

Lecture 6 - Fine Structure Corrections.

02/16/16

1. Level Splitting

2. Hyperfine splitting

1. Level Splitting

2S, 2p Caused by the Dirac equation $\Delta E_{fs} \sim \alpha^4 mc^2$

2P_{3/2}

10,950 MHz

2S_{1/2}

2P_{1/2} } 1,057 MHz

The Lamb Shift, λ

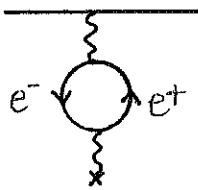
Feynman Diagrams: λ -Shift

e^- →



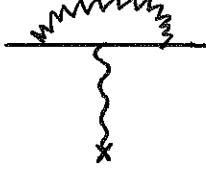
virtual photon; gauge boson of the electromagnetic interaction

this electron/photon interaction can be described as Vacuum Polarization



the photon splits into an e^-/e^+ pair then recombines
this process/phenomenon screens the field of the nucleus → stronger screening
(field screening actually infinite)

This screening causes a sort of "self-energy"



→ the self-energy shifts the quantized energy levels of the electron by adding extra jiggling.

* We must use renormalization theory to quantify these values
(i.e., see them as non-infinite)

extra jiggle = δr^i

$$\delta r^i = 0$$

↑ average jiggle ↑ jiggle in 2D

but the $(\text{average jiggle})^2 \neq 0$

$$\delta r^i \delta r^j = \frac{1}{3} \delta \bar{w}$$

$\rightarrow \delta r^2$ is the size of the extra jiggling

Energy of the jiggling:

$$U(\vec{r} + \delta\vec{r}) = U(\vec{r}) + \frac{1}{2} (\partial_i \partial_j U) \underbrace{\delta r^i \delta r^j}_{\frac{1}{3} \delta^4 r^2} + \dots \text{ (Taylor expansion)}$$

$$= U(\vec{r}) + \frac{\delta r^2}{6} \underbrace{\nabla^2 U}_{\propto}$$

$\propto e_0 \delta^3(\vec{r})$, raises energy of S-states
(and only at the origin)

$$\Delta E_{\text{lamb}} \propto \alpha^5 mc^2$$

* note: vacuum polarization lowers the energy of S-states

2. Hyperfine splitting / structure

We must consider the dipole moment of the H-atom

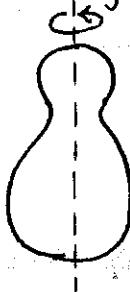
Nuclear magnetic moment - about a factor of 2,000 smaller than

$$\vec{M} = \frac{ze_0 g_N}{2M_N c} = \vec{I} \quad \text{the electronic magnetic moment}$$

$\propto \frac{1}{2M_N c}$ total angular momentum of the nucleus

$\propto \frac{1}{m_N}$ mass of the nucleus

Suppose you have any axes-symmetric volume [LIKE A PEANUT!] spinning on the axis:



angular momentum $\int \rho_m \vec{v} dV$; momentum is mass-

$$\vec{L} = \int \vec{r} \times \vec{p} dV \quad \text{density dependent}$$

electronic magnetic moment

$$\vec{\mu} = \frac{1}{c} \int \vec{r} \times \vec{j} dV$$

$\propto \int \rho_e \vec{v} dV$; current is charge-density dependent

Comparing the two:

$$\vec{\mu} = \frac{1}{c} \left(\frac{\rho_e}{\rho_m} \right) \vec{J}$$

\propto the g-factor, which accounts for the fact that e.g., an electron, does not have $\vec{\mu} \propto \frac{e}{m} \vec{J}$ but actually $\propto 2\vec{J}$

$$\Rightarrow \mu_e = \frac{-e_0 \cdot 2}{2m_e c} \frac{\hbar}{2}$$

How does this \vec{M} affect the energy levels?

$$\underbrace{\vec{M}_N \rightarrow \vec{B}_N}_{\text{turns into a little magnet}}$$

this leads to a hyperfine splitting of:

$$\rightarrow H_{hyp} = \frac{e_0}{mc} \underbrace{\vec{S} \cdot \vec{B}_N}_{\text{dipole moment}}$$

dipole moment (spin interaction with the generated magnetic field)

Proton magnetic moment μ_p electron magnetic moment μ_e

For the correction on the spatial part of the wavefunction,
(use perturbation theory)

$$\langle n, j=\frac{1}{2}, l=0 | H_{hyp} | n, j=\frac{1}{2}, l=0 \rangle = \frac{4}{3} g_N \left(\frac{m}{M_N} \right) (z\alpha)^4 mc^2 \frac{1}{n^3} \frac{\vec{S} \cdot \vec{I}}{\hbar^2} \quad \begin{matrix} \text{must diagonalize} \\ \text{this term} \end{matrix}$$

to diagonalize $\vec{S} \cdot \vec{I}$, we must first calculate \vec{I}

total angular momentum of the atom

$$F = \vec{S} + \vec{I} \quad \begin{matrix} \text{electronic angular momentum} \\ - \text{nuclear angular momentum} \end{matrix}$$

$$F^2 = (\vec{S} + \vec{I})^2 = S^2 + I^2 + 2\vec{S} \cdot \vec{I}$$

$$\frac{\vec{S} \cdot \vec{I}}{\hbar^2} = \frac{1}{2} \left(\frac{F^2}{\hbar^2} - \frac{3}{2} \right) ; \quad \frac{F^2}{\hbar^2} \leftrightarrow F(F+1)$$

hyperfine splitting is sized according to this value of F

$$\Delta E_{hyp} \propto \frac{m}{M} \alpha^5 mc^2$$