

12/7/12

Partial Wave Analysis

Recap: One-Dimensional Schrödinger Equation

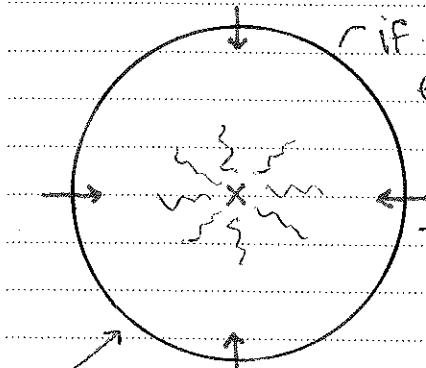
$$① \left[-\frac{\hbar^2}{2m} \frac{d^2}{dr^2} + V(r) + \frac{\hbar^2 l(l+1)}{2mr^2} \right] U_l(r) = E U_l(r)$$

Boundary condition: $U_l(r=0) = 0$

$$② U_l(r) \xrightarrow[r \rightarrow \infty]{} A \sin(kr - \frac{l\pi}{2} + \delta_l(k)) \text{ phase shift}$$

$$③ \frac{d\sigma}{d\Omega} = |f(\theta)|^2 = \left| \sum_{l=0}^{\infty} \frac{(2l+1)}{k} (e^{i2\delta_l} - 1) P_l(\cos(\theta)) \right|^2$$

$$\sigma = \frac{4\pi}{k^2} \sum_{l=0}^{\infty} (2l+1) \sin^2(\delta_l)$$



if there is no potential, wavefunction will evolve in time like it shrinks and then expands (free particle)

incoming/outgoing
spherical wave

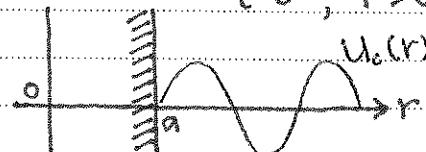
- if there is a potential, complicated things may happen, but at infinity the wavefunction will look like the free particle because of unitarity, but it will have picked up some phase shift, $\delta_l(k)$
* note that this is still the TISE!

But our initial condition is a plane wave...

→ a plane wave is just a linear combination of spherical waves, so instead of a single ' l ', we now have all ' l 's' (different partial waves will have different phases)

Example: Hard Sphere (S-wave $\Rightarrow l=0$)

$$-\frac{\hbar^2}{2m} U_0'' + \begin{cases} \infty, & r < a \\ 0, & r > a \end{cases}, \quad U_0(r) = E_0 U_0(r)$$



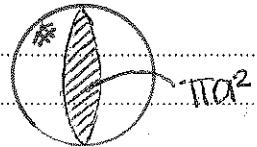
$$\Leftrightarrow U_0(r) = A \sin(kr + S_0)$$

Boundary condition: $U_0(r=a) = 0$

$$\Leftrightarrow ka + S_0 = 0$$

$$S_0 = -ka$$

$$\sigma = \frac{4\pi}{k^2} \sin^2(ka) \xrightarrow{ka \ll 1} 4\pi a^2$$



\downarrow 4 times the geometric cross-section of the sphere
 small energy approximation

If we were to go to higher l -values, we would need SBF's.

Things to know about spherical bessel functions:

1. Asymptotic behavior
2. Where the zeroes are

Example: Soft Sphere (S-Wave)

$$\text{Diagram of a soft sphere potential } V(r) = \begin{cases} 0 & r < a \\ V_0 & r > a \end{cases}$$

$$U(r) = \begin{cases} A \sin(kr) & r < a \\ B \sin(kr + \delta) & r > a \end{cases}$$

$$\downarrow k = \frac{1}{\hbar} \sqrt{2m(E + V_0)}$$

$$\uparrow k = \frac{1}{\hbar} \sqrt{2mE}$$

Boundary condition: $U(a_+) = U(a_-)$

$$\rightarrow A \sin(ka) = B \sin(ka + \delta)$$

Boundary condition: $U'(a_+) = U'(a_-)$

$$\rightarrow A k \cos(ka) = B k \cos(ka + \delta)$$

\downarrow although there are many unknowns
 we only need to find this
 (normalization arbitrary)

$$\Rightarrow \tan(\delta) = \frac{\frac{B}{A} \tan(ka) - \tan(ka)}{1 + \frac{B}{A} \tan(ka) \tan(ka)}$$

$$\xrightarrow{ka \ll 1} \delta \approx \tan(\delta) \approx ka \left[\frac{\tan(ka)}{ka} - 1 \right]$$

$$\tan \rightarrow \infty$$

$$\text{There's an energy bound } ka = \frac{\pi}{2} \Rightarrow \delta = \frac{\pi}{2}$$

bound states appear

\uparrow largest cross section possible!

$$\sin^2(\frac{\pi}{2}) = 1$$

\rightarrow You reach the unitarity bound!