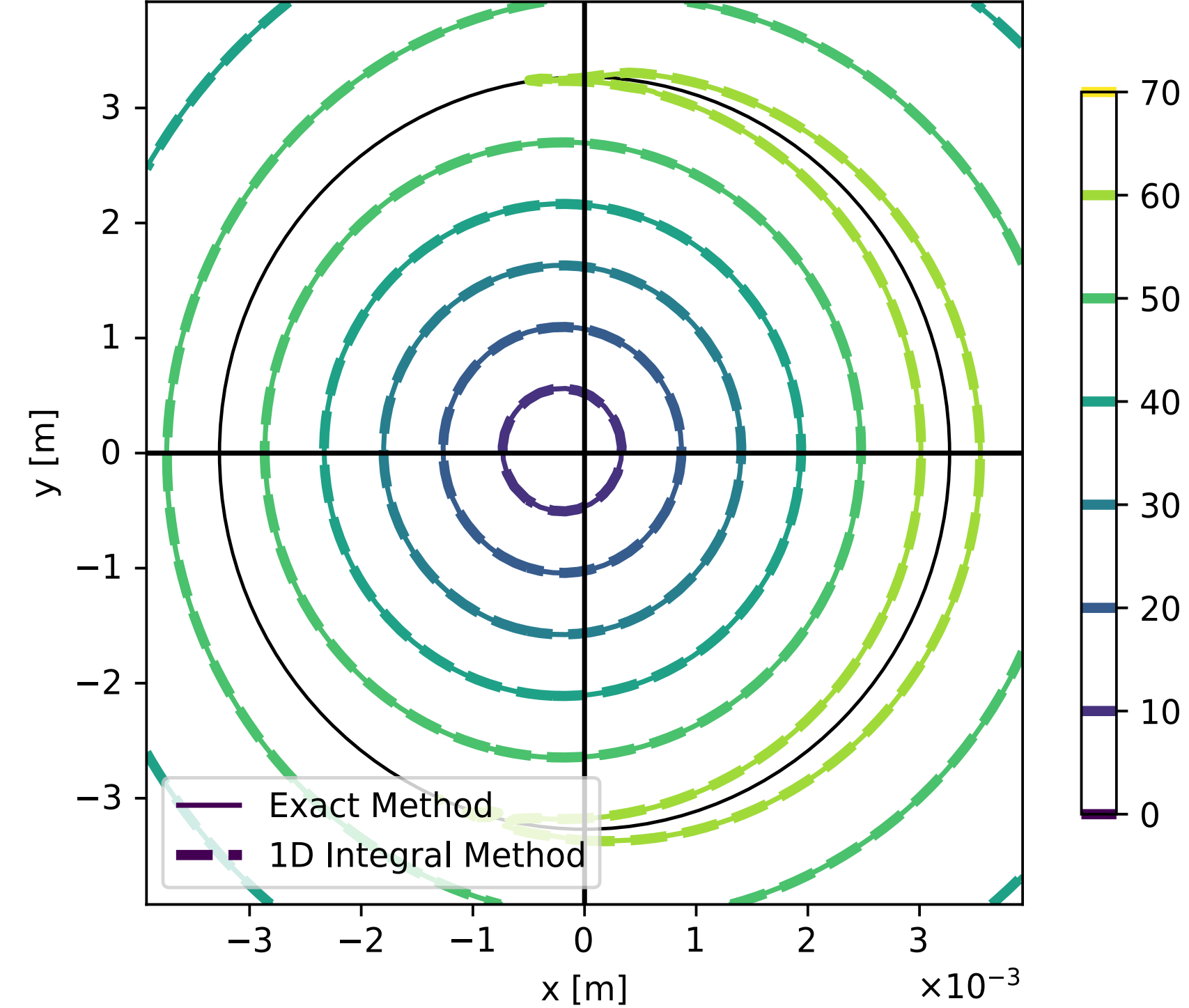
$$\mathbf{B}(\mathbf{r}) = \mathbf{B}_{reg} + \mathbf{B}_{cyl} + \mathbf{B}_{hi}$$

$$\mathbf{B}_{reg} \equiv \frac{\mu_0 I}{4\pi} \int_0^{2\pi} d\tilde{\phi} \frac{\tilde{\mathbf{r}}'_c \times \Delta\mathbf{r}_c}{(|\Delta\mathbf{r}_c|^2 + a^2/\sqrt{e})^{3/2}}$$

$$\mathbf{B}_{cyl} \equiv (-\sin \theta \mathbf{e}_2 + \cos \theta \mathbf{e}_3) \frac{\mu_0 I}{2\pi} \begin{cases} s/a^2, & s < a \\ 1/a, & s > a \end{cases}$$

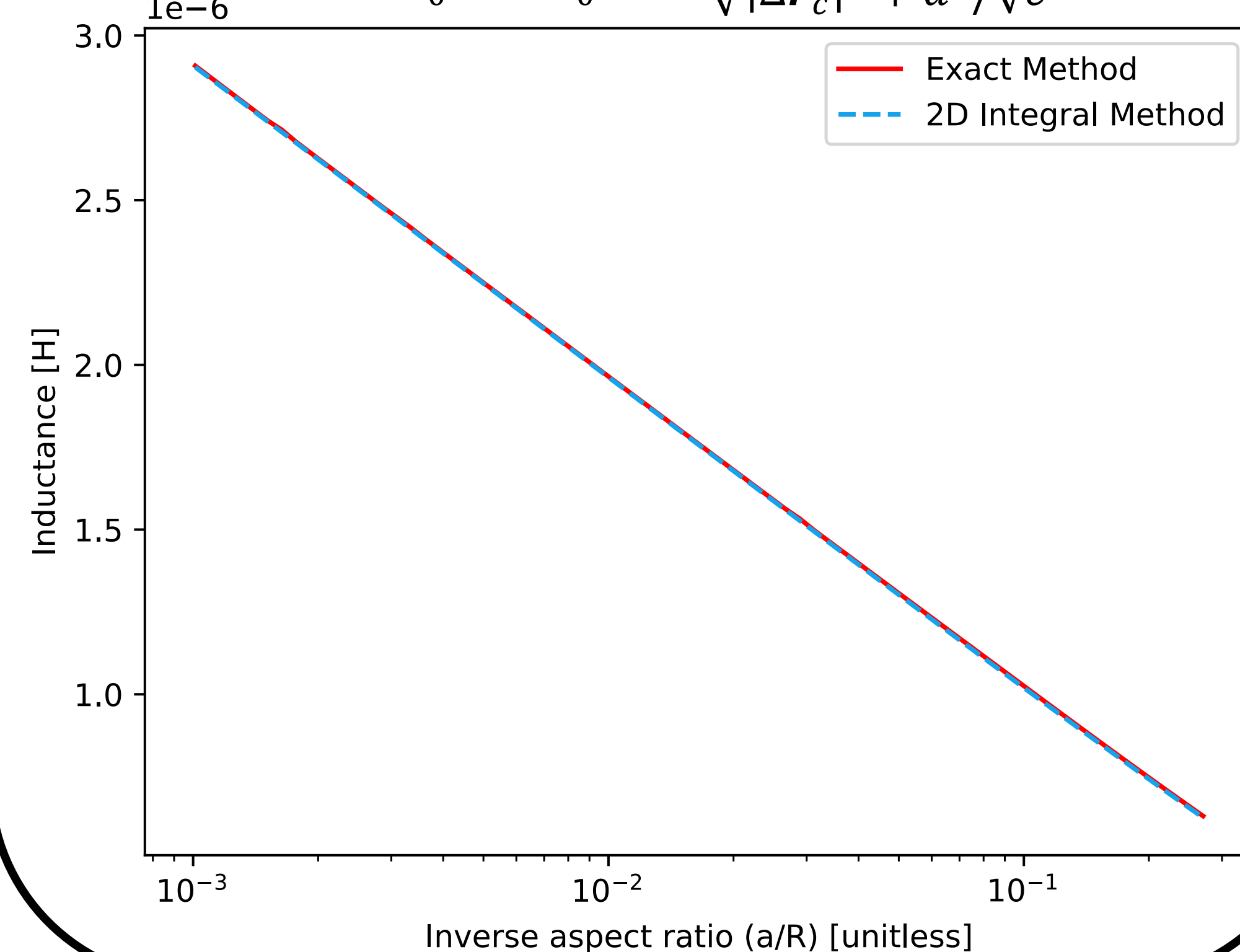
$$\mathbf{B}_{hi} \equiv \frac{\mu_0 I \kappa}{8\pi} \begin{cases} -\frac{s^2 \sin 2\theta}{2a^2} \mathbf{e}_2 + \left(\frac{3}{2} + \frac{s^2}{a^2} \left(\frac{\cos 2\theta}{2} - 1 \right) \right) \mathbf{e}_3, & s < a \\ \left(\frac{a^2}{2s^2} - 1 \right) \sin 2\theta \mathbf{e}_2 + \left(\frac{1}{2} - 2 \ln \frac{s}{a} - \left(\frac{a^2}{2s^2} - 1 \right) \cos 2\theta \right) \mathbf{e}_3, & s > a \end{cases}$$

$$|\mathbf{B}| \text{ [Tesla]}, a/R = 10^{-2} \quad \leftarrow R \equiv \int_0^{2\pi} d\phi |\mathbf{r}'_c| / 2\pi$$



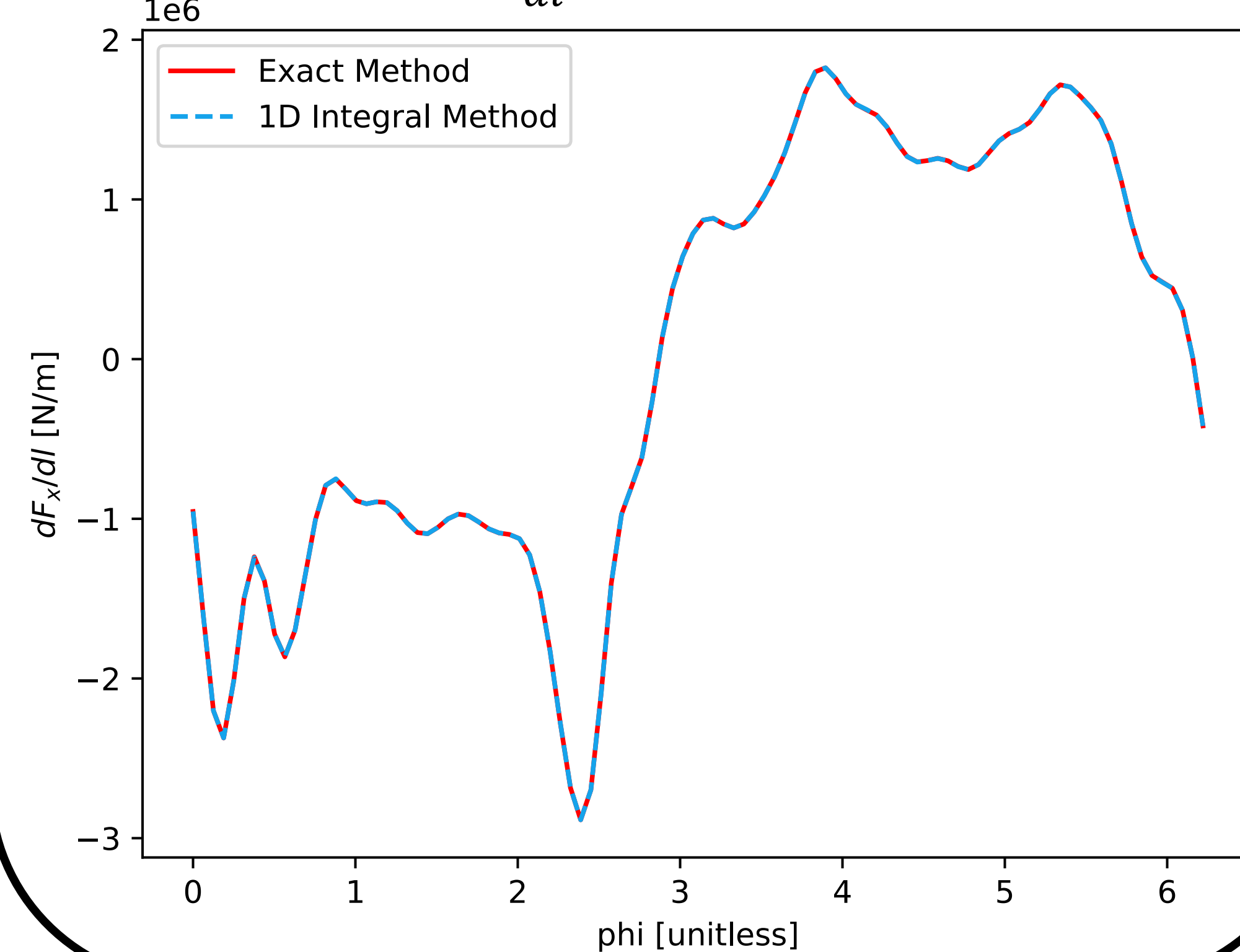
Self-Inductance

$$L = \frac{\mu_0}{4\pi} \int_0^{2\pi} d\tilde{\phi} \int_0^{2\pi} d\phi \frac{\mathbf{r}'_c \cdot \tilde{\mathbf{r}}'_c}{\sqrt{|\Delta\mathbf{r}_c|^2 + a^2/\sqrt{e}}}$$



Self-Force

$$\frac{d\mathbf{F}}{dl} = I \mathbf{e}_1 \times \mathbf{B}_{reg}$$



This work was supported by a grant from the Simons Foundation (No. 560651, T. A.). This work was also supported by the U.S. Department of Energy under Contract DE-FG02-93ER54197.

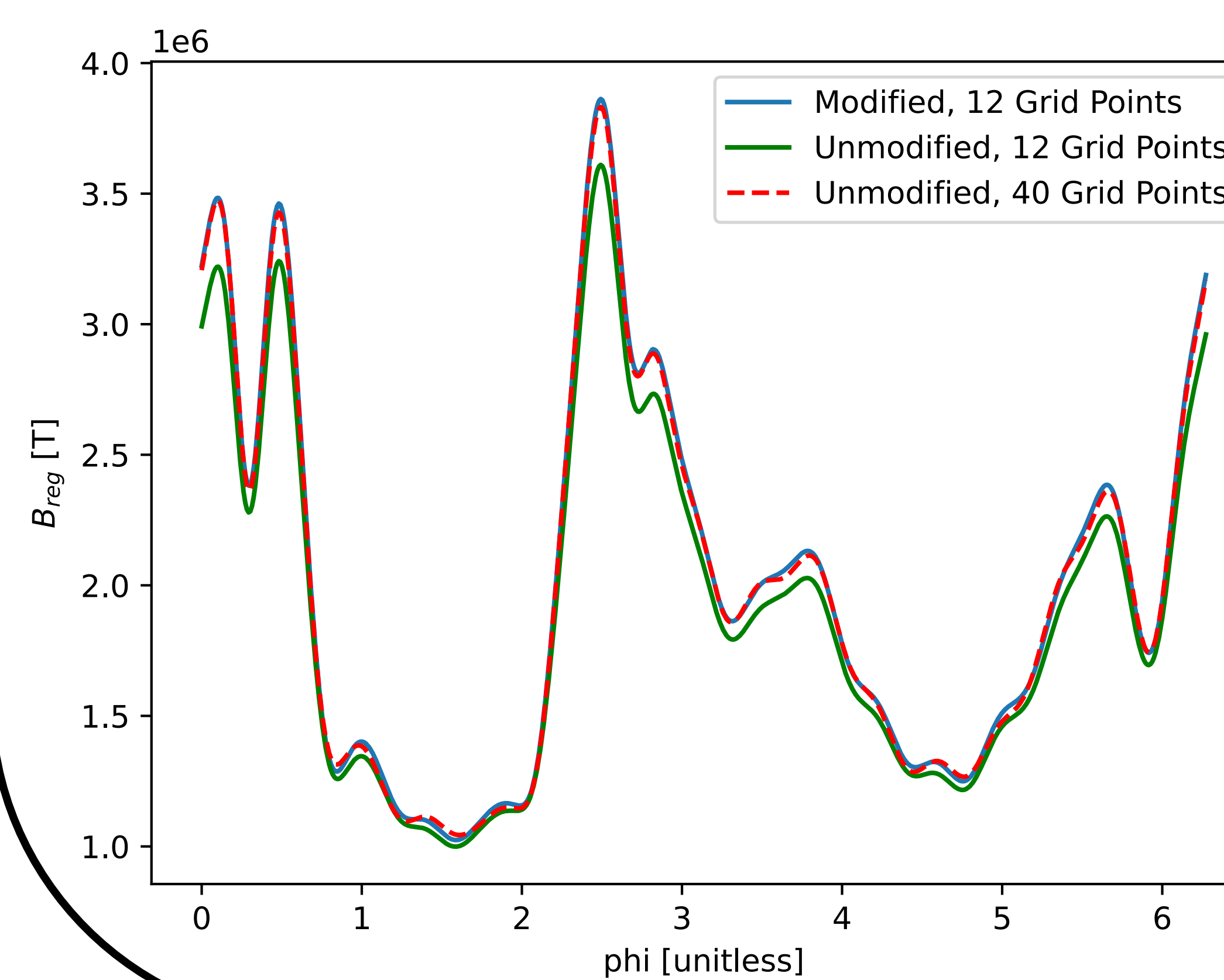
Modified Methods for Self-Force and Self-Inductance

- Evaluating $d\mathbf{F}/dl$ and L by quadrature for $a/\mathcal{L} \ll 1$ is inefficient as the respective integrands become sharp as $a \rightarrow 0$.
- Instead, modify the equations for \mathbf{B}_{reg} and L by:
 - finding a quantity, Λ , with a known antiderivative that fits the integrand well near $\tilde{r} = r$.
 - subtracting Λ from the integrand, then adding it back in separately and analytically integrating this piece

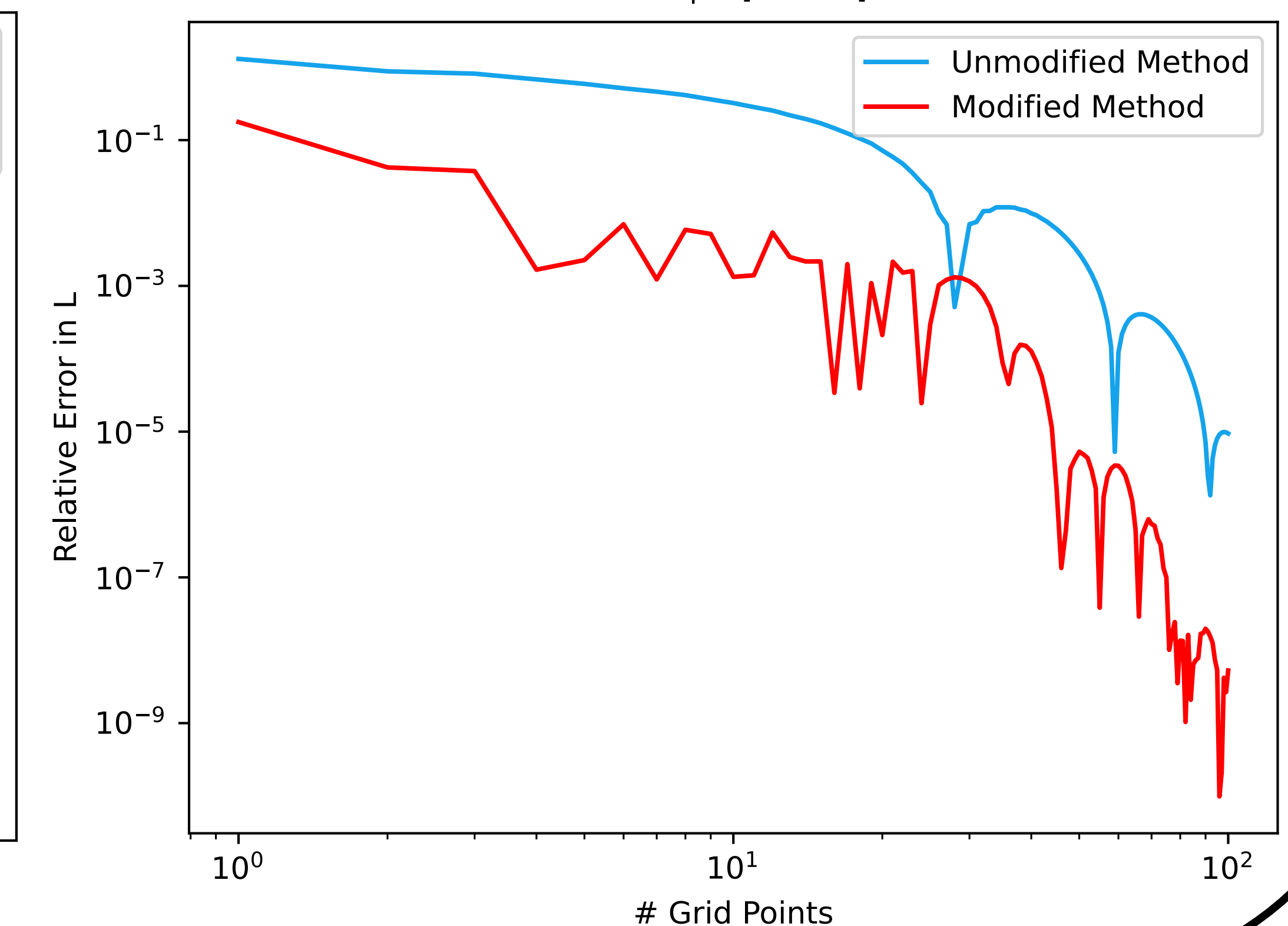
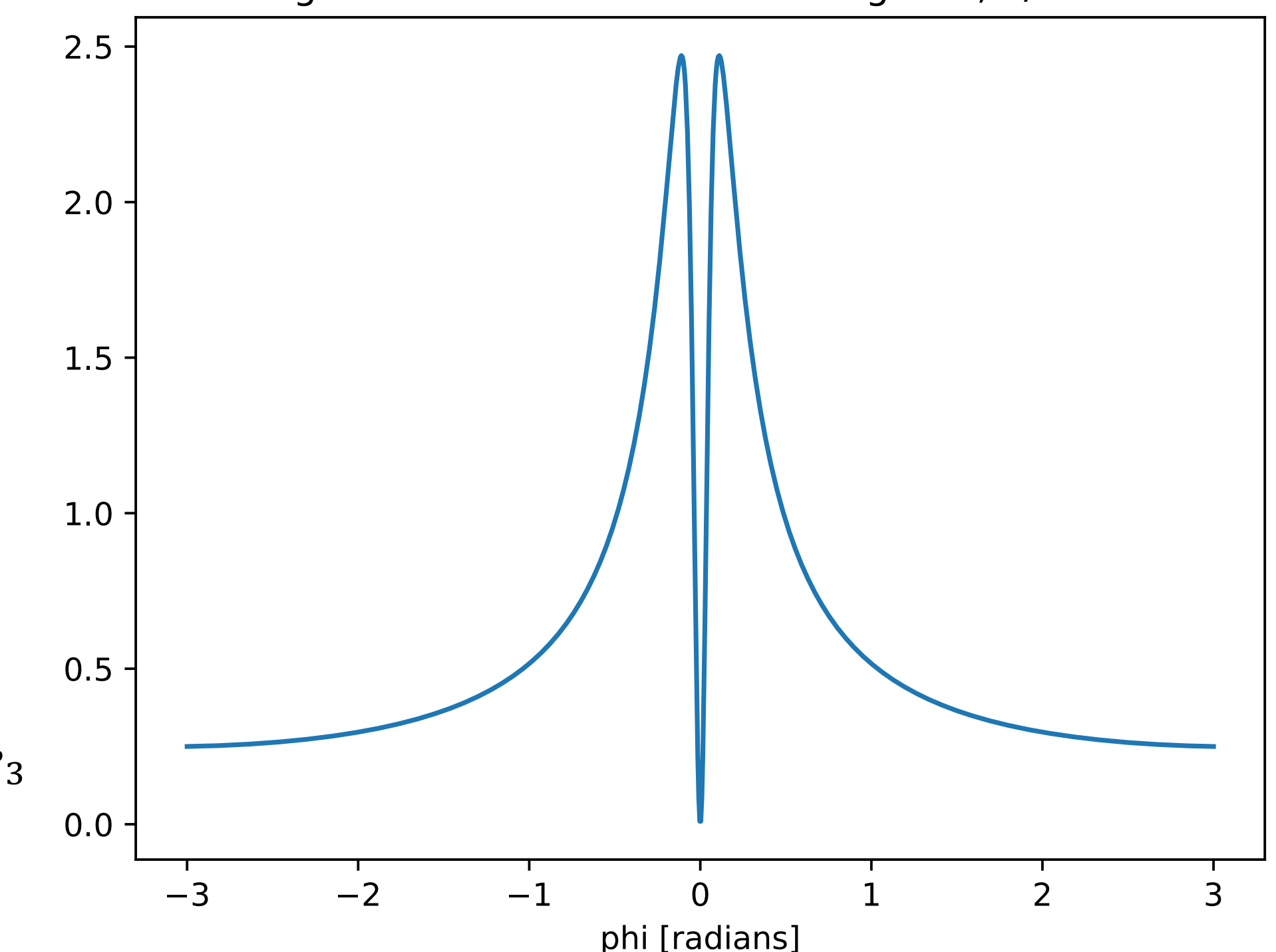
- For the force,

$$\Lambda \equiv \frac{\kappa(1 - \cos \Delta\phi) \mathbf{e}_3}{2^{3/2} \left((1 - \cos \Delta\phi) + a^2/2\sqrt{e} |\mathbf{r}'_c|^2 \right)^{3/2}}$$

$$\Rightarrow \mathbf{B}_{reg} = \frac{\mu_0 I}{4\pi} \int_0^{2\pi} d\tilde{\phi} \left(\frac{\tilde{\mathbf{r}}'_c \times \Delta\mathbf{r}_c}{(|\Delta\mathbf{r}_c|^2 + a^2/\sqrt{e})^{3/2}} - \Lambda \right) + \frac{\mu_0 I \kappa}{4\pi} \left(\ln \frac{8|\mathbf{r}'_c|}{a} - \frac{3}{4} \right) \mathbf{e}_3$$



Regularized Biot-Savart Law Integrand, $a/R=0.1$



References

- Hurwitz, S., Landreman, M., & Antonsen Jr, T. M. (2023). Efficient calculation of self magnetic field, self-force, and self-inductance for electromagnetic coils. *arXiv preprint arXiv:2310.12087*.
- Landreman, M., Hurwitz, S., & Antonsen Jr, T. M. (2023). Efficient calculation of self magnetic field, self-force, and self-inductance for electromagnetic coils. II. Rectangular cross-section. *arXiv preprint arXiv:2310.12087*.

