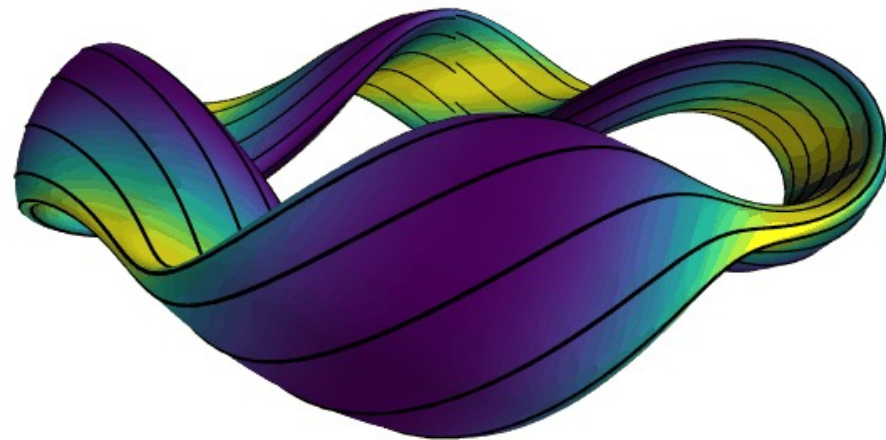
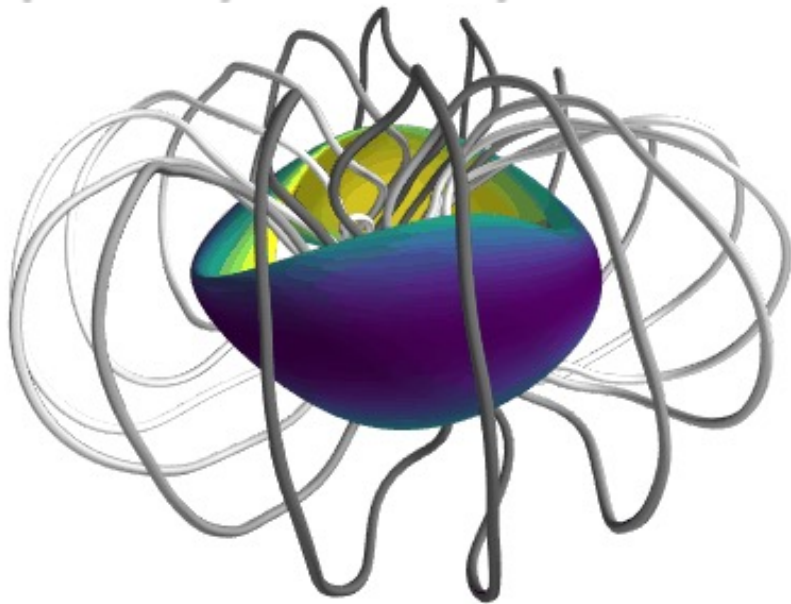


New stellarator configurations with precise quasisymmetry and energetic particle confinement

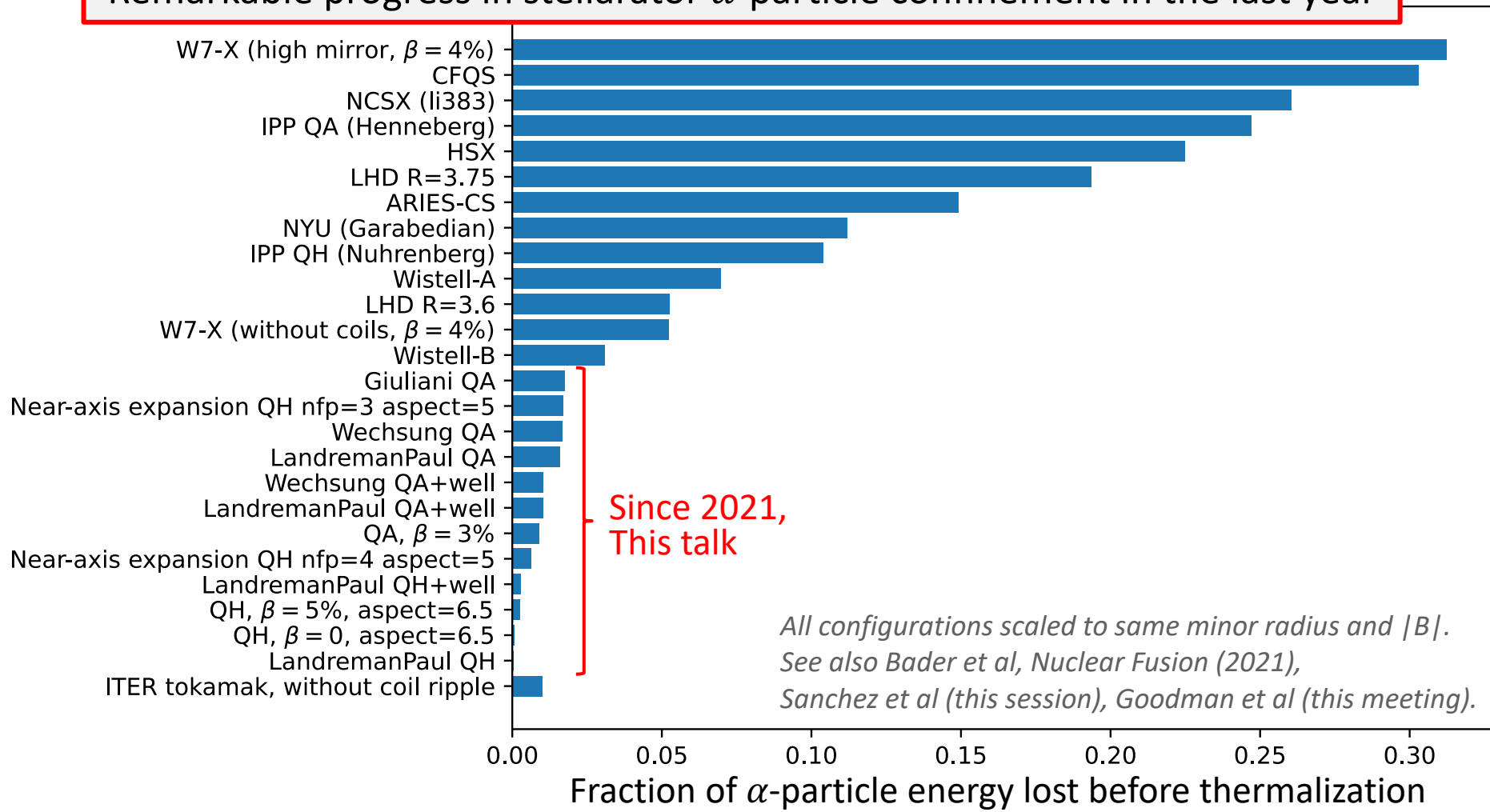


M Landreman^a, S Buller^a, A Cerfon^b, M Drevlak^c, A Giuliani^b, B Medasani^d,
M Padidar^e, E J Paul^d, G Stadler^b, F Wechsung^b, C Zhu^e

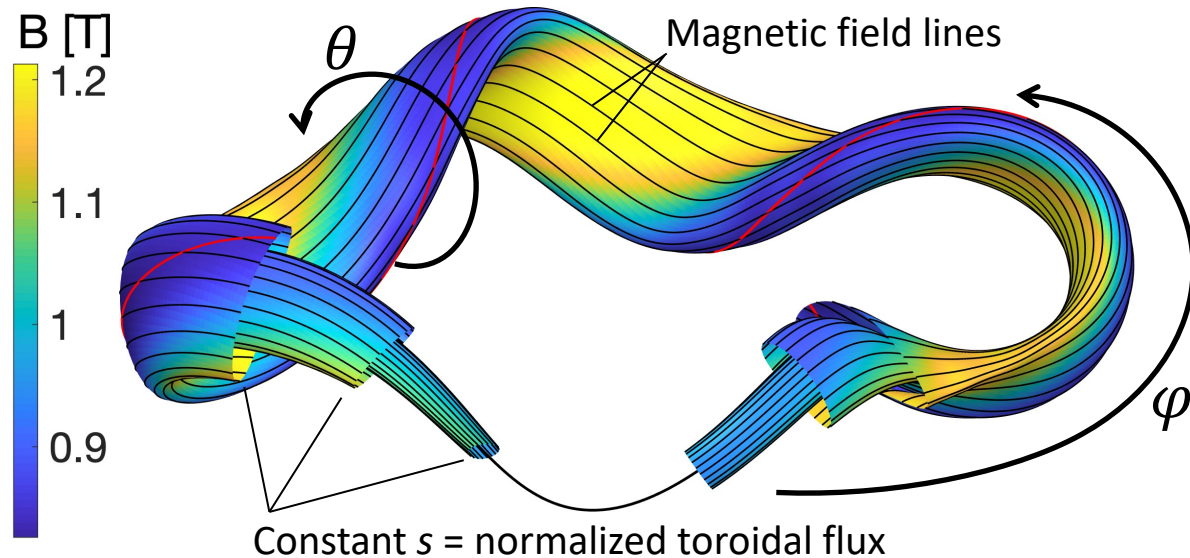
^a U of Maryland, ^b New York U, ^c Max Planck Institute, ^d PPPL, ^e U of Science & Technology of China, ^e Cornell

ML & Paul, PRL (2022), Wechsung et al, PNAS (2022), Giuliani et al arXiv (2022), ML, Buller, & Drevlak arXiv (2022)

Remarkable progress in stellarator α -particle confinement in the last year



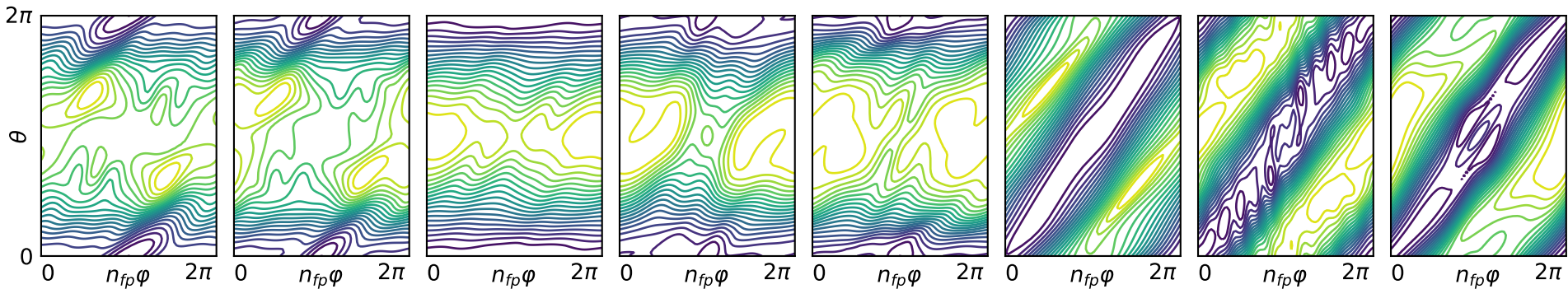
These new configurations with good alpha confinement use the principle of *quasisymmetry*.



$$B = B(s, \theta - N \varphi)$$

$\swarrow \quad \searrow$
 Boozer angles

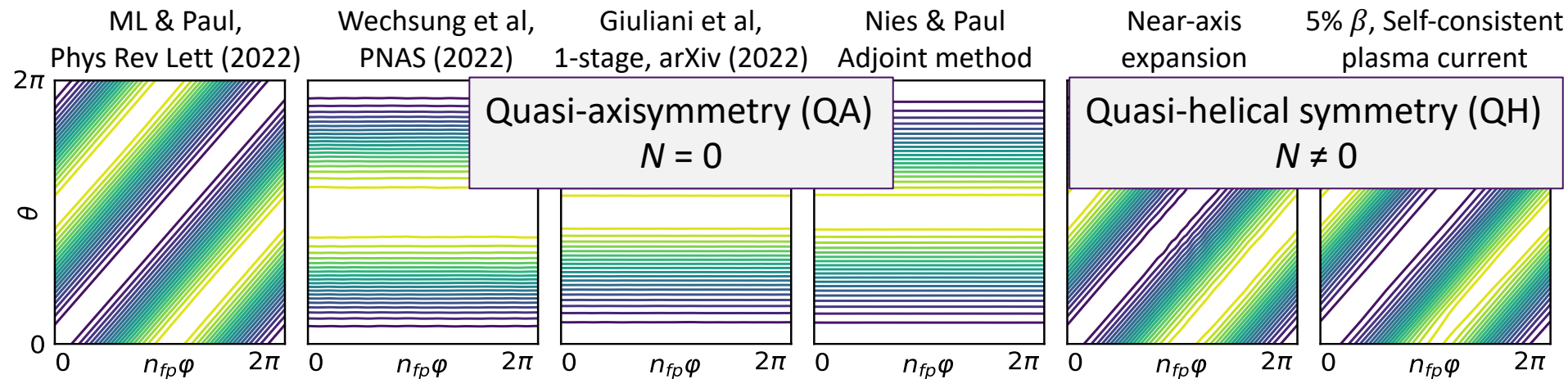
$$\Rightarrow \oint (\mathbf{v}_d \cdot \nabla s) dt = 0$$



Goal: $B = B(s, \theta - N \varphi)$



Since 2021



Also Dudt et al (this afternoon), Zarnstorff et al (yesterday)

- Minimal optimization recipe (low β)
- Self-consistent bootstrap current at high β
- Future directions

- Minimal optimization recipe (low β)
- Self-consistent bootstrap current at high β
- Future directions

Optimization problem

- 2 stage approach, as for W7-X: First optimize shape of boundary surface, then coils.

- Objective functions:

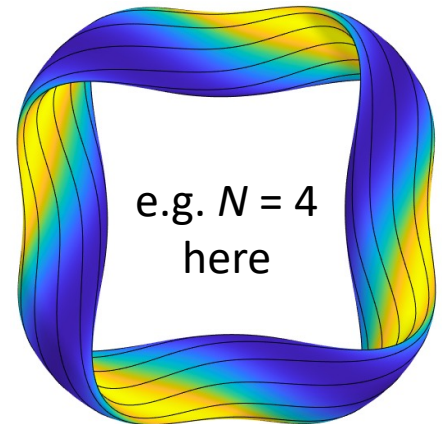
$$f_{QS} = \int d^3x \left(\frac{1}{B^3} \left[(N - \iota) \mathbf{B} \times \nabla B \cdot \nabla \psi - (G + NI) \mathbf{B} \cdot \nabla B \right] \right)^2$$

$$f_{QH} = \left(\underset{\substack{\uparrow \\ \text{Boundary aspect ratio}}}{A - A_*} \right)^2 + f_{QS} \qquad f_{QA} = \left(A - A_* \right)^2 + \left(\iota_* - \int_0^1 \iota ds \right)^2 + f_{QS}$$

Goal: $B = B(s, \theta - N \varphi)$.

For quasi-axisymmetry,
 $N = 0$.

For quasi-helical symmetry,
 N is the number of field periods,



Optimization problem

- 2 stage approach, as for W7-X: First optimize shape of boundary surface, then coils.

- Objective functions:

$$f_{QS} = \int d^3x \left(\frac{1}{B^3} \left[(N - \iota) \mathbf{B} \times \nabla B \cdot \nabla \psi - (G + NI) \mathbf{B} \cdot \nabla B \right] \right)^2$$

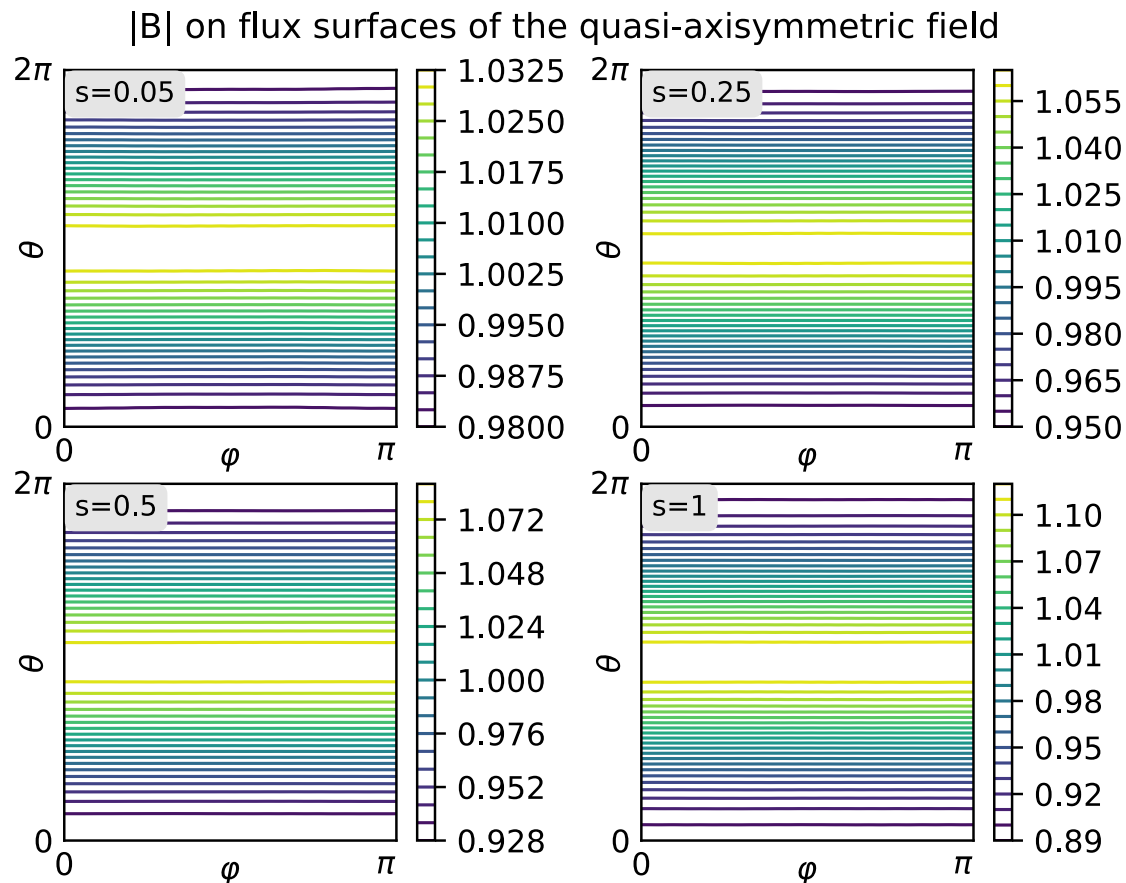
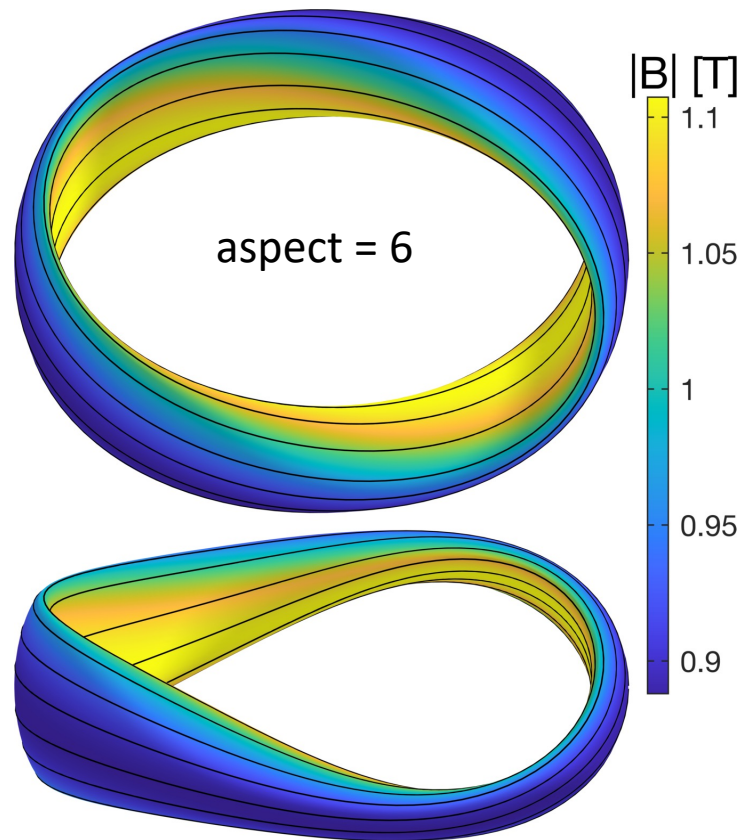
$$f_{QH} = \left(\underset{\substack{\uparrow \\ \text{Boundary aspect ratio}}}{A - A_*} \right)^2 + f_{QS} \qquad f_{QA} = \left(A - A_* \right)^2 + \left(\iota_* - \int_0^1 \iota ds \right)^2 + f_{QS}$$

- Parameter space: $R_{m,n}$ & $Z_{m,n}$ defining a toroidal boundary

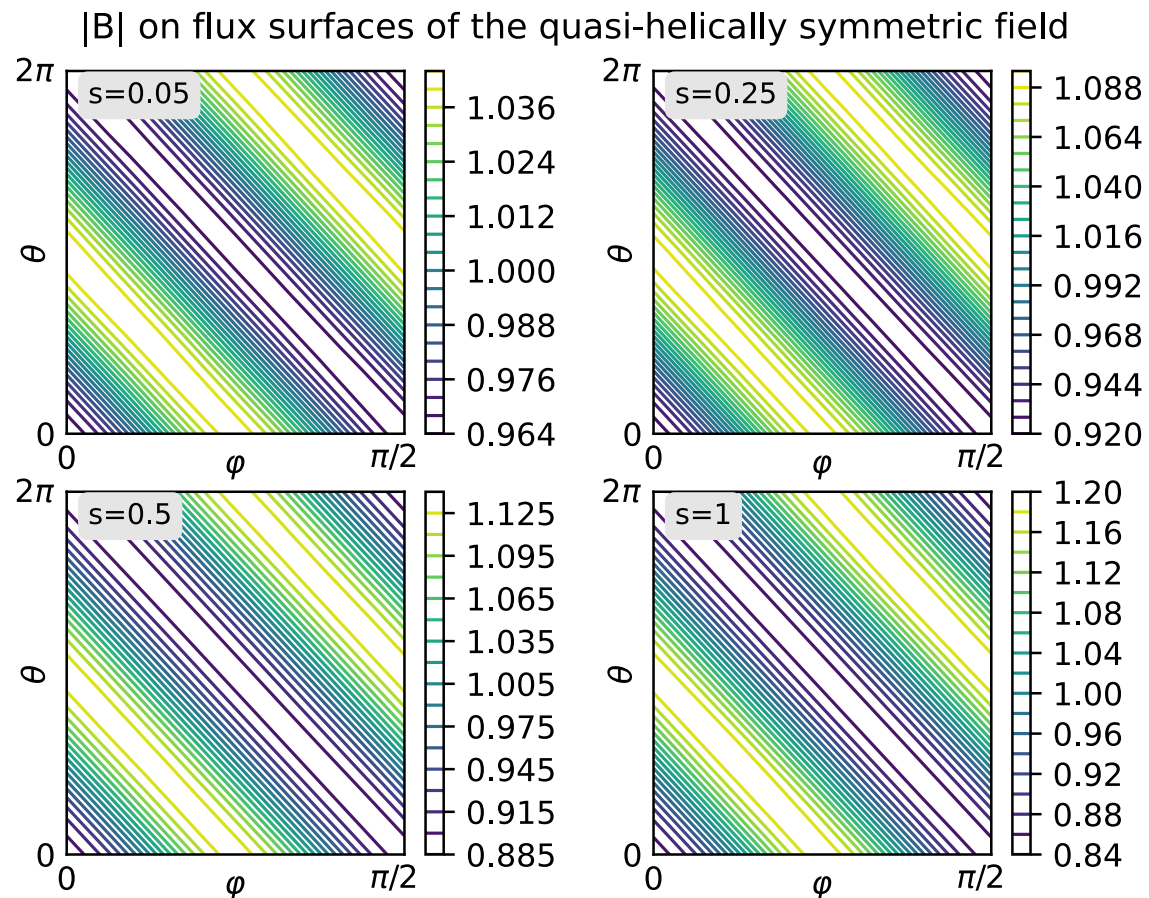
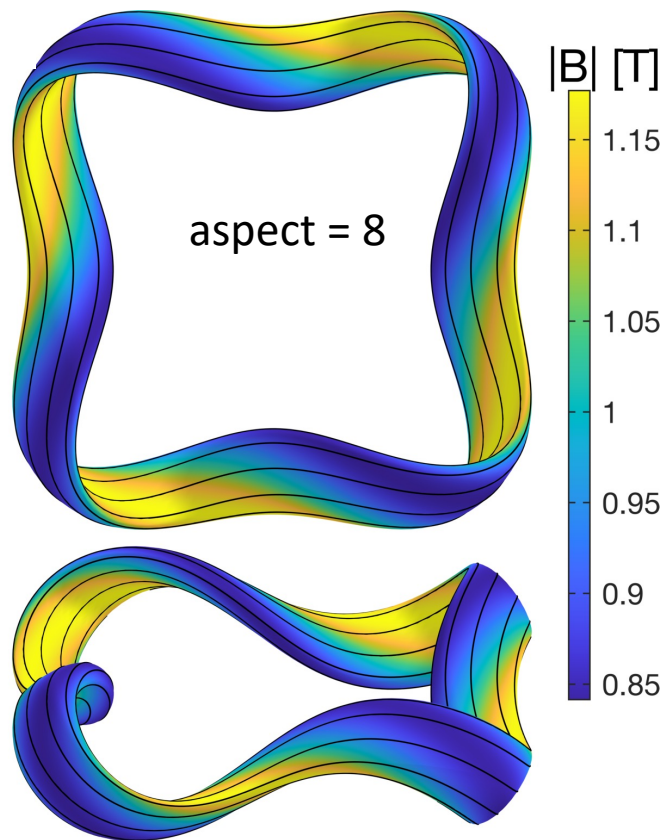
$$R(\theta, \phi) = \sum_{m,n} R_{m,n} \cos(m\theta - n\phi), \quad Z(\theta, \phi) = \sum_{m,n} Z_{m,n} \sin(m\theta - n\phi)$$

- Codes used: SIMSOPT with VMEC
- Cold start: circular cross-section torus
- Vacuum fields ($\nabla \times \mathbf{B} = 0$) at first, allowing precise checks
- Algorithm: default for least-squares in scipy (trust region reflective)
- 6 steps: increasing # of modes varied & equilibrium resolution
- Run many optimizations, pick the best

Straight $|B|$ contours are possible for quasi-axisymmetry

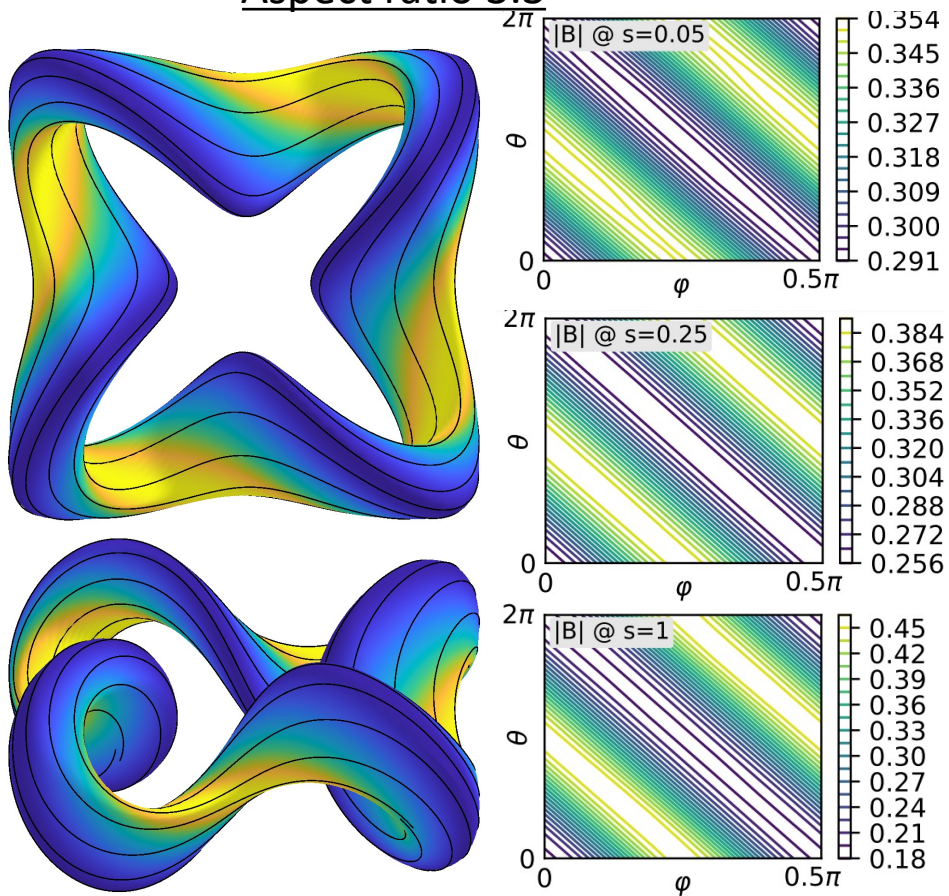


Straight $|B|$ contours are possible for quasi-helical symmetry



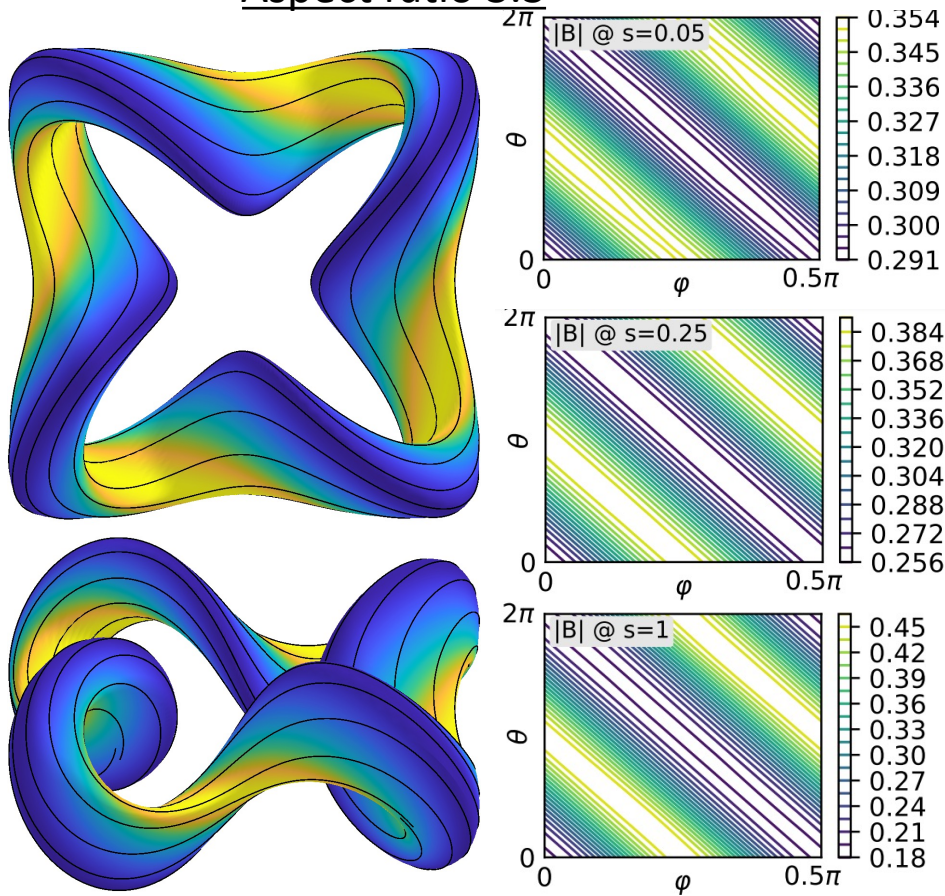
Nearly as good quasisymmetry exists also at lower aspect ratio

Aspect ratio 3.3

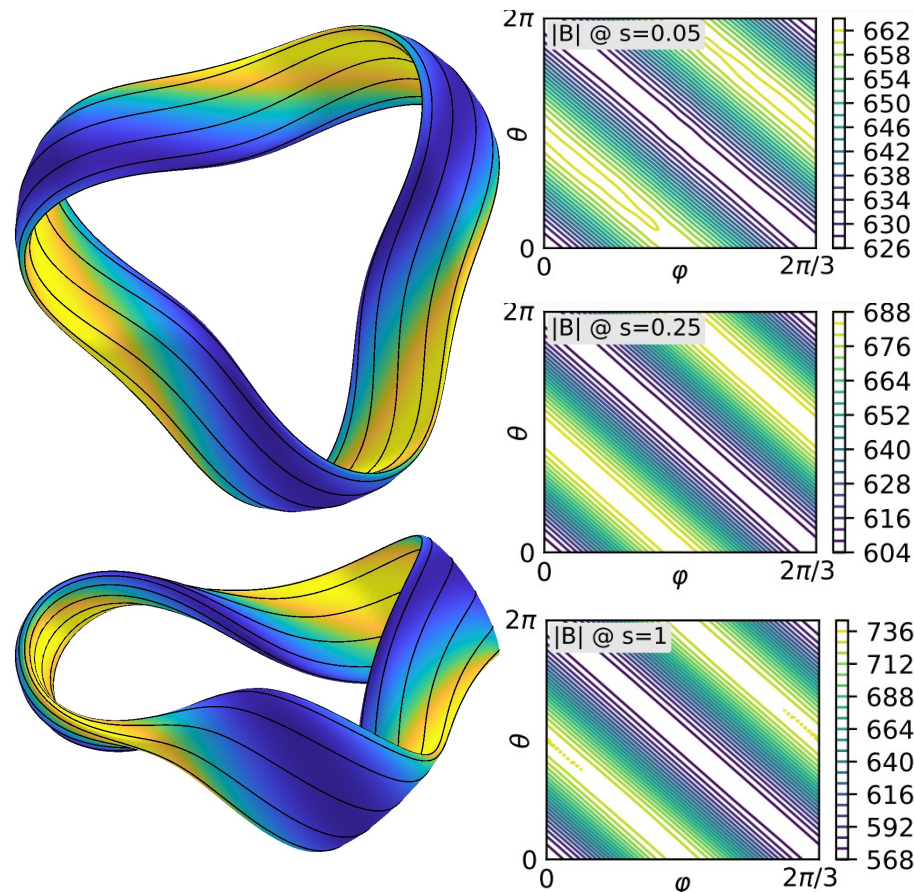


Nearly as good quasisymmetry exists also at lower aspect ratio or different # of field periods

Aspect ratio 3.3

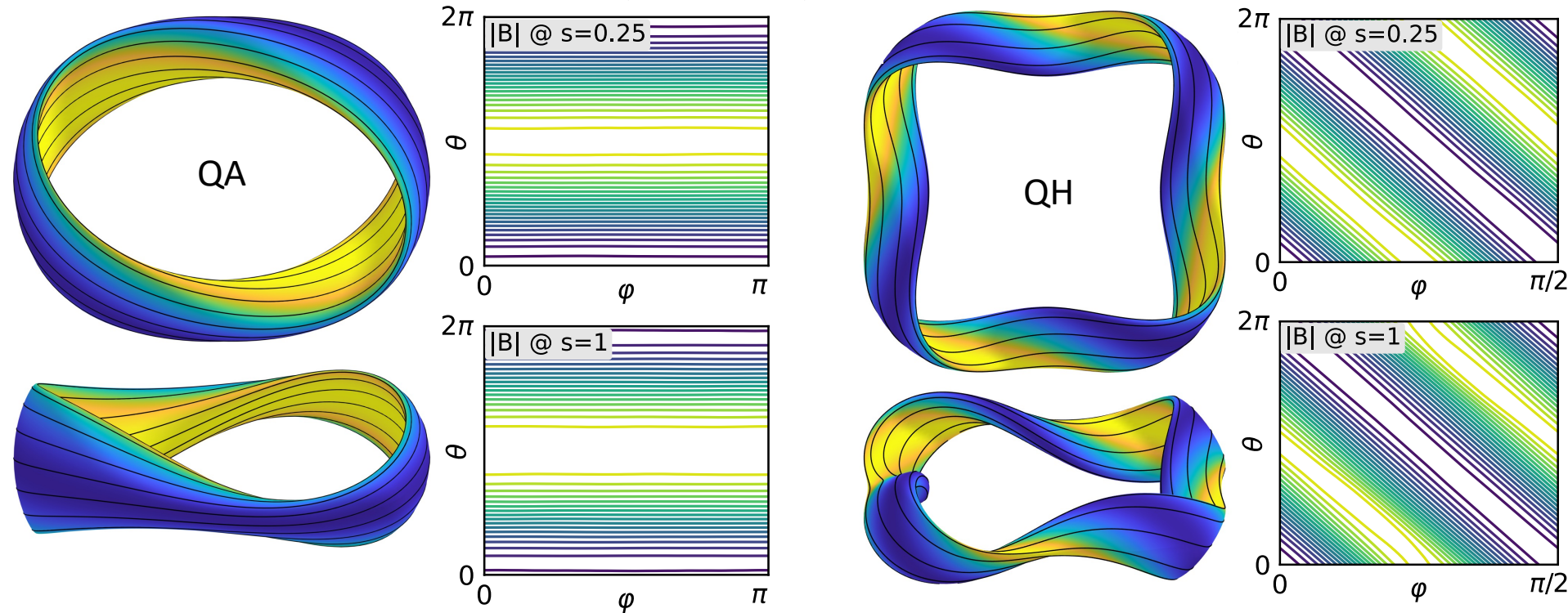


3 field periods, aspect ratio 6

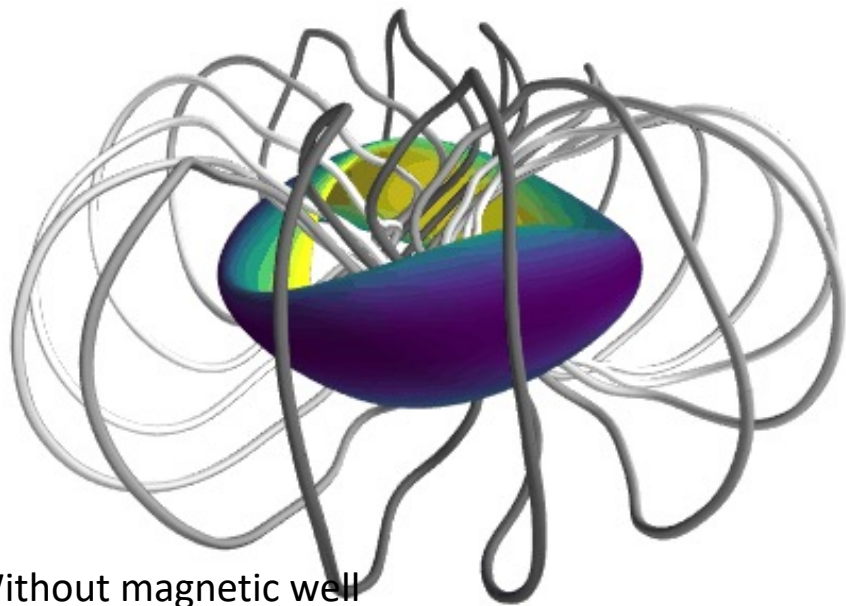


Good symmetry also exists with magnetic well

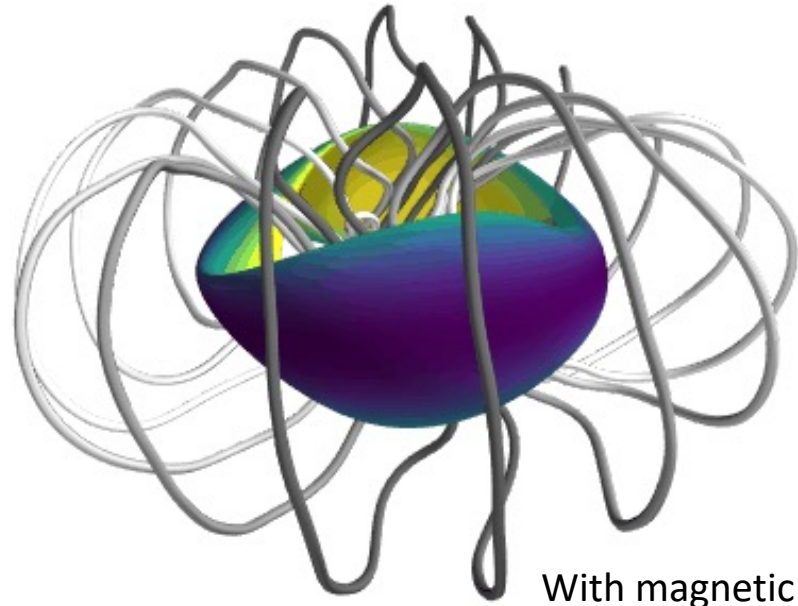
$$\frac{d^2(\text{flux surface volume})}{d(\text{toroidal flux})^2} < 0 \text{ everywhere}$$



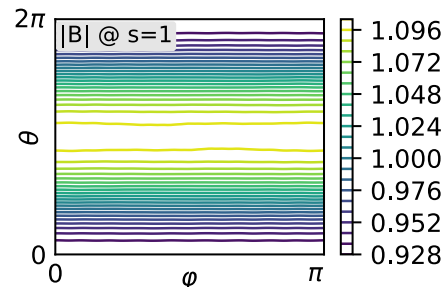
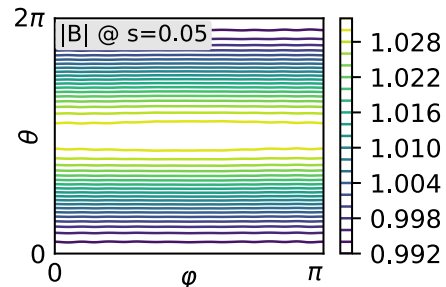
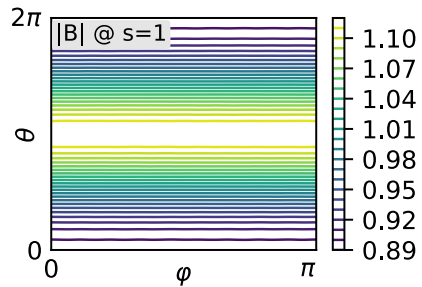
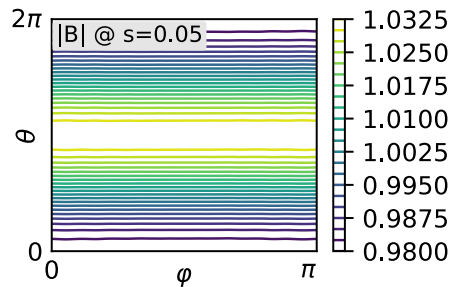
16-coil solutions have been found for the quasi-axisymmetric configurations



Without magnetic well



With magnetic well

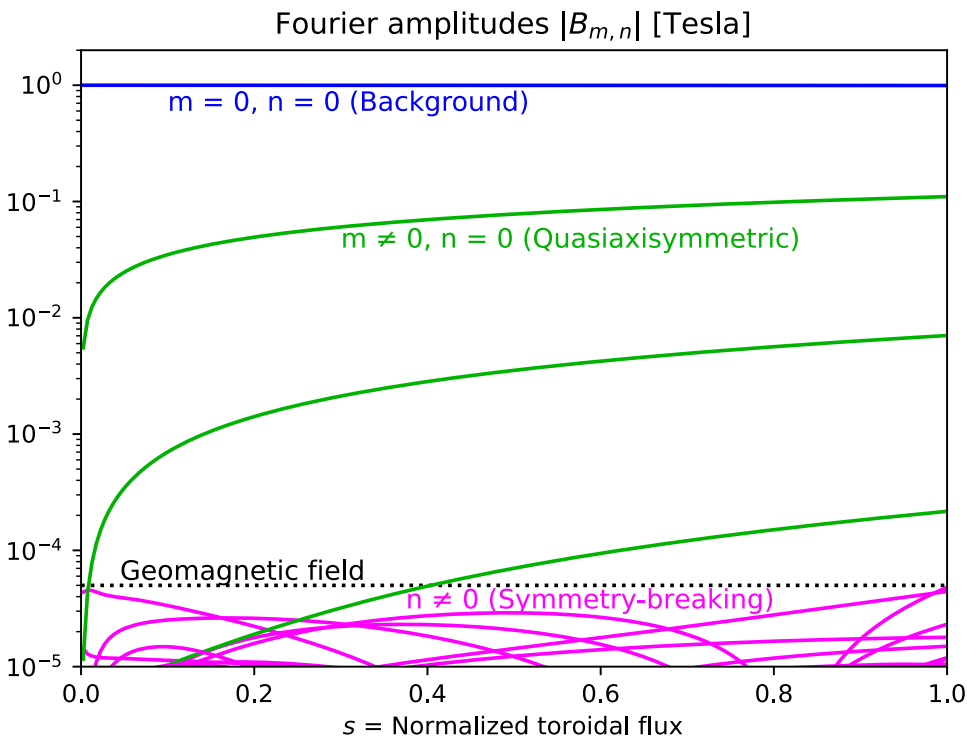


Wechsung et al, PNAS (2022).

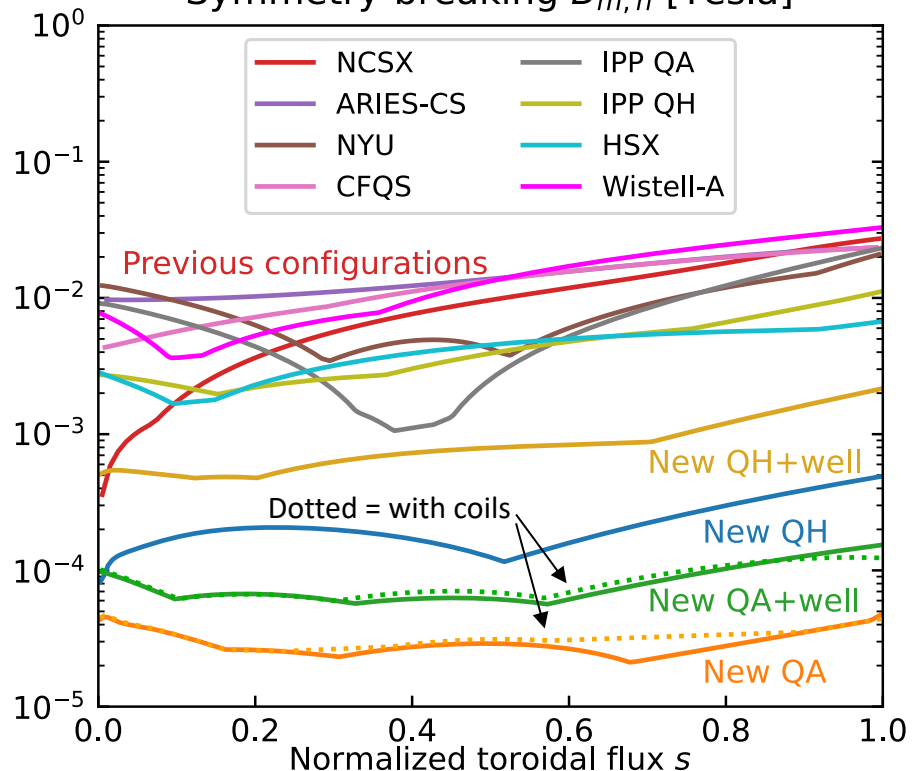
$\langle R \rangle / 10$ between filament centers.

Symmetry-breaking modes can be made extremely small

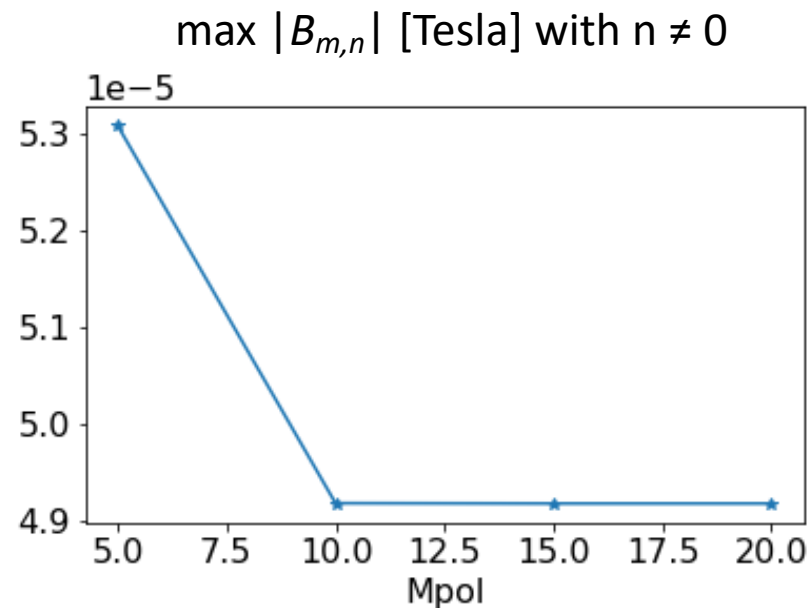
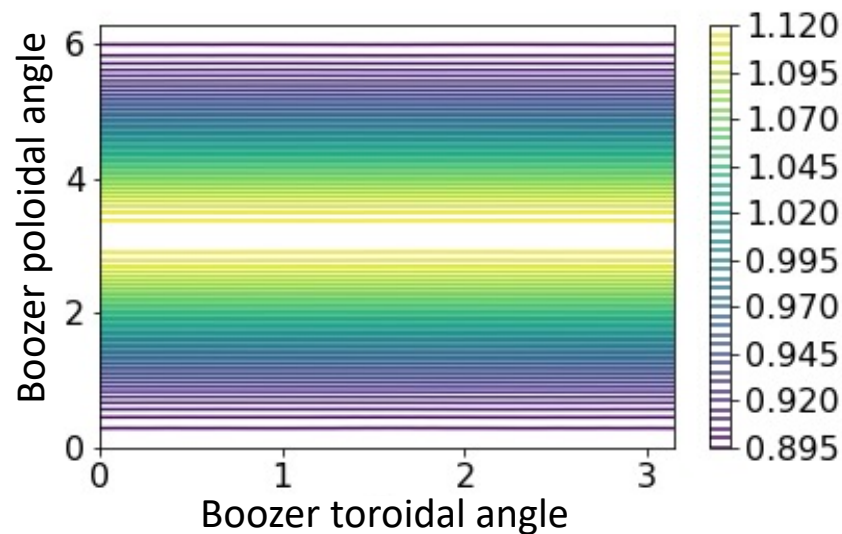
New QA configuration



Symmetry-breaking $B_{m,n}$ [Tesla]



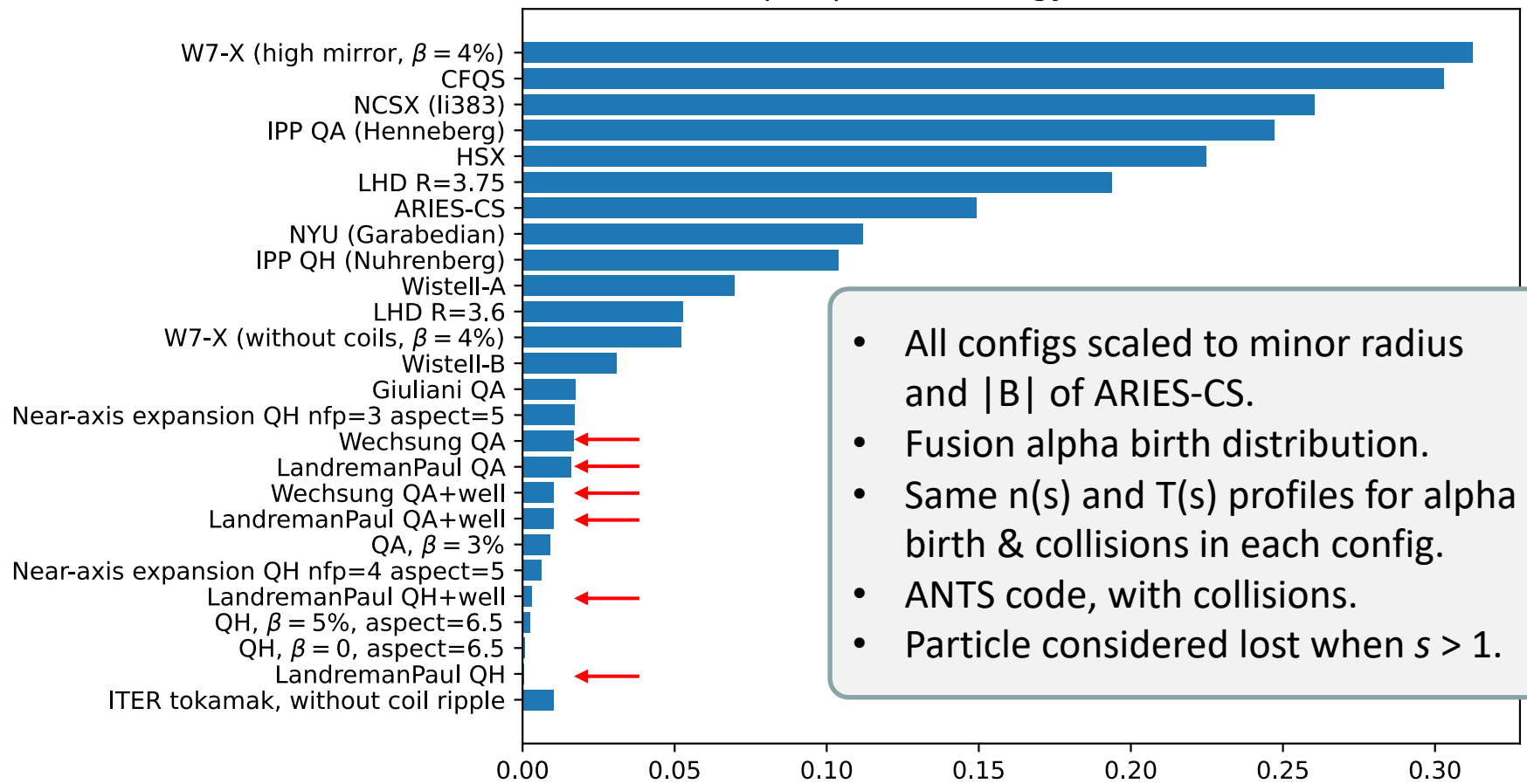
|B| in Boozer coordinates was verified by independent SPEC calculations



($N_{tor} = M_{pol}$, $L_{rad} = M_{pol} + 4$)

Quasisymmetry works: alpha particle confinement is significantly improved

Fraction of alpha particle energy lost before thermalization

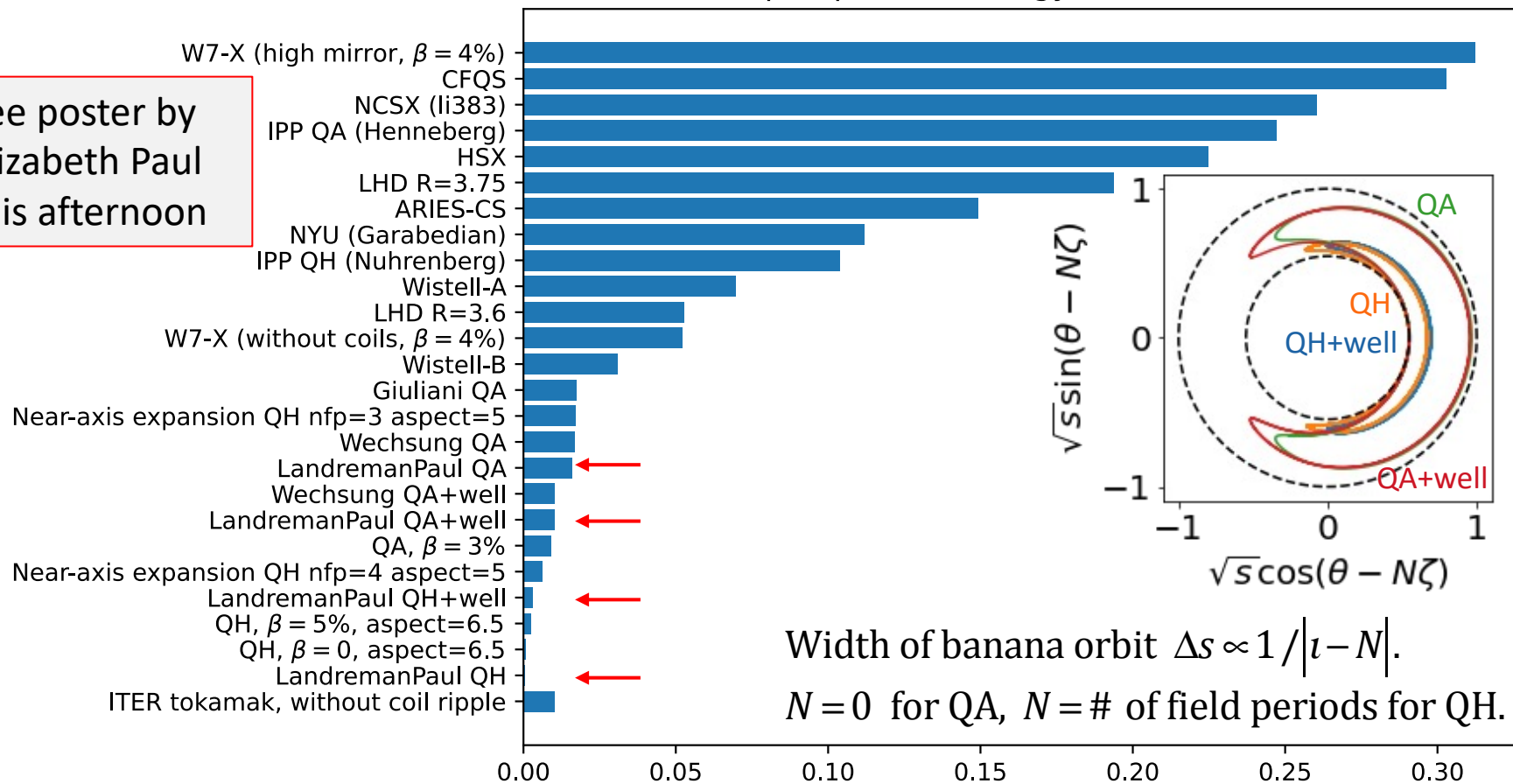


- All configs scaled to minor radius and $|B|$ of ARIES-CS.
- Fusion alpha birth distribution.
- Same $n(s)$ and $T(s)$ profiles for alpha birth & collisions in each config.
- ANTS code, with collisions.
- Particle considered lost when $s > 1$.

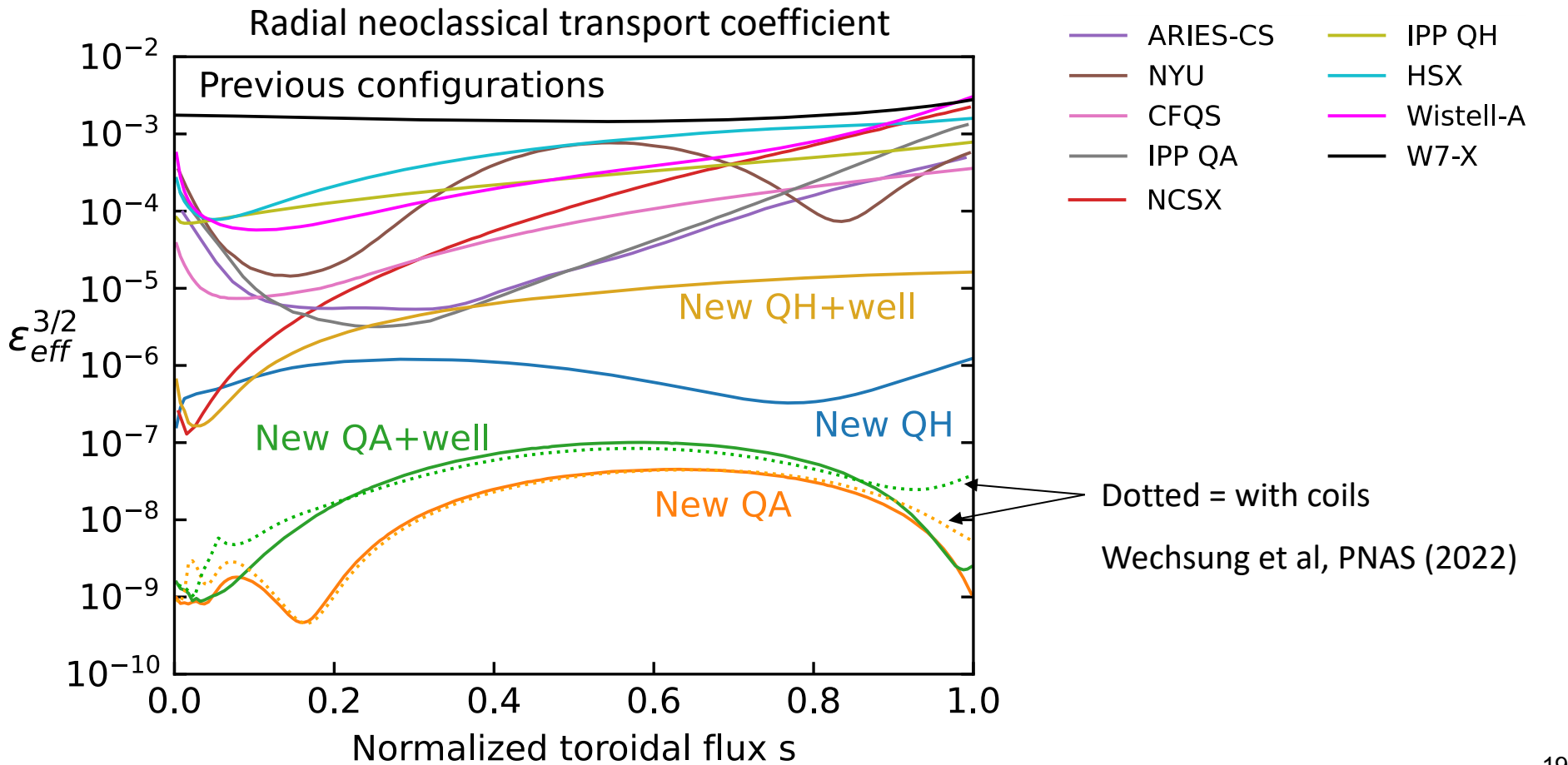
Alpha confinement in quasi-helical stellarators can be better than in a tokamak due to thinner bananas

See poster by
Elizabeth Paul
this afternoon

Fraction of alpha particle energy lost before thermalization



The symmetry also yields extremely low collisional transport for a thermal plasma



Even better quasisymmetry and ϵ_{eff} is achieved by refinement with combined plasma-and-coil optimization

Giuliani et al, arXiv:2203.03753 (2022)

Objective:

$$f = \sum_{\text{surfaces}} (\text{QS error})^2 + (\bar{l} - \bar{l}_*)^2 + \left(\begin{array}{l} \text{coil length, mean curvature,} \\ \text{max curvature, coil-coil distance terms} \end{array} \right)$$

Parameter space: coil shapes

Derivatives: analytic + adjoint method

QA

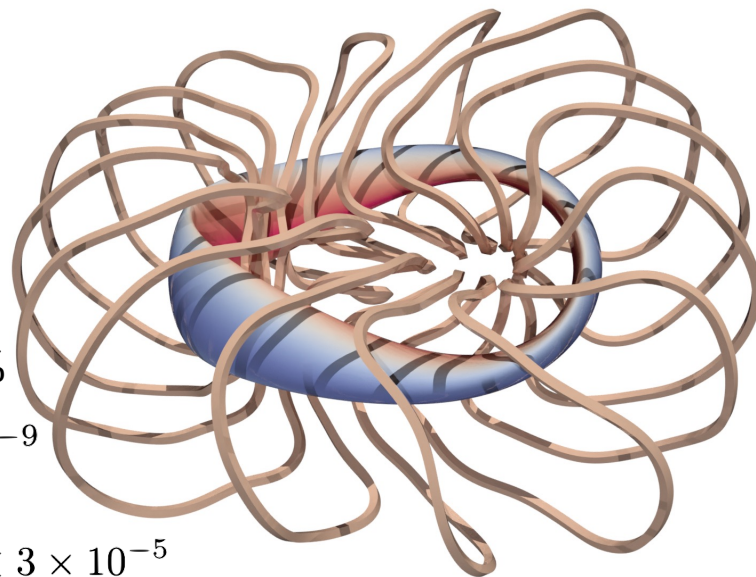
16 coils

$\iota > 0.4$

α losses < 2%

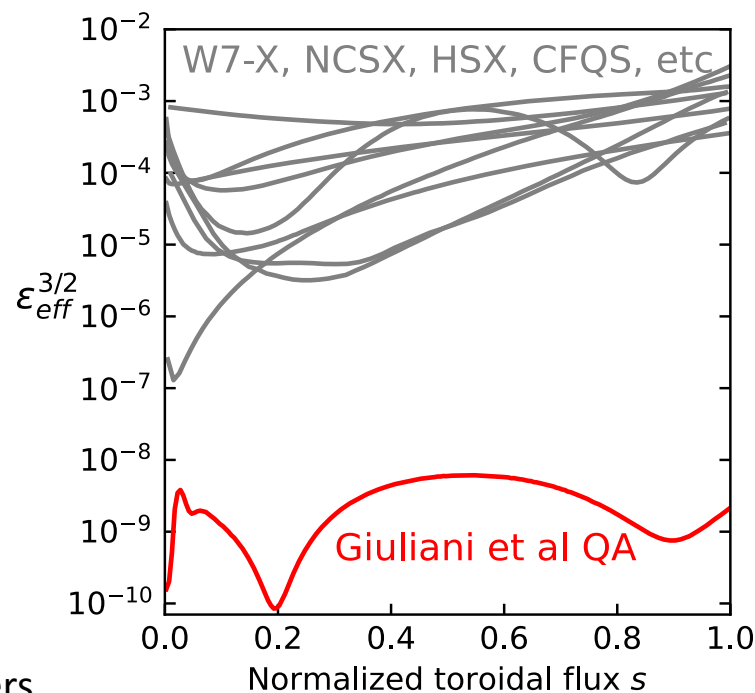
$$\epsilon_{\text{eff}}^{3/2} < 5 \times 10^{-9}$$

$$\left| \frac{B_{m,n}^{\text{non-QS}}}{B_0} \right| < 3 \times 10^{-5}$$



$\langle R \rangle / 10$ between filament centers

Collisional heat flux



- Minimal optimization recipe (low β)
- Self-consistent bootstrap current at high β
- Future directions

How can bootstrap current be included self-consistently in stellarator optimization?

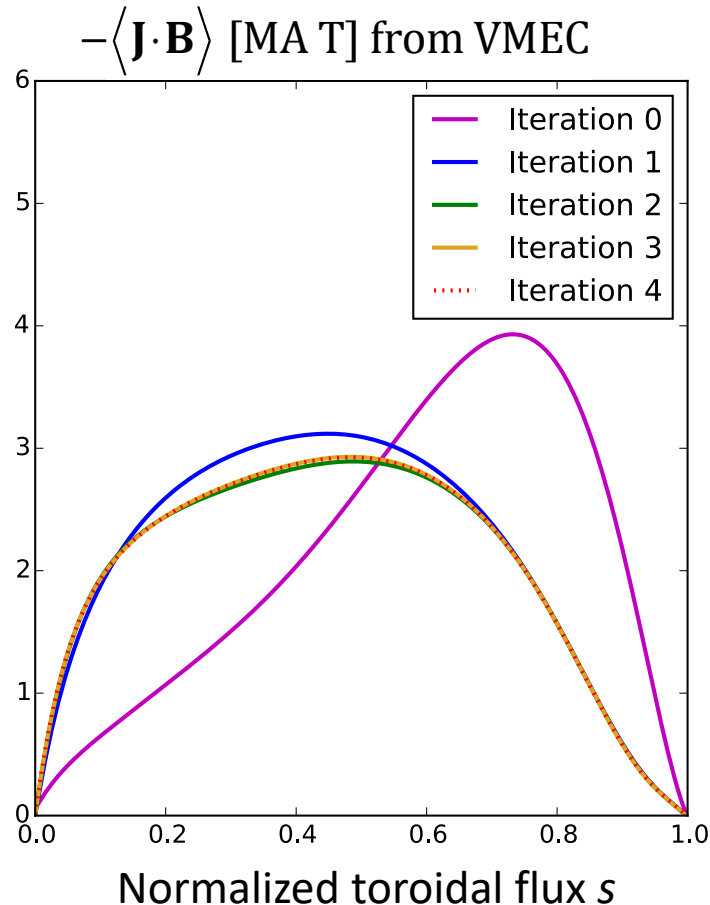
- Need *self-consistency* between MHD equilibrium and drift-kinetic equation.
- Previous method: fixed-point iteration, only after an optimization.

MHD
equilibrium
code

Drift-kinetic
code

→ VMEC: given $I_0(s)$, determine \mathbf{B}_0 .
→ SFINCS: given \mathbf{B}_0 , determine $I_1(s)$.
VMEC: given $I_1(s)$, determine \mathbf{B}_1 .
SFINCS: given \mathbf{B}_1 , determine $I_2(s)$.
...

- Accurate drift-kinetic bootstrap calculations in stellarators are computationally expensive. Preferably not in the optimization loop.



New idea: exploit quasisymmetry & use analytic expressions for tokamaks

Pytte & Boozer (1981), Boozer (1983):

$$l \rightarrow l - N$$

Bootstrap current (& other quantities) in quasisymmetry are the same as in axisymmetry, up to some substitutions:

Should be accurate for the new precisely quasisymmetric configurations.

A new set of analytical formulae for the computation of the bootstrap current and the neoclassical conductivity in tokamaks

Cite as: Phys. Plasmas **28**, 022502 (2021); doi: [10.1063/5.0012664](https://doi.org/10.1063/5.0012664)

Submitted: 6 May 2020 · Accepted: 11 December 2020 ·

Published Online: 2 February 2021



[View Online](#)

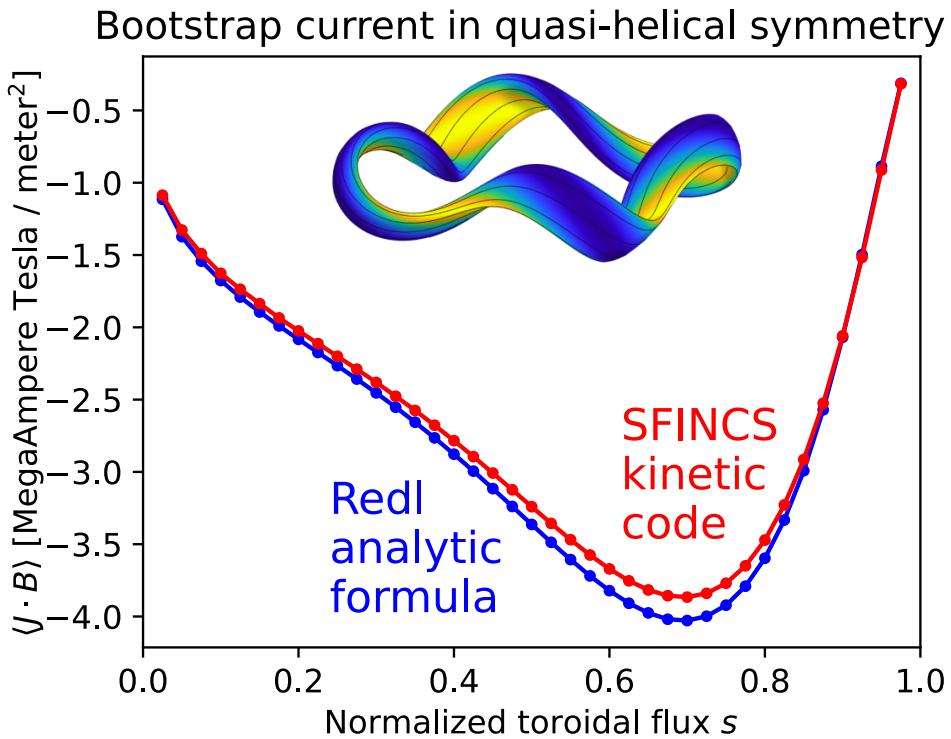
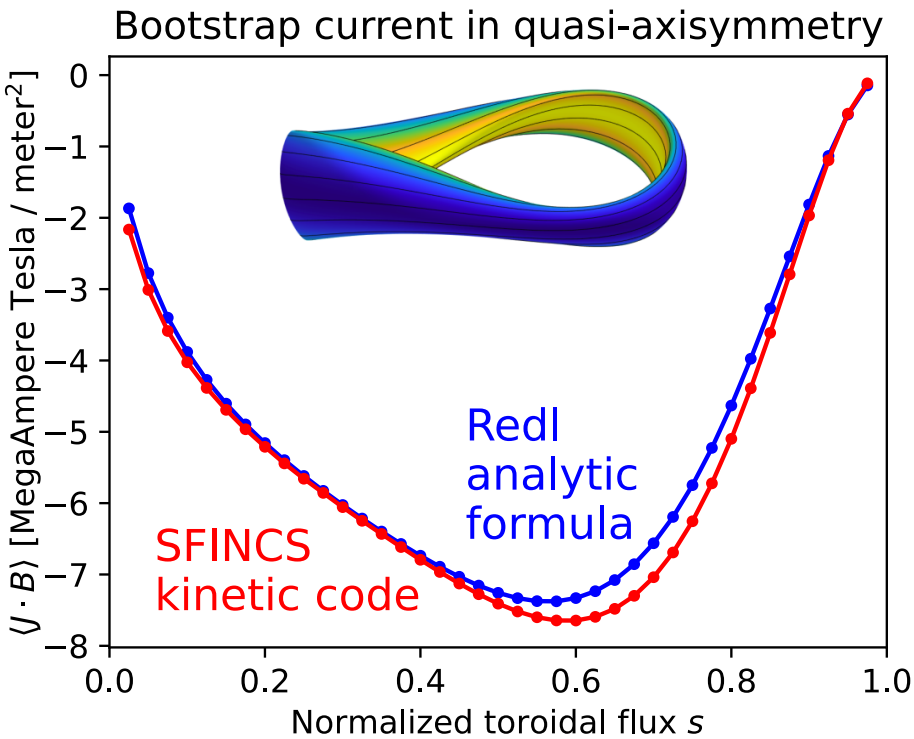


[Export Citation](#)

A. Redl,^{1,2,a)}  C. Angioni,¹  E. Belli,³  O. Sauter,⁴  ASDEX Upgrade Team^{b)} and EUROfusion MSTI Team^{c)}

Before doing new optimizations: Redl formula is accurate in previous QA & QH stellarators

$$n_e = (1 - s^5) 4 \times 10^{20} \text{ m}^{-3}, \quad T_e = T_i = (1 - s) 12 \text{ keV}$$



(Not self-consistent yet)

Optimization recipe

- Objective function: $f = f_{QS} + f_{bootstrap} + \underbrace{(A - 6.5)^2}_{\text{Boundary aspect ratio}} + \underbrace{(a - a_{\text{ARIES-CS}})^2}_{\text{Minor radius}} + (\langle B \rangle - \langle B \rangle_{\text{ARIES-CS}})^2$

$$f_{QS} = \int d^3x \left(\frac{1}{B^3} \left[(N - I) \mathbf{B} \times \nabla B \cdot \nabla \psi - (G + NI) \mathbf{B} \cdot \nabla B \right] \right)^2$$

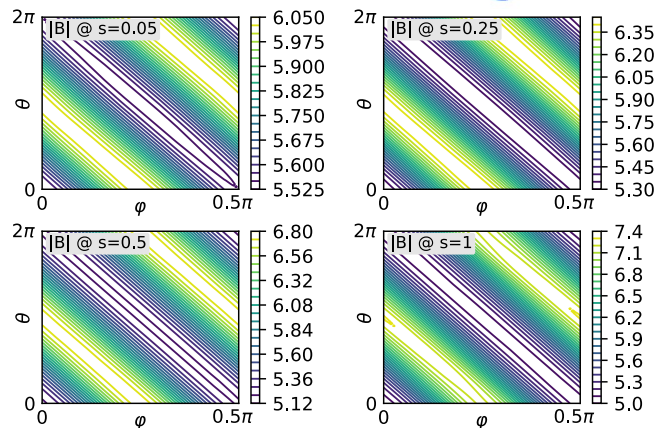
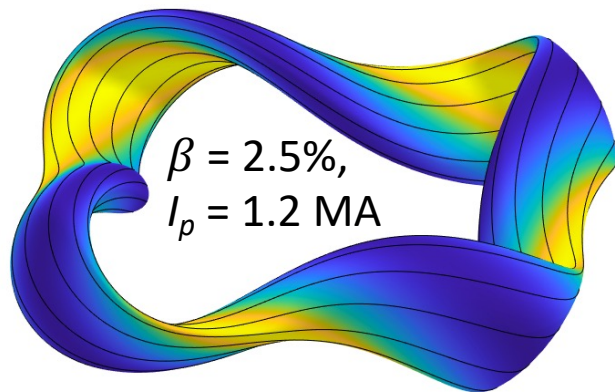
$$f_{bootstrap} = \frac{\int_0^1 ds \left[\langle \mathbf{j} \cdot \mathbf{B} \rangle_{\text{vmec}} - \langle \mathbf{j} \cdot \mathbf{B} \rangle_{\text{Redl}} \right]^2}{\int_0^1 ds \left[\langle \mathbf{j} \cdot \mathbf{B} \rangle_{\text{vmec}} + \langle \mathbf{j} \cdot \mathbf{B} \rangle_{\text{Redl}} \right]^2}$$

- Parameter space: $\{R_{m,n}, Z_{m,n}, \text{toroidal flux, current spline values}\}$
or $\{R_{m,n}, Z_{m,n}, \text{toroidal flux, iota spline values}\}$
- Cold start
- Algorithm: default for least-squares in scipy (trust region reflective)
- Steps: increasing # of modes varied: m and |n/nfp| up to j in step j

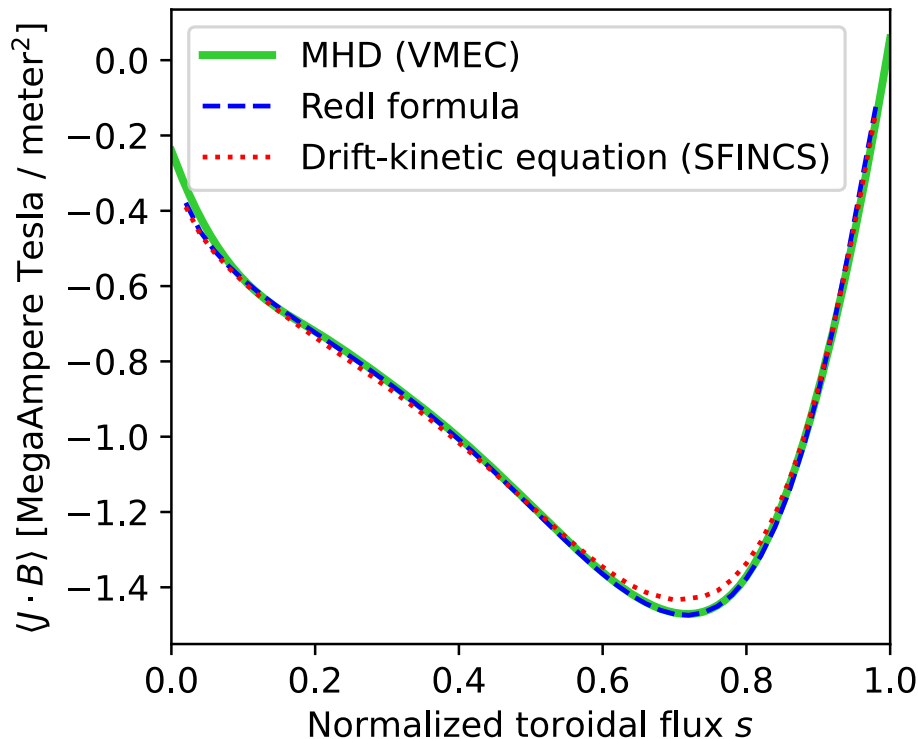
Example of optimization with self-consistent bootstrap current

$$n_{e0} = 2.2e20/\text{meters}^3$$

$$T_{e0} = T_{i0} = 10 \text{ keV}$$



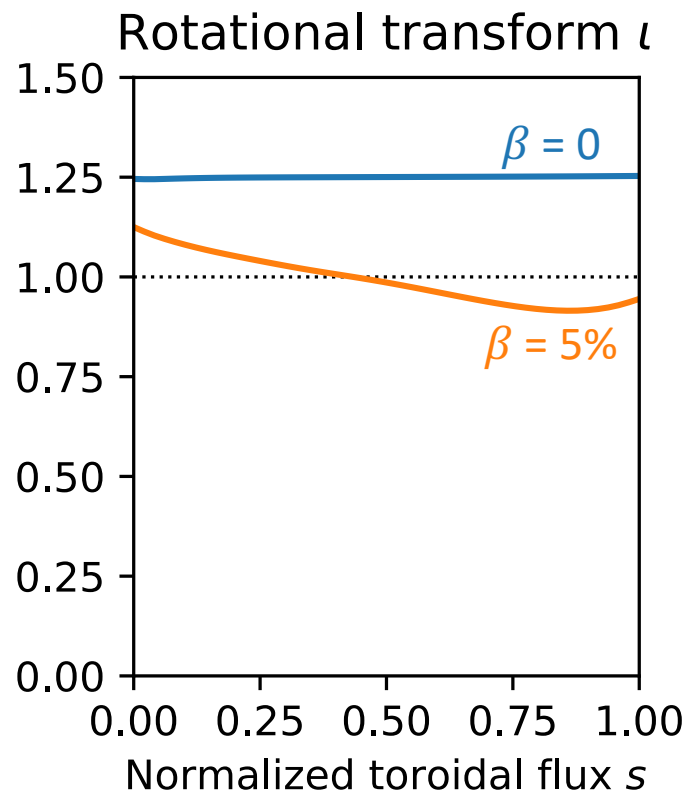
Bootstrap current profile



All input/output files and optimization scripts online at
doi.org/10.5281/zenodo.6520103

To reach reactor-relevant 5% beta in QH without crossing $\iota=1$, a constraint on ι can be included

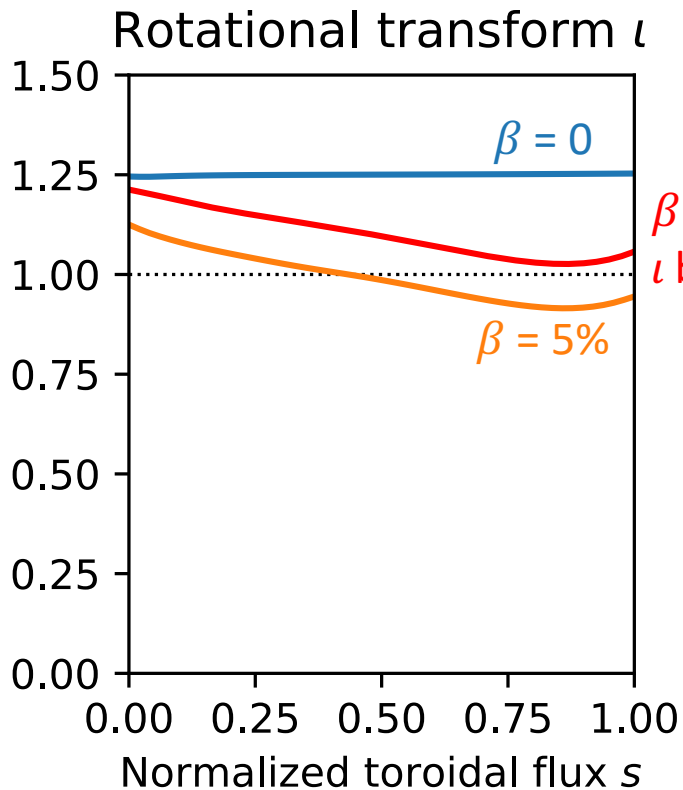
Crossing $\iota=1$, the worst resonance, is probably unacceptable.



$$n_{e0} = 3 \times 10^{20} / \text{meters}^3, T_{e0} = T_{i0} = 15 \text{ keV}$$

To reach reactor-relevant 5% beta in QH without crossing $\iota=1$, a constraint on ι can be included

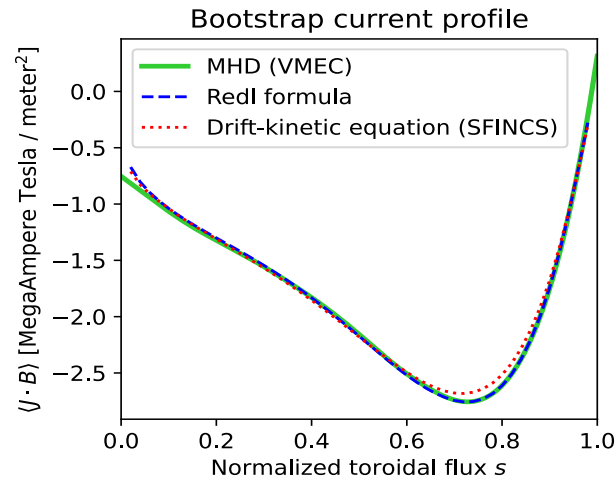
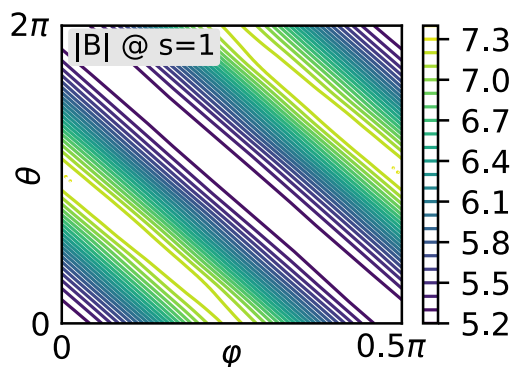
Crossing $\iota=1$, the worst resonance, is probably unacceptable.



Solution: Add barrier term in objective

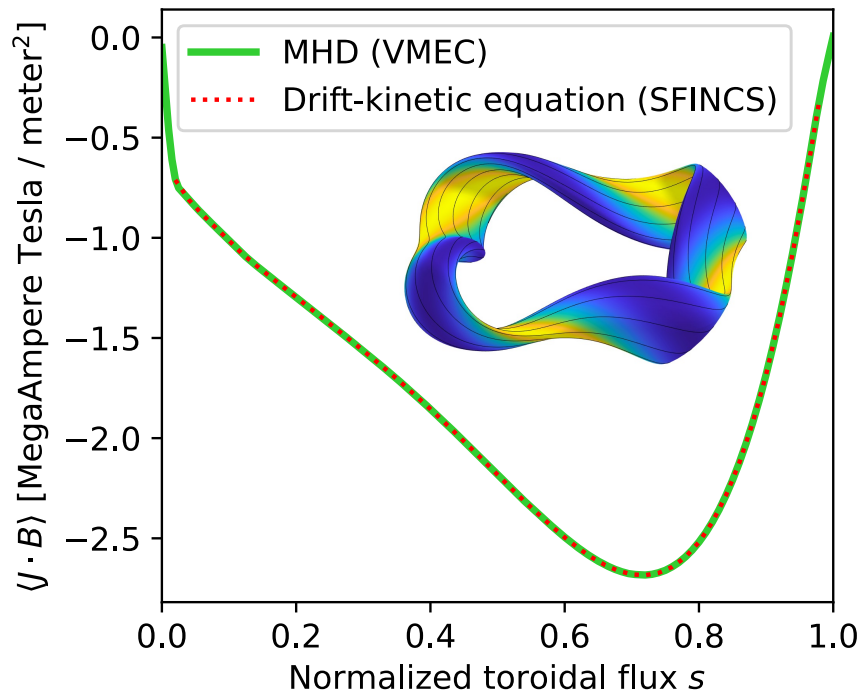
$$f += \int_0^1 ds \left[\min(|\iota(s)| - 1.03, 0) \right]^2$$

Quasisymmetry & bootstrap consistency remain good:



If you want *perfectly* self-consistent current,
you can do a few fixed-point iterations at the end

Bootstrap current profile

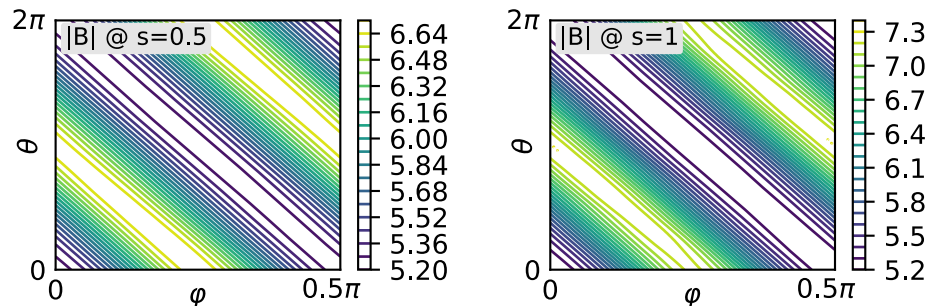


$$\langle \beta \rangle = 5\%, \quad \varepsilon_{eff}^{3/2} < 6 \times 10^{-5}$$

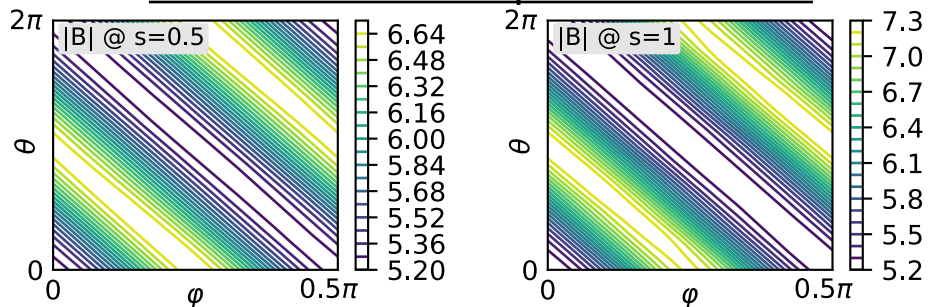
α -particle losses < 0.3%

No significant degradation in quasisymmetry:

Optimization with Redl current

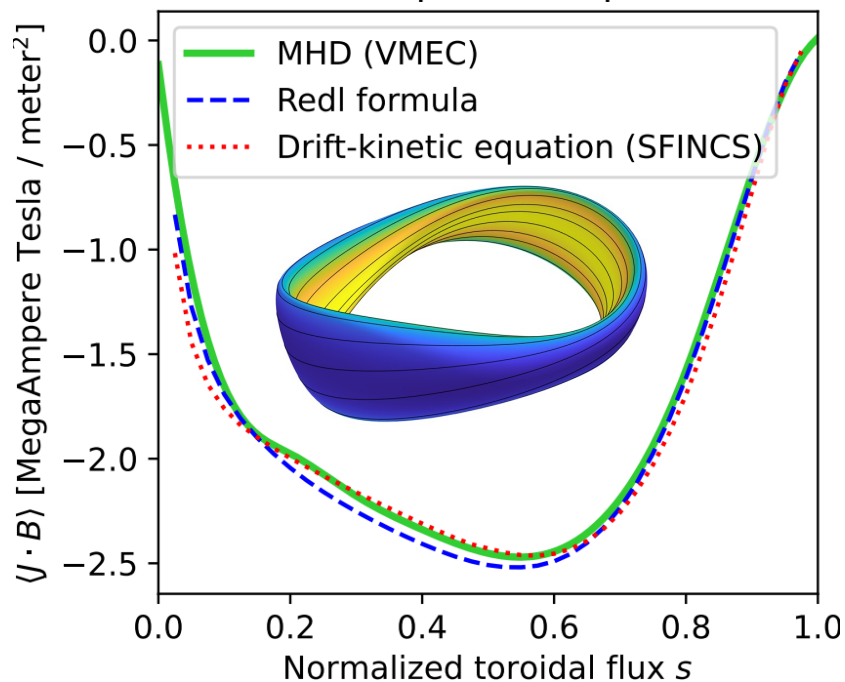


After SFINCS fixed-point iterations



The optimization with self-consistent bootstrap current also works for quasi-axisymmetry

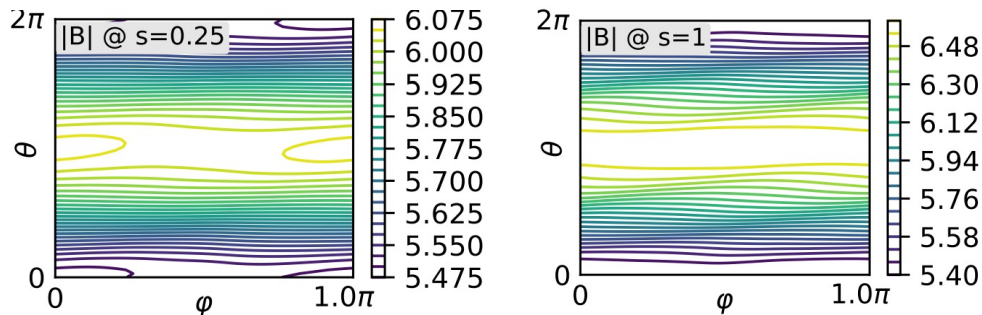
Bootstrap current profile



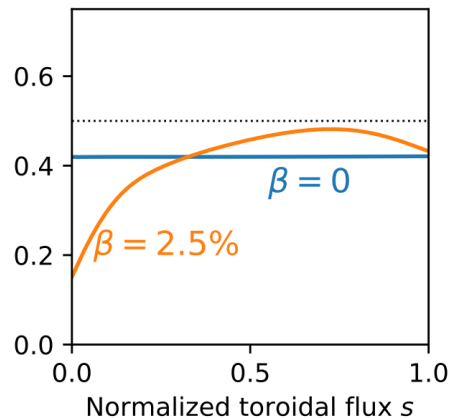
$$\langle \beta \rangle = 2.5\%, \quad \epsilon_{eff}^{3/2} < 2 \times 10^{-5}$$

α -particle losses < 1.5%

Symmetry is not as good as for vacuum, but sufficient for excellent confinement

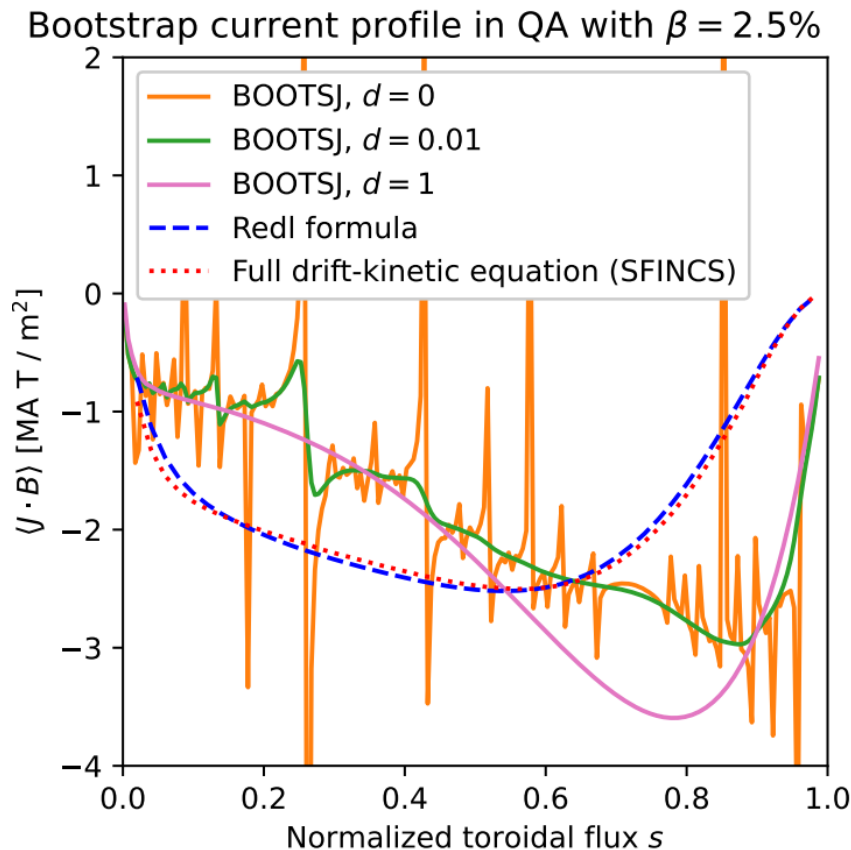


Rotational transform ι



Possible islands where $\iota = 1/4, 2/5, 1/3$?

Redl formula is more accurate than long-mean-free-path stellarator bootstrap formula, & free of resonances



Stellarator bootstrap formulae for long-mean-free-path (low collisionality):

Shaing & Callen (1983),

Shaing et al (1989),

Helander, Parra & Newton (2017)

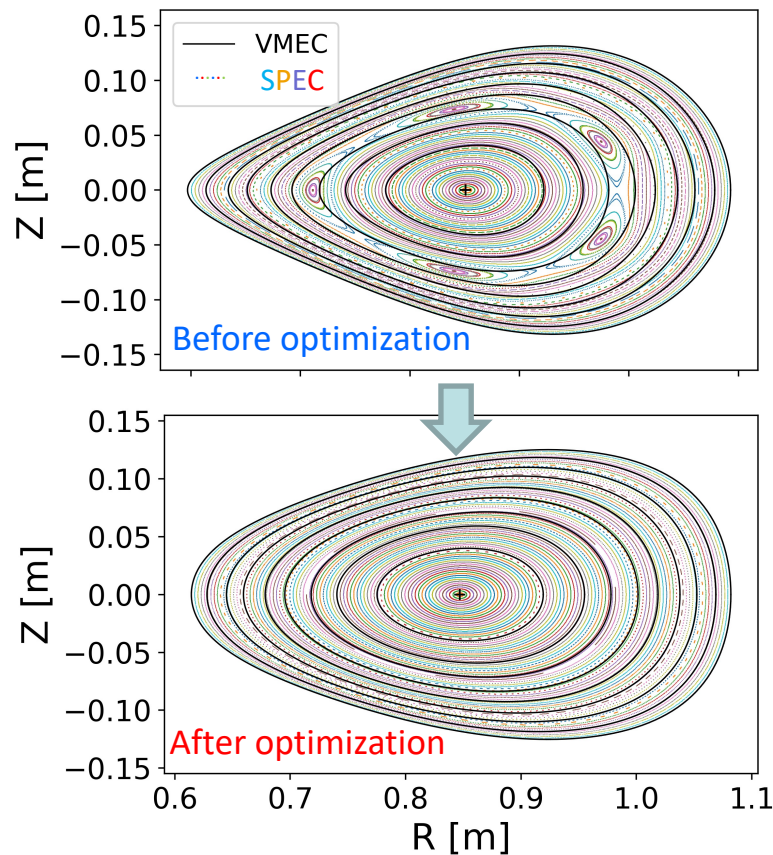
BOOTSJ ad-hoc smoothing:

$$\frac{1}{m - n/\iota} \rightarrow \frac{m - n/\iota}{(m - n/\iota)^2 + m^2 d^2}$$

- Minimal optimization recipe (low β)
- Self-consistent bootstrap current at high β
- Future directions

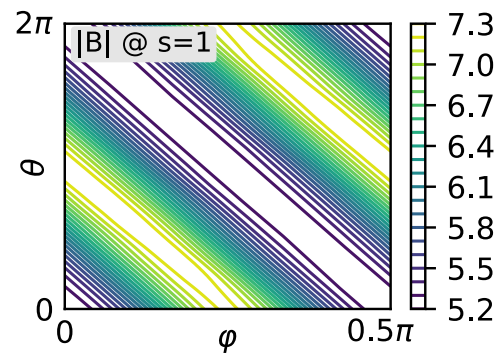
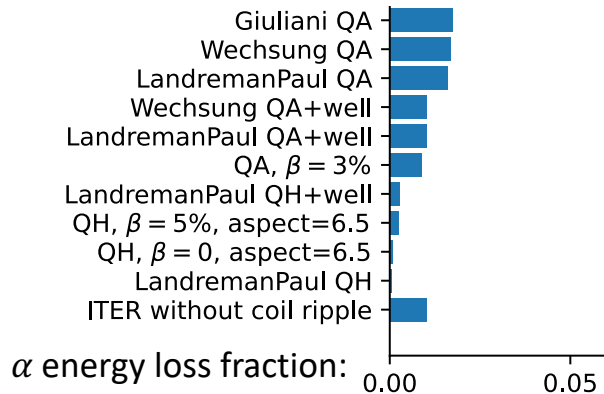
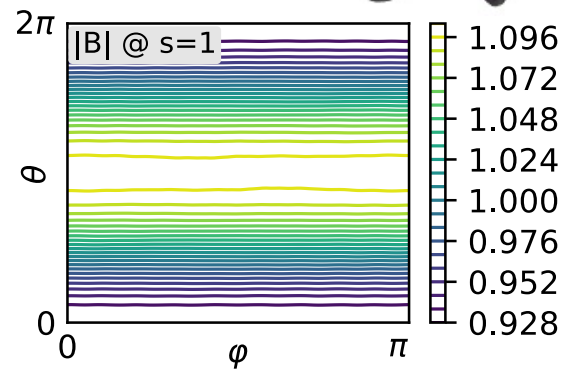
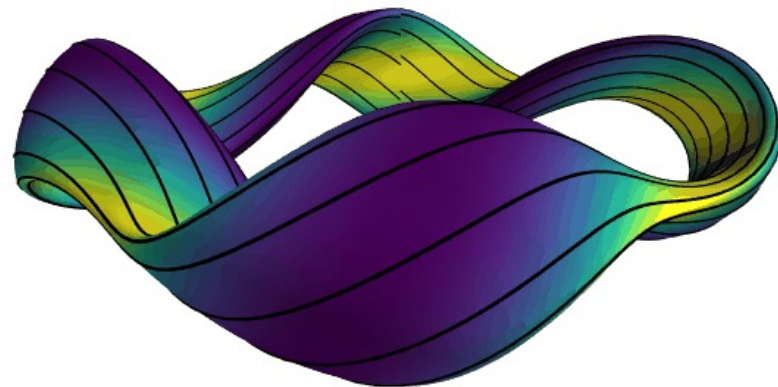
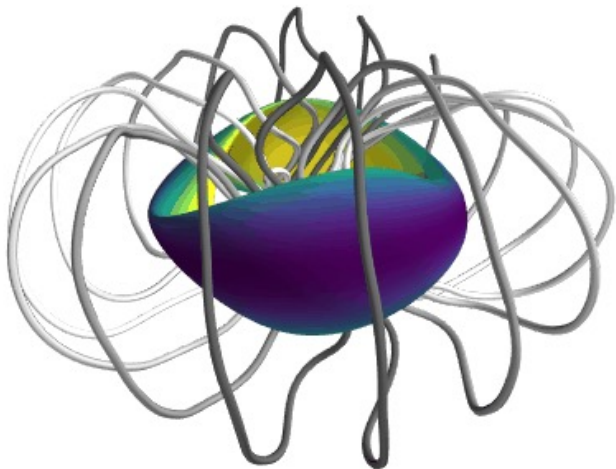
Future directions

- For the high β configurations, check surface quality, & eliminate any islands.
- Coils & MHD stability for the high β configurations.
- Check robustness to uncertainty in the pressure profile.
- Similar recipes for quasi-poloidal symmetry or quasi-isodynamic?



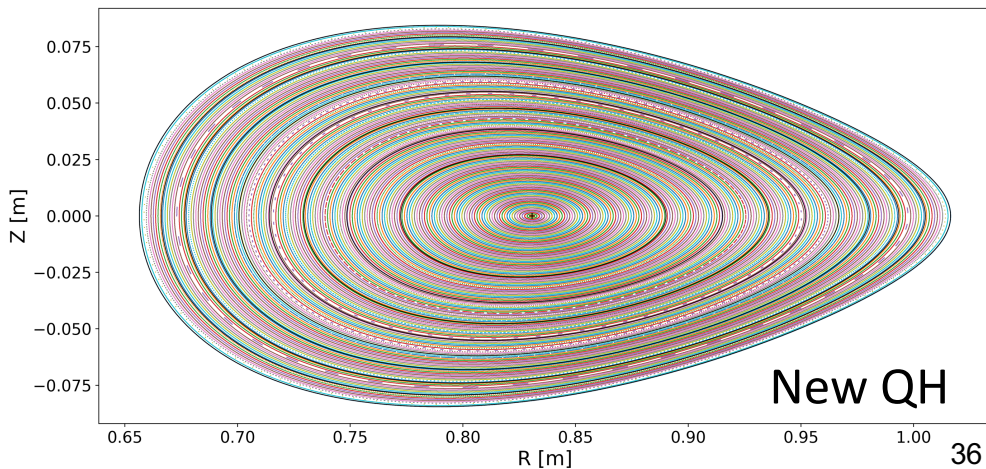
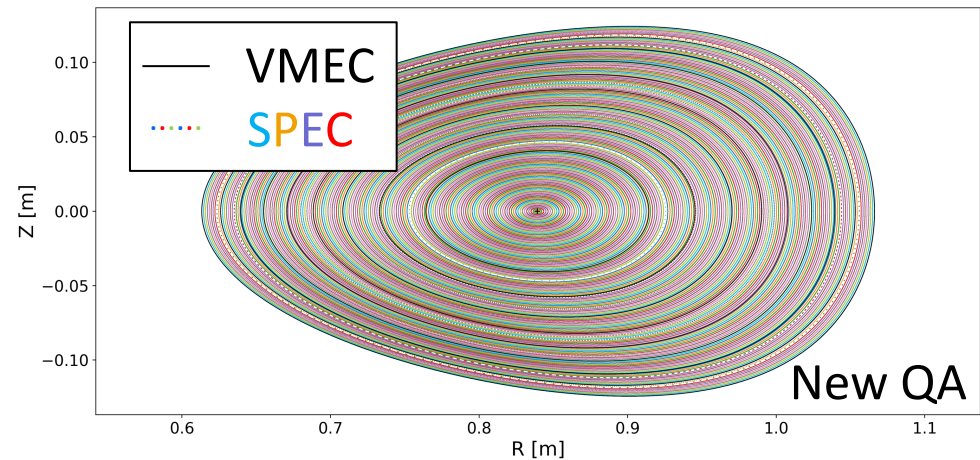
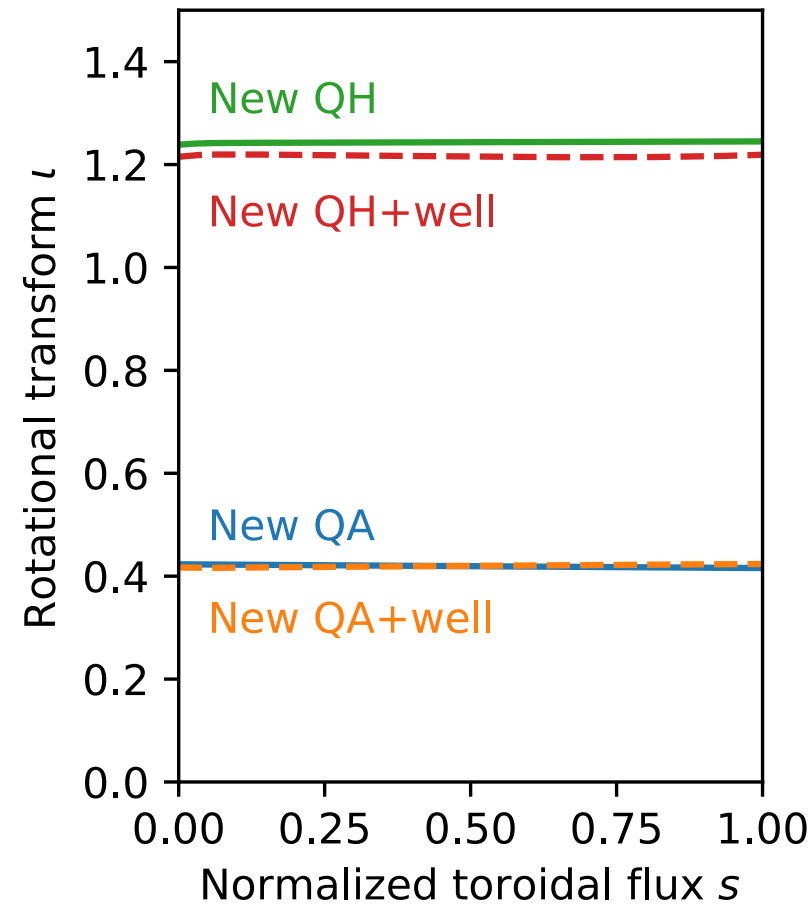
*ML, Medasani & Zhu (2021),
Baillod et al (2022)*

It is now possible to design stellarators with alpha confinement close to or better than a tokamak.



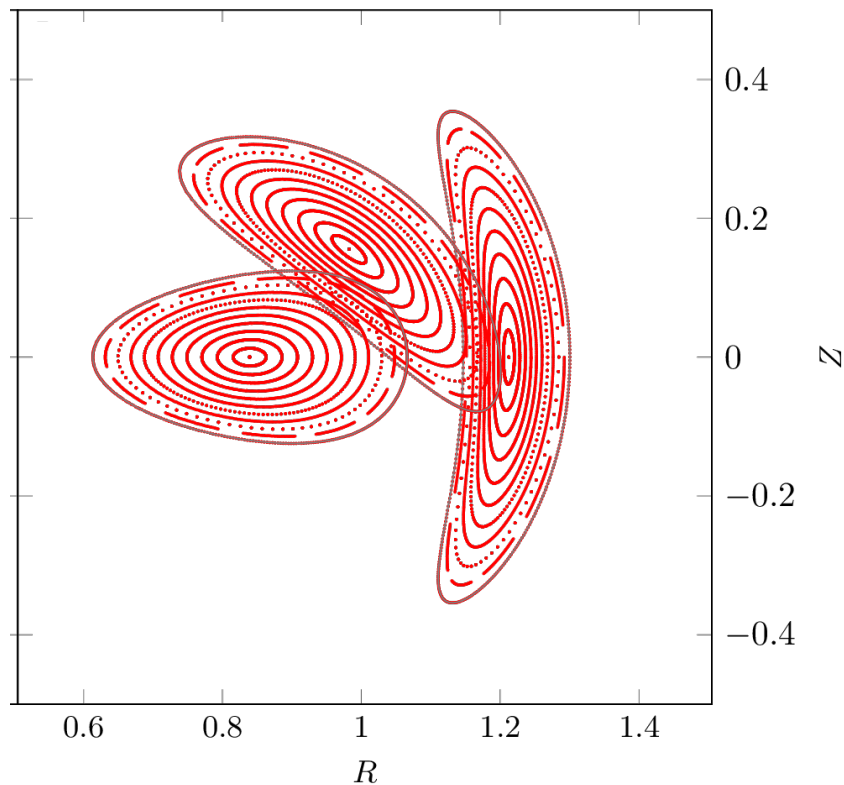
Extra slides

The new configurations have small magnetic shear

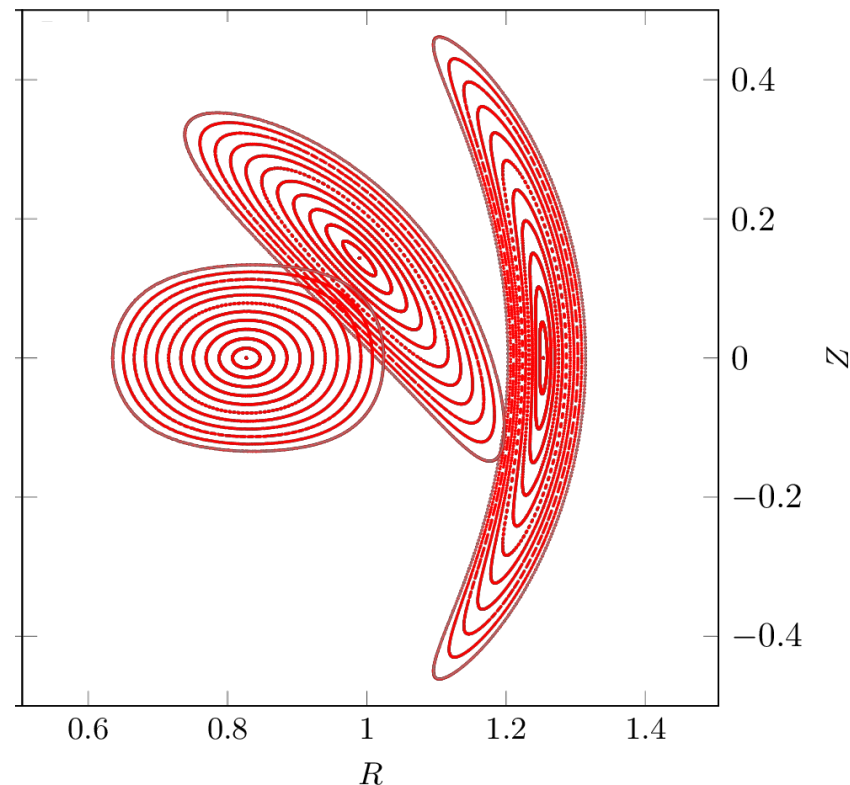


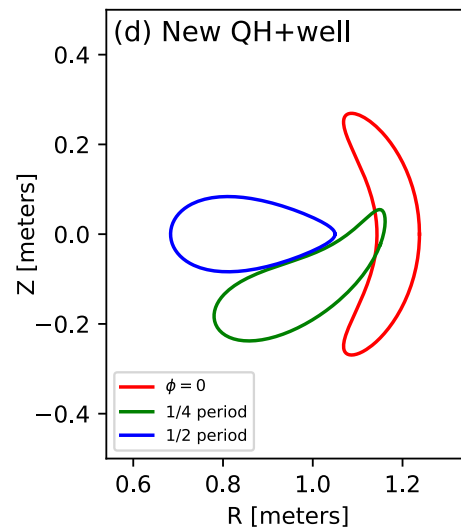
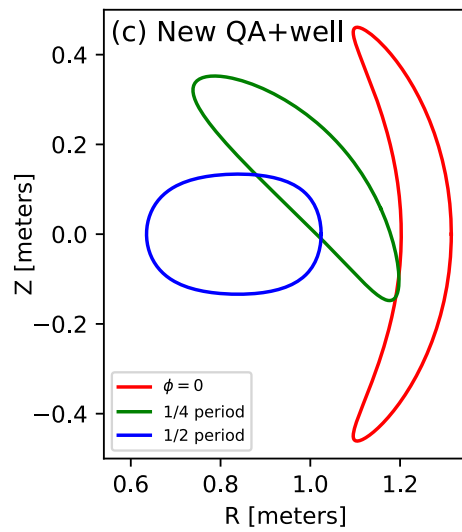
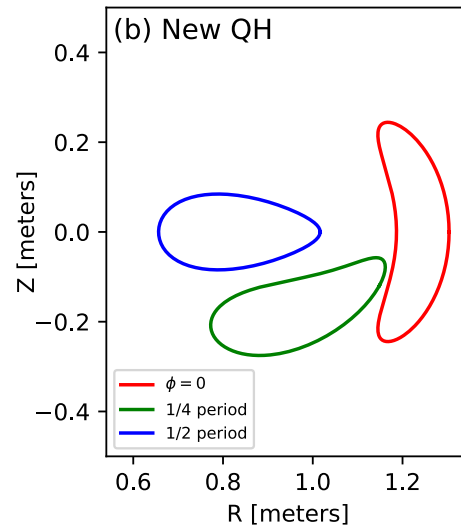
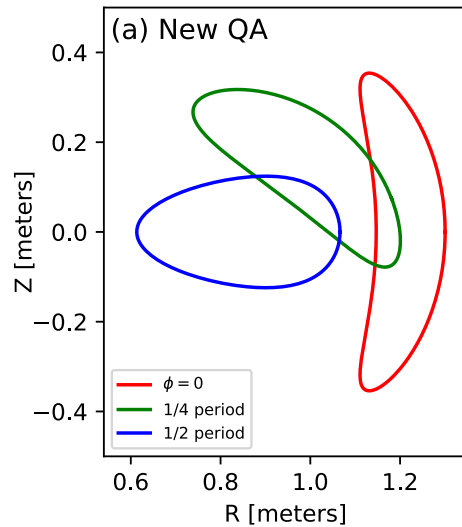
Good flux surface exist with coils

New QA

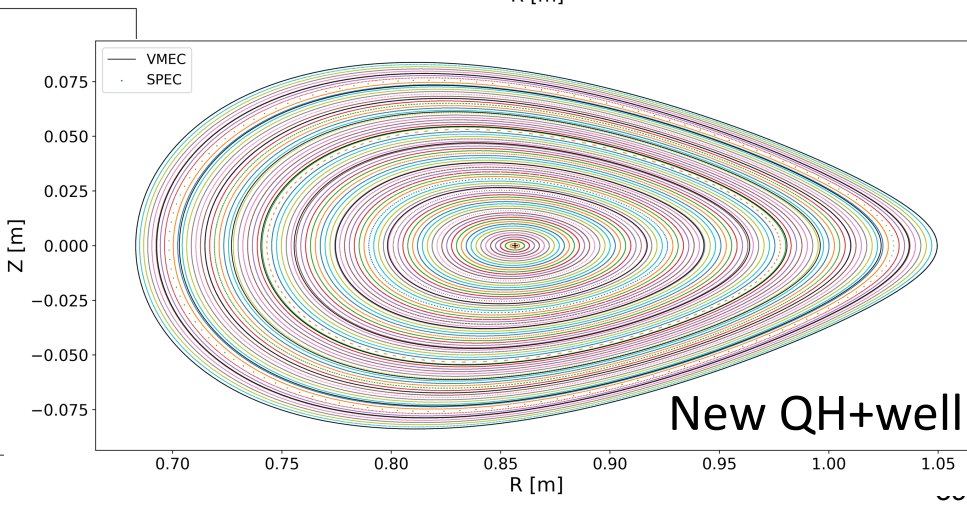
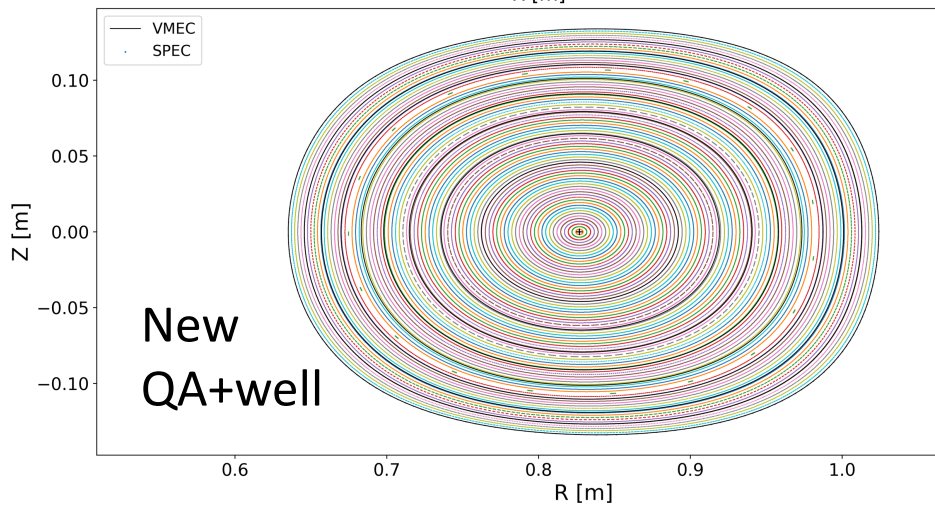
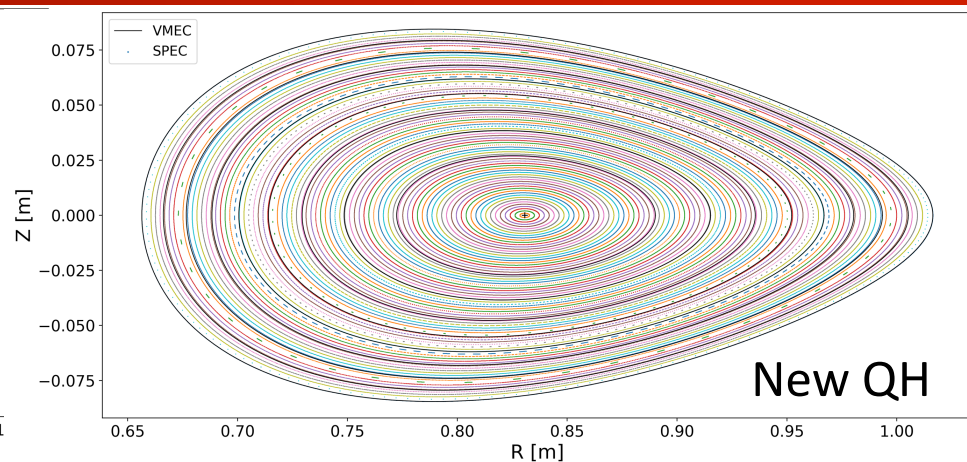
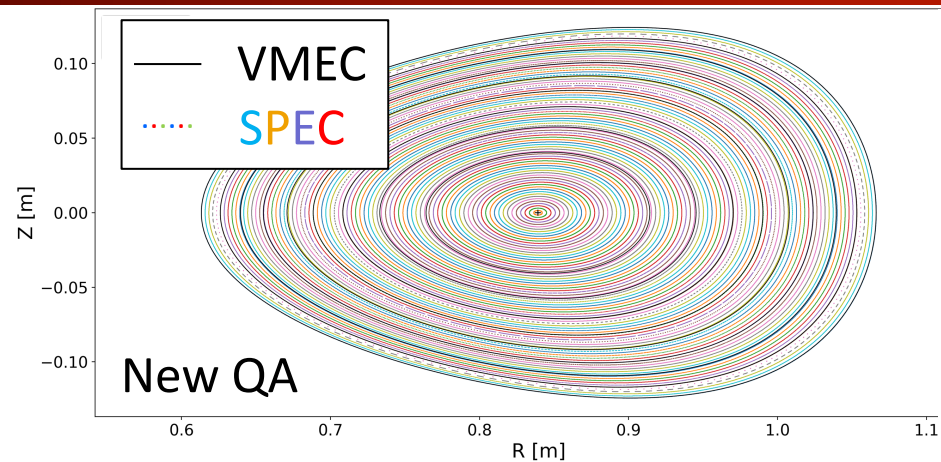


New QA+well

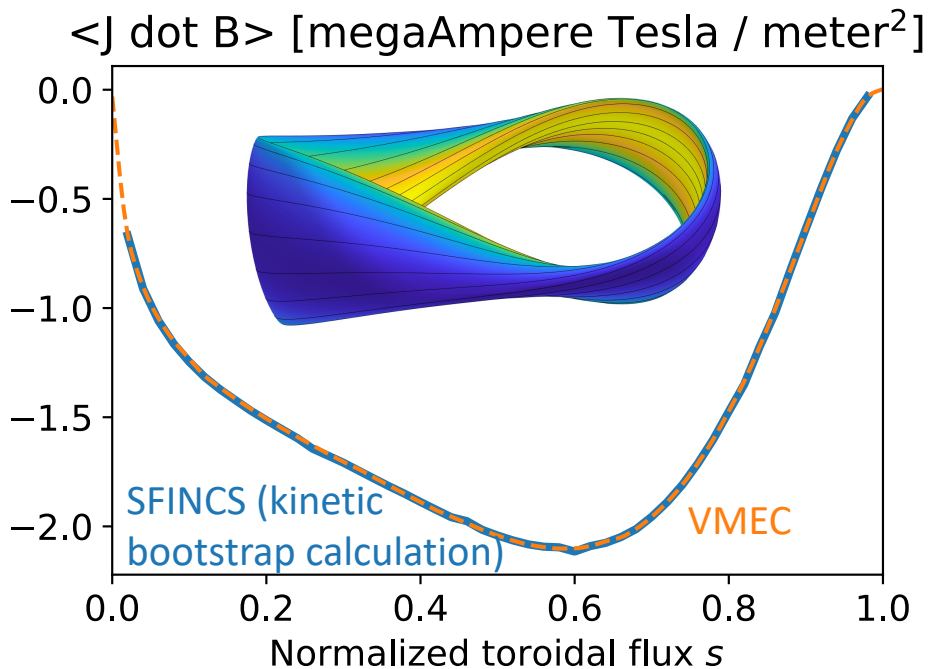




SPEC confirms the new QA/QH configurations have good surfaces



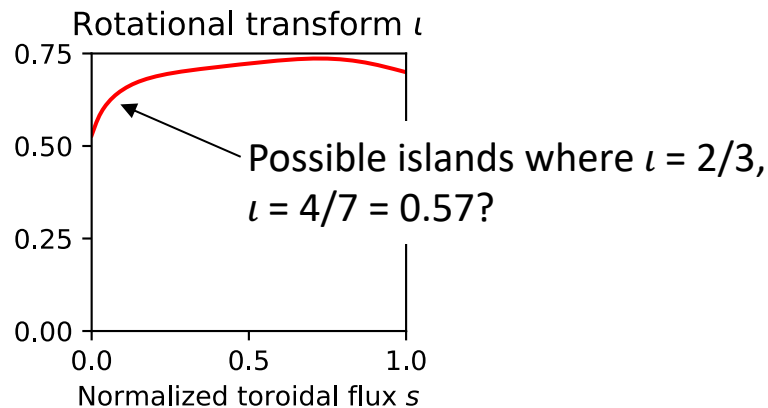
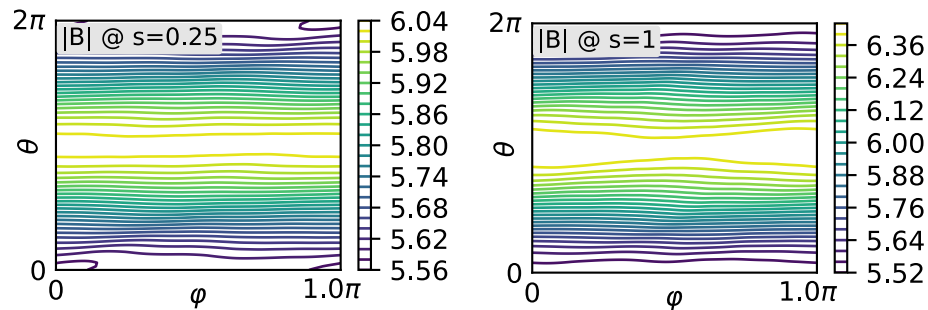
The optimization with self-consistent bootstrap current also works for quasi-axisymmetry



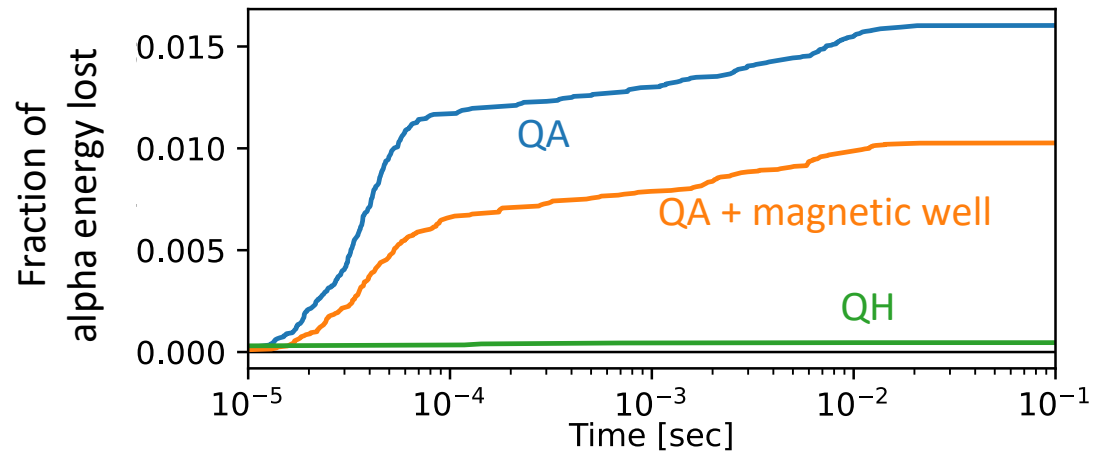
$$\langle \beta \rangle = 3\%, \quad \varepsilon_{eff}^{3/2} < 7 \times 10^{-6}$$

α -particle losses < 1%

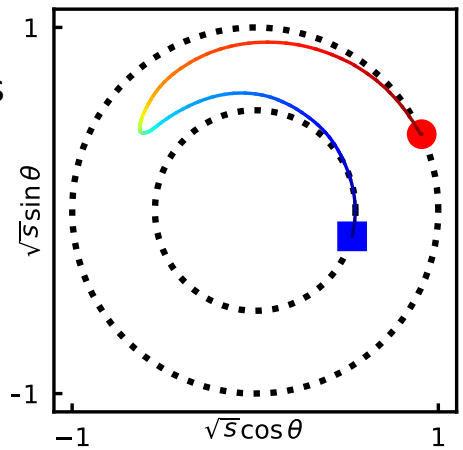
Symmetry is not as good as for vacuum, but sufficient for excellent confinement



Why the configurations with best quasisymmetry not have the best trajectory confinement?



Lost trajectories in the new QA look like this:

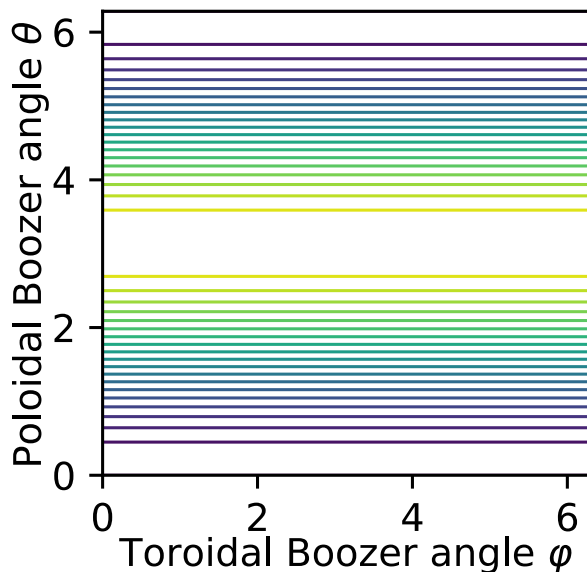


Width of banana orbit $\Delta s \approx \left| \frac{mvR\sqrt{2r\bar{\eta}}}{(\iota - N)\psi_{edge}Ze} \right| \propto \frac{1}{|\iota - N|}$

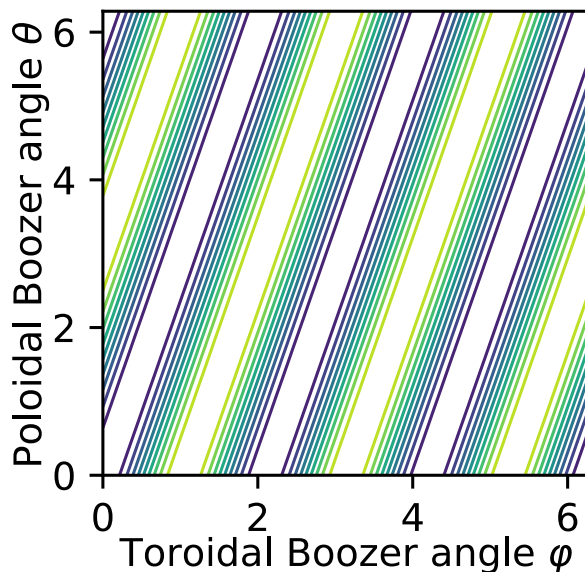
For fixed minor radius, $\frac{\Delta s_{QA}}{\Delta s_{QH}} \sim 4$

2 types of quasisymmetry

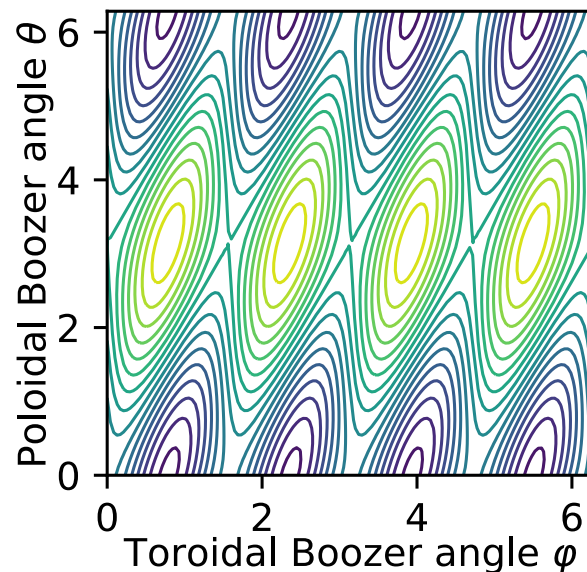
Quasi-axisymmetry
(QA): $B = B(r, \theta)$



Quasi-helical symmetry
(QH): $B = B(r, \theta - N\phi)$

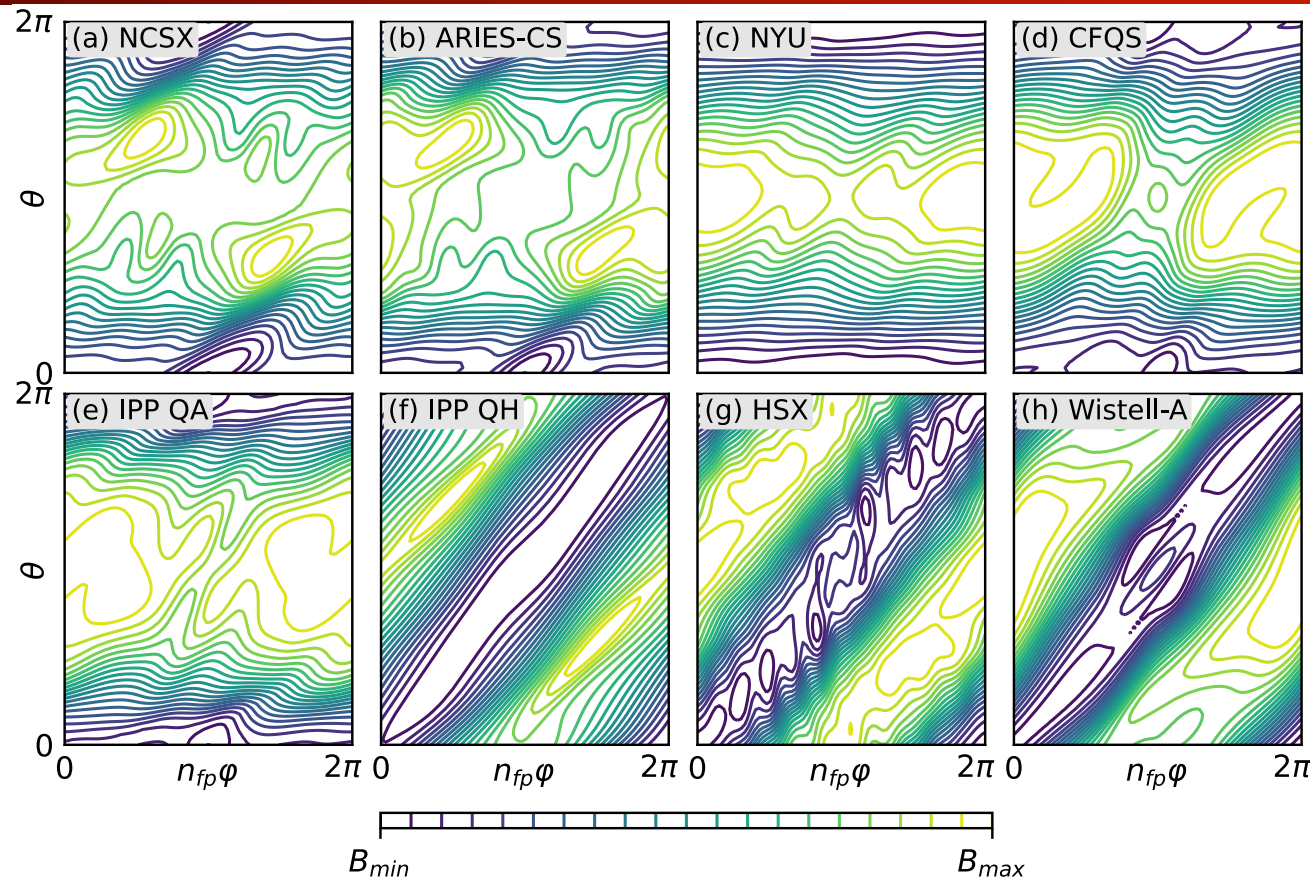


General stellarator
(not symmetric)



Contours of $B = |\mathbf{B}|$: B_{min}  B_{max}

Previous quasisymmetric configurations



- (a) Zarnstorff et al (2001)
- (b) Najambadi et al (2008)
- (c) Garabedian (2008)
- (d) Liu et al (2018)
- (e) Henneberg et al (2019)
- (f) Nuhrenberg & Zille (1988)
- (g) Anderson et al (1995)
- (h) Bader et al (2020)

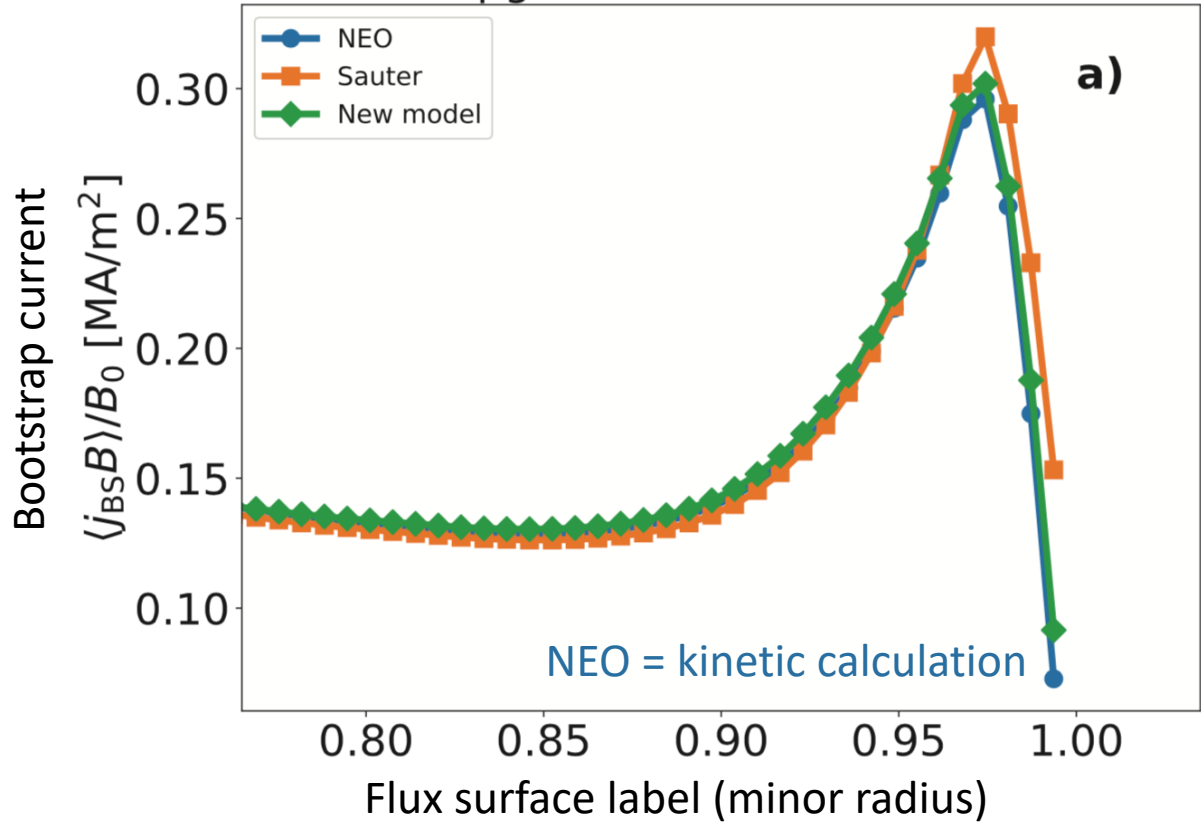
We want
 $B = B(r, \theta - N \varphi)$

Is there an optimization recipe that can give consistently straight $|B|$ contours?

New idea: exploit quasisymmetry & use analytic expressions for tokamaks

Redl (2021)

ASDEX Upgrade #33173, time = 4.75sec



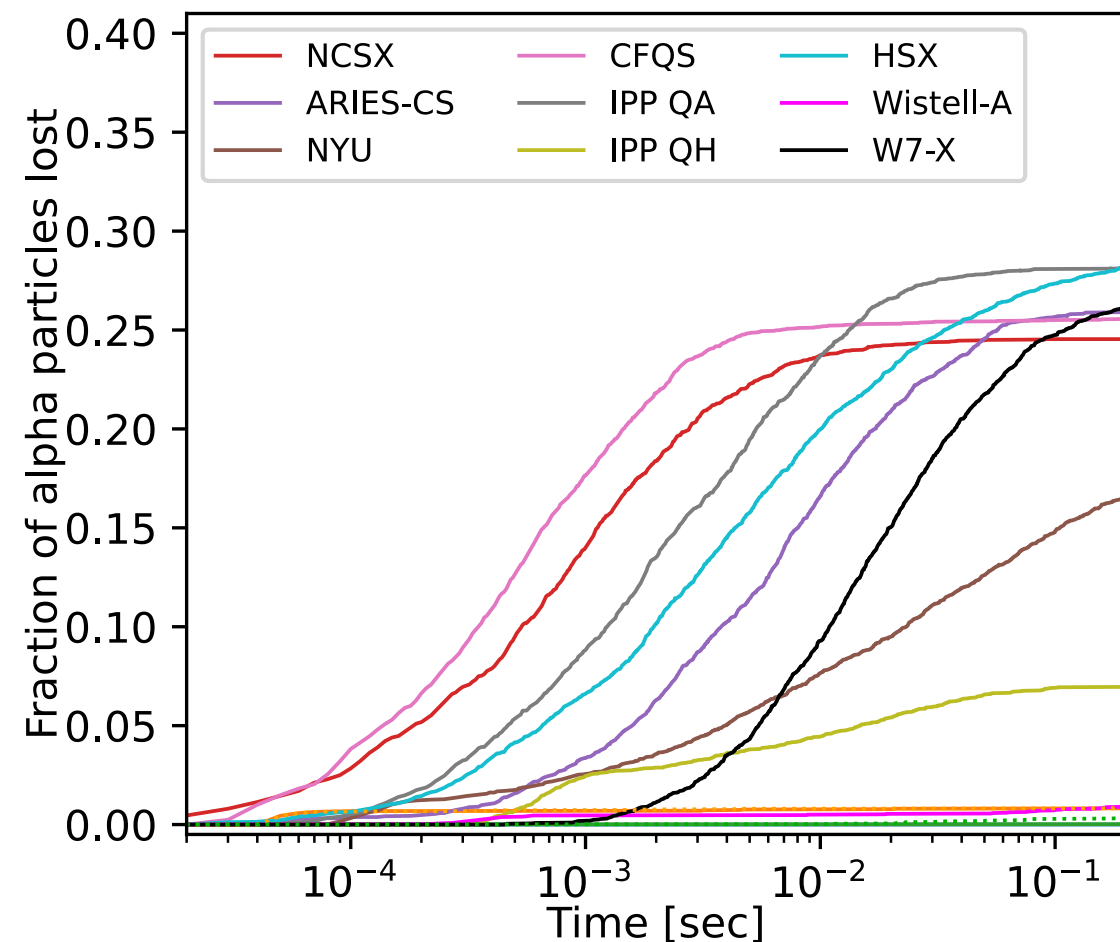
Geometry enters through

$$f_t = 1 - f_c = 1 - \frac{3}{4} \langle B^2 \rangle \int_0^{1/B_{\max}} \frac{\lambda d\lambda}{\langle \sqrt{1 - \lambda B} \rangle}$$

$$\nu_{e*} = 6.921 \cdot 10^{-18} \frac{q R n_e Z \ln \Lambda_e}{T_e^2 \epsilon^{3/2}},$$

$$\nu_{i*} = 4.90 \cdot 10^{-18} \frac{q R n_i Z^4 \ln \Lambda_{ii}}{T_i^2 \epsilon^{3/2}},$$

The symmetry yields extremely good confinement of collisionless trajectories

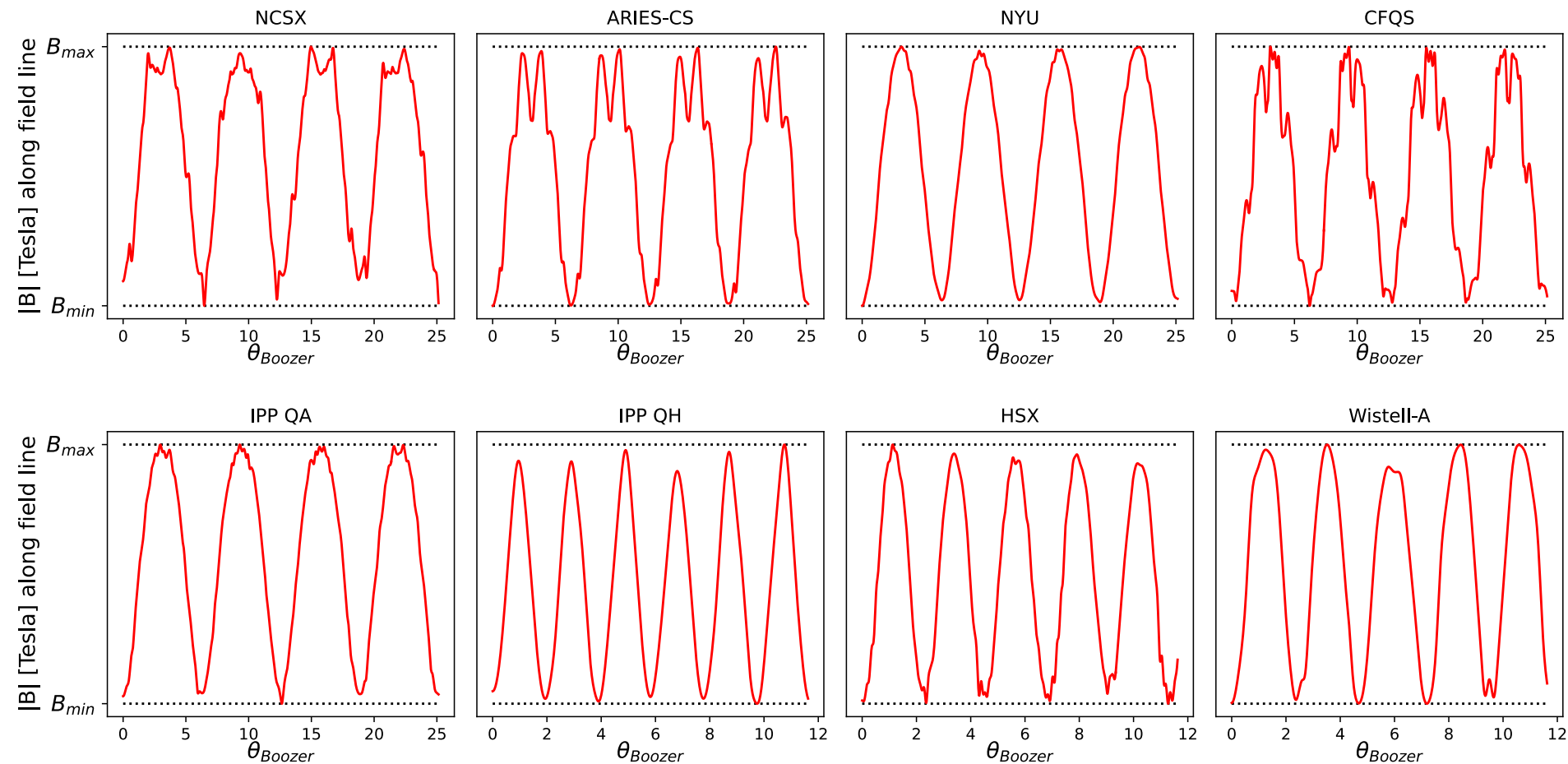


All configurations scaled to ARIES-CS minor radius (1.7 m) and $|B|$ (5.7 T).

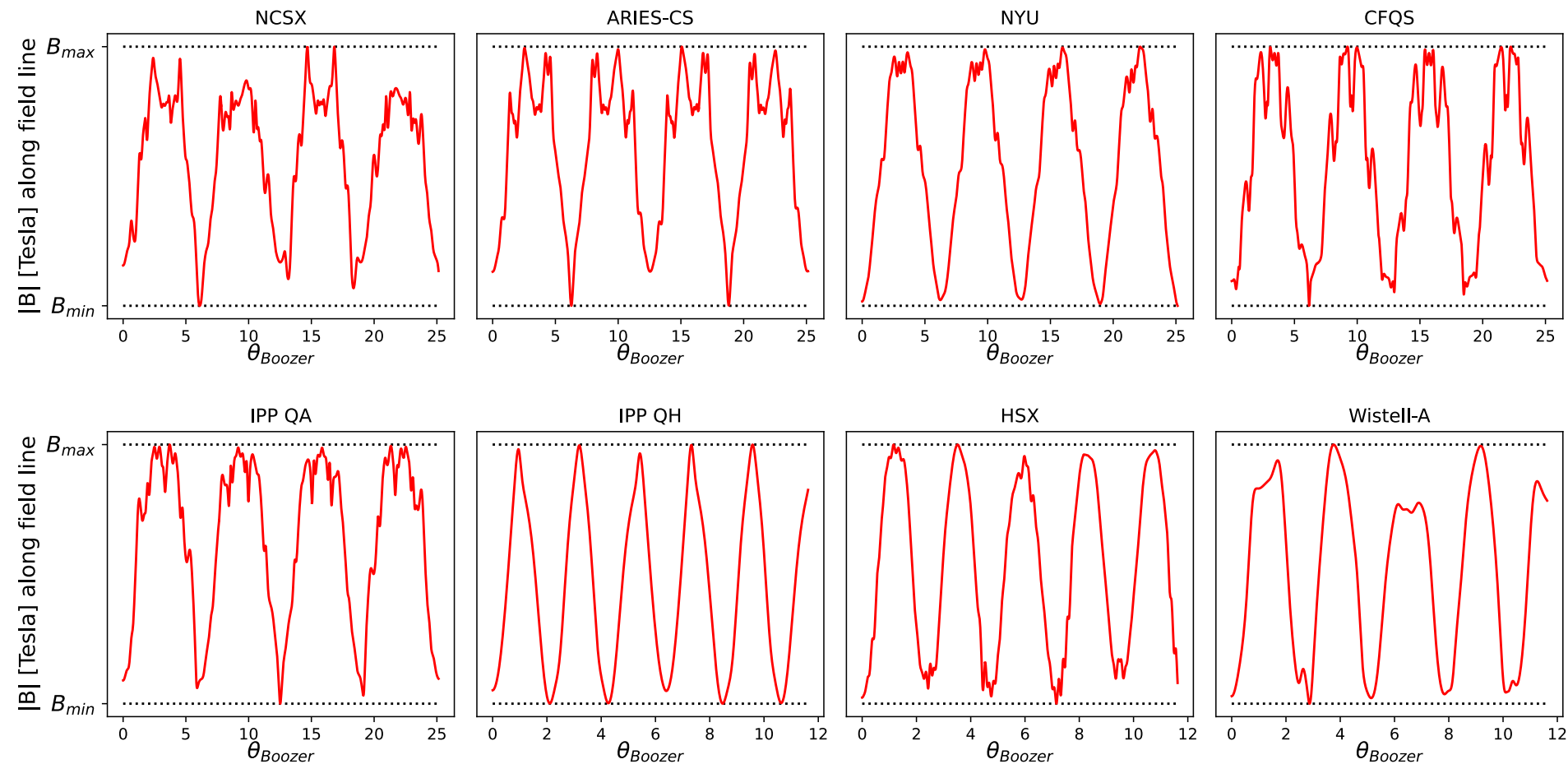
5000 alpha particles initialized isotropically at $s=0.3$.

SIMPLE code: Albert et al, JCP (2020).

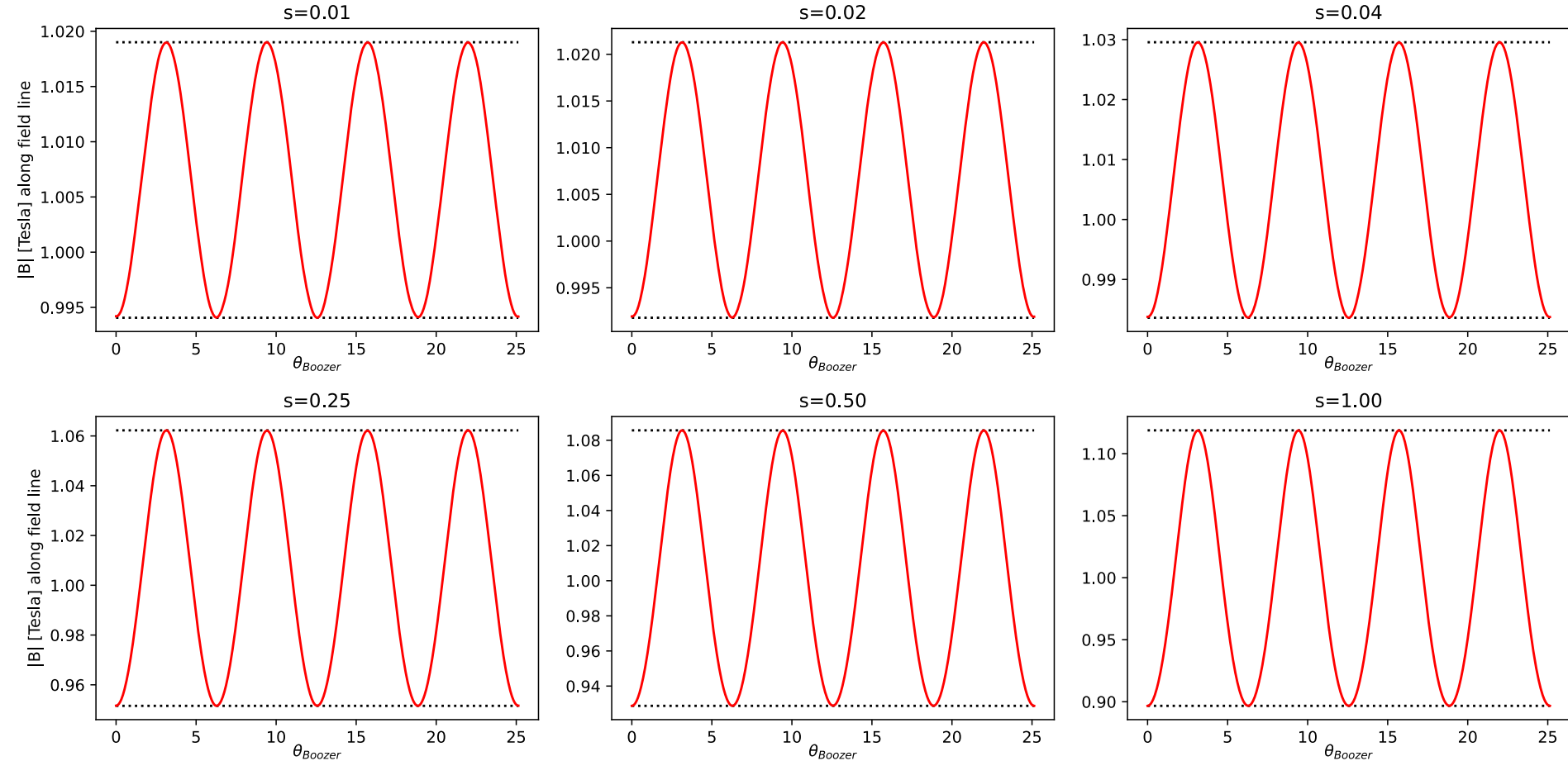
Previous quasisymmetric configurations (s=0.5)



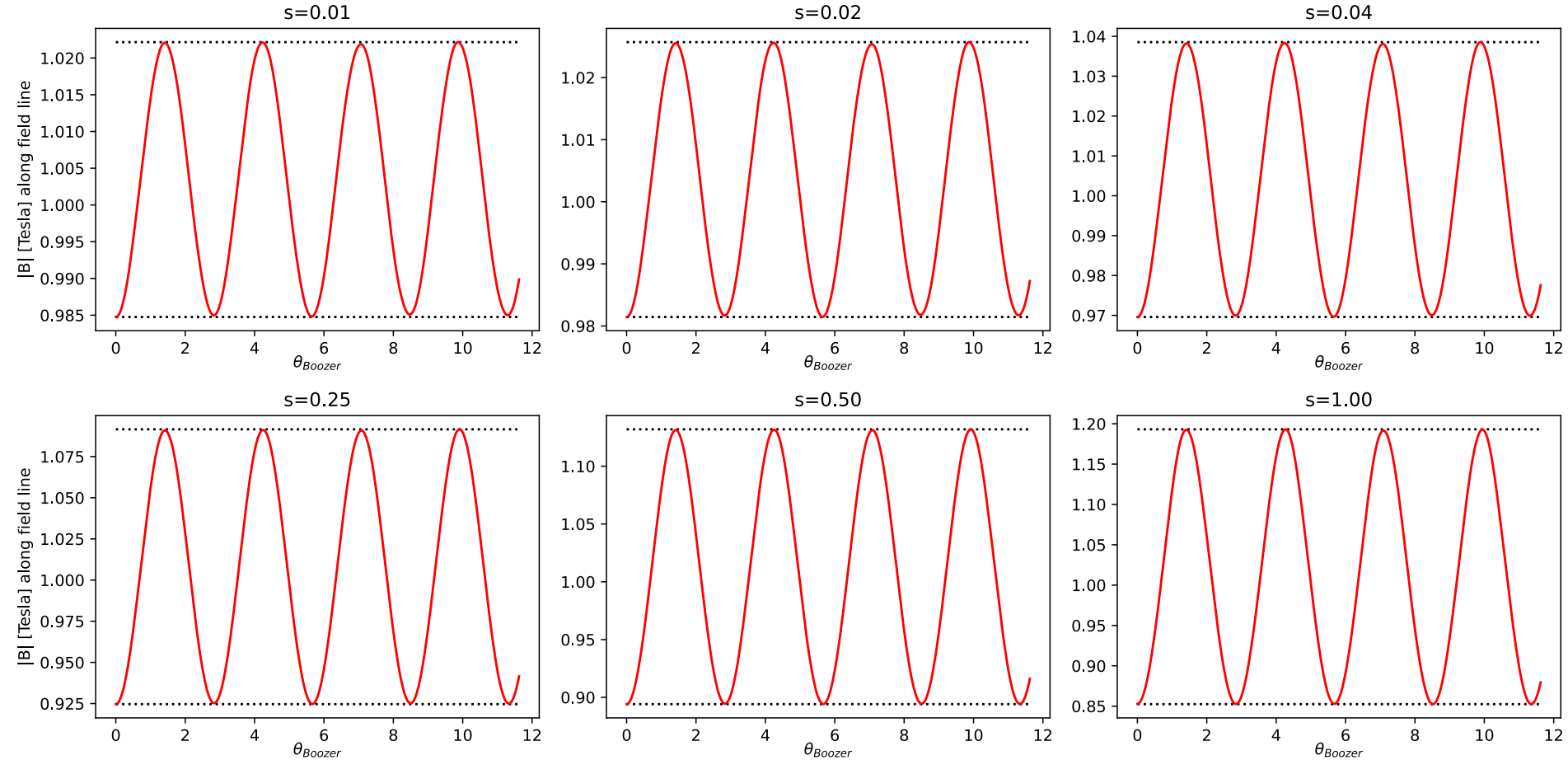
Previous quasisymmetric configurations (s=1)



|B| along a field line for new QA



|B| along a field line for new QH



|B| along a field line for new QA with magnetic well

