

Lecture 3 - Building Up To The Hilbert Space

09/04/15

1. Problem #2

2. Beam Splitters

3. Matrix Methods

(normalization/unitary)

4. Problem #3

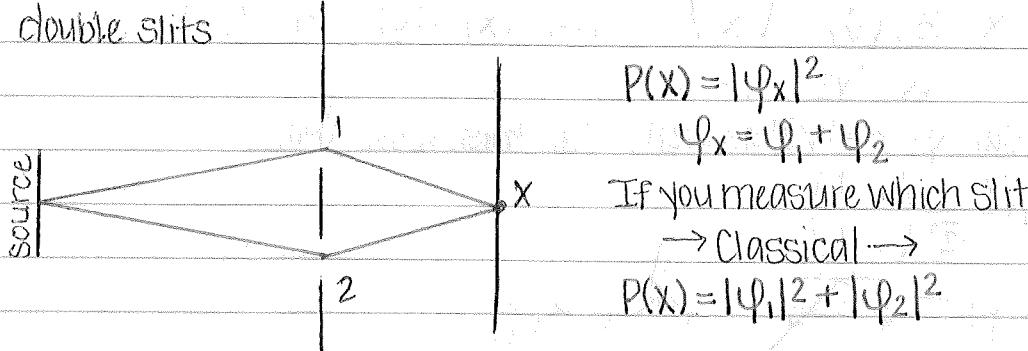
5. Testing Bombs

6. Spin \rightarrow Stern-Gerlach

7. Problem #4

1. Problem #2 Interference

double slits



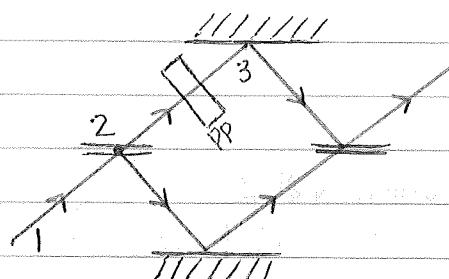
$$P(X) = |\psi_X|^2$$

$$\psi_X = \psi_1 + \psi_2$$

If you measure which slit
 \rightarrow Classical \rightarrow

$$P(X) = |\psi_1|^2 + |\psi_2|^2$$

2. Beam Splitters



A wave packet, $\psi = \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$, as
a function of time:

$$@1(0) \quad @2(\alpha) \quad @3(\alpha') \quad @4(\beta')$$

a definite known state
(no superpositions, etc.)

From Lecture 2...

Components of an interferometer:

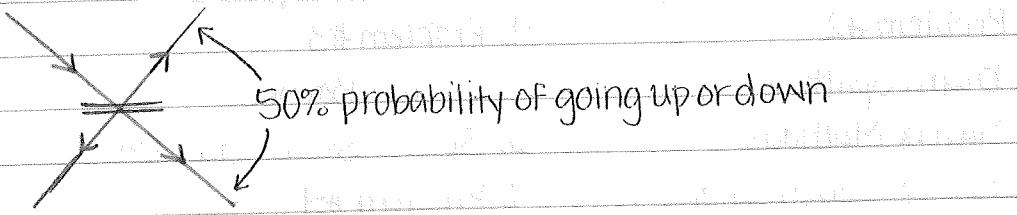
① Detector $P = |\psi|^2$

② Phase plate $\begin{pmatrix} e^{i\phi} & 0 \\ 0 & 1 \end{pmatrix}$ or $\begin{pmatrix} 1 & 0 \\ 0 & e^{i\phi} \end{pmatrix}$

③ Beam Splitter $\begin{pmatrix} w & y \\ x & z \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \begin{pmatrix} \alpha' \\ \beta' \end{pmatrix}$

Matrix for "balanced beam splitter"

\rightarrow Splits photons 50/50



From the top:

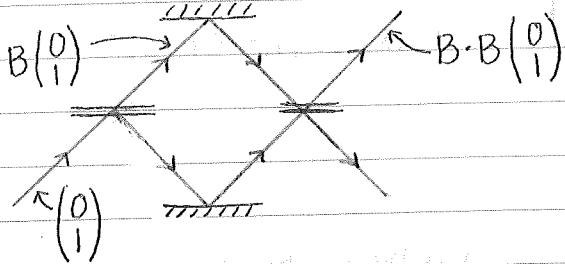
$$\begin{pmatrix} W & Y \\ X & Z \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} W \\ X \end{pmatrix} \Rightarrow |W|^2 = |X|^2 = \frac{1}{2} \rightarrow \text{mod} = \frac{1}{\sqrt{2}}$$

$|W| = |X| = |Y| = |Z| = \frac{1}{\sqrt{2}}$

$W = \frac{1}{\sqrt{2}}; X = \frac{1}{\sqrt{2}}$

- Can $Y = Z = \frac{1}{\sqrt{2}}$ as well? Can this work? (no)

$$B = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$$



$$B \cdot B \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

Probability of 2
NOT POSSIBLE (must be < 1)

There must be a sign change somewhere!

$$\Rightarrow B = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{pmatrix}$$

* Must constrain norms and norms of sums
in Quantum Mechanics problems

Matrices must "conserve probability"

3. Matrix Methods.

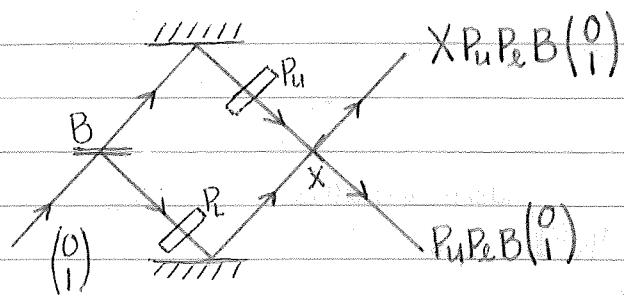
(normalization/unitary)

X matrix

↳ direction change / phase shift

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

How do we combine all the interferometer components?



$$\text{for } \psi = \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$

$$\psi \rightarrow B\psi \rightarrow P_u P_e B\psi \rightarrow X P_u P_e B\psi$$

- these are associative!

$$X(P_u P_e(B\psi)) = (X P_u P_e B)\psi$$

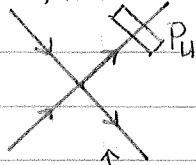
represents/describes the interferometer

$$P_u = \begin{pmatrix} e^{i\delta_u} & 0 \\ 0 & 1 \end{pmatrix} \quad P_e = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\delta_e} \end{pmatrix}$$

$$P_u P_e = \begin{pmatrix} e^{i\delta_u} & 0 \\ 0 & e^{i\delta_e} \end{pmatrix} = P_e P_u$$

the order of your phase plates does not matter

BUT, these are not commutative



$$P_u X \neq X P_u \Rightarrow \text{Observables do not commute}$$

cross matrix \Rightarrow beams intersect; no beam splitter

Normalization

- Can ψ be any matrix? (constraints on 2-state systems)

$$\psi = \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \Rightarrow |\alpha|^2 + |\beta|^2 = 1$$

state vector normalization constraint

$$(\alpha^* \beta^*) \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \psi^* \psi = 1$$

t: dagger operator = transpose & complex conjugate

constraint: ψ must normalize to 1

- Can the interferometer be any matrix?

for an interferometer matrix R

$$\psi \rightarrow R\psi = \psi'$$

$$\psi^\dagger \psi = 1$$

$$\Rightarrow \psi^\dagger = \psi^\dagger R^\dagger$$

$$\psi^\dagger \psi = (\psi^\dagger R^\dagger)(R\psi)$$

$$1 = \psi^\dagger (R^\dagger R) \psi$$

\hookrightarrow true IFF $R^\dagger R = 1$

identity matrix

$$1 \Rightarrow R^\dagger R = 1 \Rightarrow \text{Unitary } R$$

4. Problem #3

→ prove the converse is true

Show that every R (that is, every unitary matrix) can be made from P_u, P_e, X, B

→ Try to get every real R , then add phase plates orthogonal

a notation you might choose for the phase plates

$$\begin{pmatrix} \cos(\theta) & -\sin(\theta) \\ \sin(\theta) & \cos(\theta) \end{pmatrix}$$

5. Testing Bombs

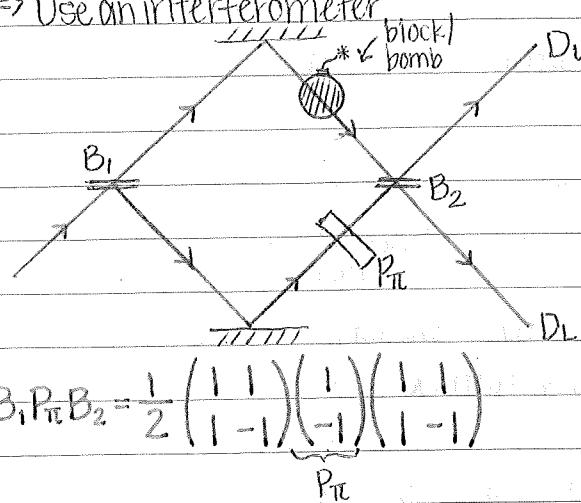
(ex. from 1993, pre-quantum computing)

There are bombs that are photon-activated

→ a single photon sent in... BOOM!!

But, some of the detectors don't work. How do you figure out which detectors work without exploding all the bombs?

⇒ Use an interferometer



$$B_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$B_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$B_1 P_{\pi} B_2 = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

P_{π}

$$= \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = R$$

Pr(B₂) ↑ unitary

only valid/applicable in
the quantum sense (ignore
when in classical conditions)

$R \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ — a photon shot in from the bottom branch comes out on top branch (nothing to D₂)

If we block the upper beam (classical → we know all photons in lower leg)

↳ 1/4 of the photons exit on the lower beam

↳ 50% probability with 50% of the photons

Now, we replace the block with a bomb (capacity unknown)

 "detector/observer if functional"

• Photon → upper leg ↳ You know it went through upper leg now

 → [Working] bomb → *BOOM* (w/ 50% chance)

• Photon → lower leg ↳ You know it went through lower leg now

 → [Working] bomb doesn't explode → D_U, D_L clicks

 1/4 1/4

By doing this experiment with the working bomb, interference is destroyed by the "observation" of which path the photon took.

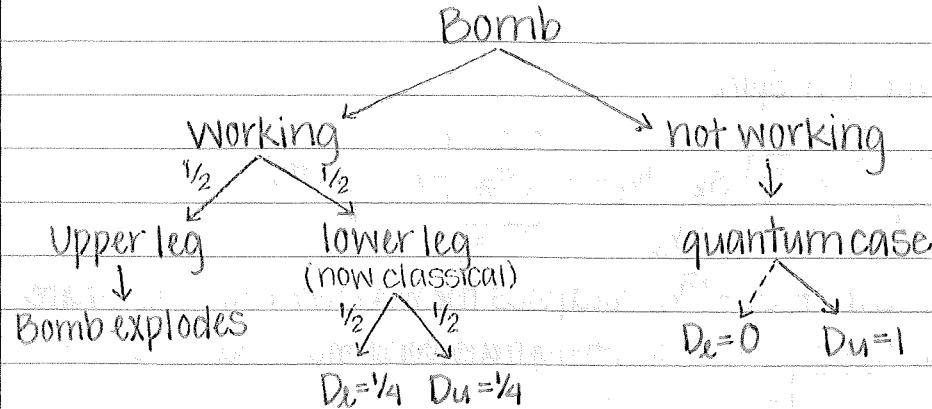
Working Bomb ⇒ forces to classical solution

 ↳ i.e. R not applicable

If the bomb is not working ⇒ quantum solution

 → click only at D_U (as a result of R... or by chance)

Possibilities



⇒ Click on D_U = definitive working bomb (without exploding it!)

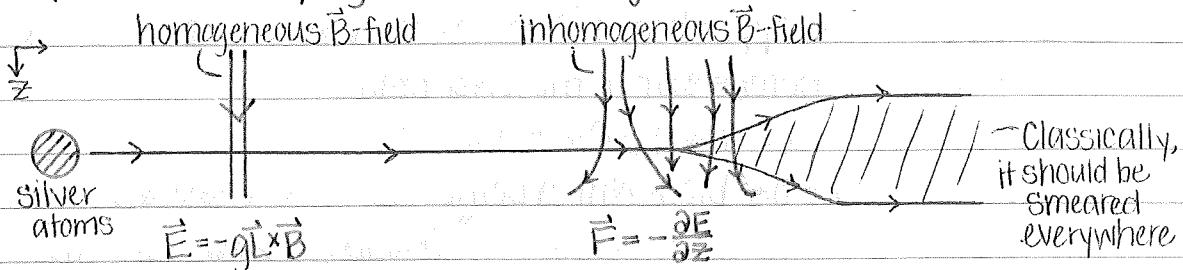
b. Spin: Stern-Gerlach

"basis independence"

Angular momentum is quantized

↳ this is a prediction of the Bohr Model of an atom

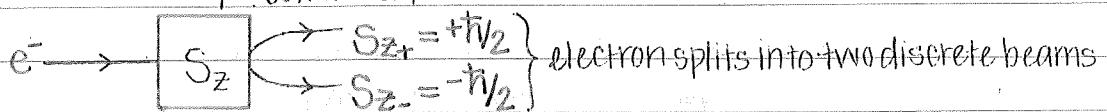
Experiment: Trying to measure angular momentum



↳ Despite classical expectations, the angular momentum was quantized into discrete bands

The Stern-Gerlach Machine

"box the experiment"



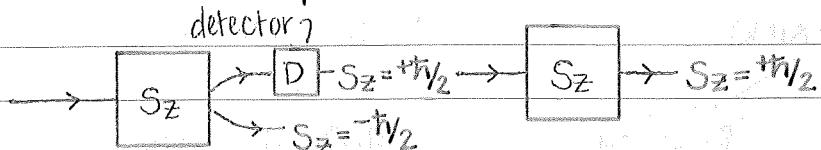
↳ S: Angular Momentum becomes Spin

$$v = \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \underbrace{\alpha(S_z = +\frac{\hbar}{2})}_{(1)} + \underbrace{\beta(S_z = -\frac{\hbar}{2})}_{(2)}$$

What's special about z ? Nothing! The Universe is rotationally invariant.

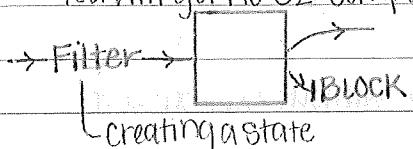
(i.e. You can make a Stern-Gerlach Machine along X, Y, etc.)

Measurement of Spin



Measurement of $S_z = +\hbar/2$ collapses the waveform into this state

⇒ You will get no S_z component from a second Stern-Gerlach



Rotate to X:

$$\rightarrow S_z = \frac{+i}{2} \rightarrow \boxed{S_x} \rightarrow \begin{cases} S_x = +\frac{i}{2} \\ S_x = -\frac{i}{2} \end{cases} \text{ even with an input of solely } S_z = \frac{+i}{2} !$$

$$Z_+ = \alpha' (S_x = +\frac{i}{2}) + \beta' (S_x = -\frac{i}{2})$$

X_+ X_-

$$|\alpha'|^2 = |\beta'|^2 = \frac{1}{\sqrt{2}} \quad (50/50 \text{ split})$$

"The notation is just going to get more inconvenient until we give into using bras & kets"

In reverse...

$$\rightarrow S_x = \frac{+i}{2} \rightarrow \boxed{S_z} \rightarrow \begin{cases} S_z = +\frac{i}{2} \\ S_z = -\frac{i}{2} \end{cases}$$

$$X_+ = \alpha (S_z = +\frac{i}{2}) + \beta (S_z = -\frac{i}{2}), \quad |\alpha|^2 = |\beta|^2 = \frac{1}{\sqrt{2}}$$

Given a state, you can choose its phase once

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} Z_+ + \frac{1}{\sqrt{2}} Z_-$$

$$X_- = \alpha_- (S_z = +\frac{i}{2}) + \beta_- (S_z = -\frac{i}{2})$$

$$|\alpha_-|^2 = |\beta_-|^2 = \frac{1}{\sqrt{2}}$$

In analogy with the beam splitter matrix:

to conserve probability, if $\alpha = \beta = \frac{1}{\sqrt{2}}$

$$\Rightarrow \alpha_- = -\beta_- = \frac{e^{i\frac{\pi}{4}}}{\sqrt{2}}$$

B₋ must be $\alpha/\sqrt{2}$

$$X_z = \frac{Z_+ \pm Z_-}{\sqrt{2}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} Z_z$$

7. Problem #4

→ repeat all this with y

hint: introducing S_y creates complex numbers

$$\rightarrow S_x = \frac{+i}{2} \rightarrow \boxed{S_y} \rightarrow \begin{cases} S_y = +\frac{i}{2} \text{ (50%)} \\ S_y = -\frac{i}{2} \text{ (50%)} \end{cases}$$

$$Y_{\pm} = (\alpha_{\pm} Z_+ + \beta_{\pm} Z_-)$$

⇒ α_{\pm} and β_{\pm} can't be wholly real

$$Y_{\pm} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ i & -i \end{pmatrix} Z_{\pm}$$

this choice is not unique for y, but there is no choice that allows them to be real

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

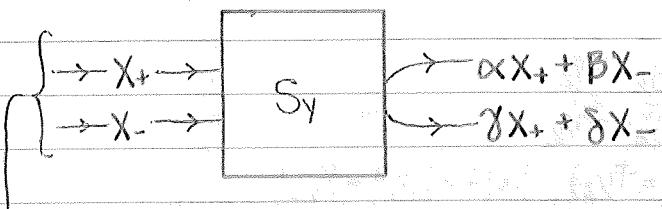
8. Orthogonality

ex: $X_+^T X_- = 0$

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = 0 \text{ orthogonal!}$$

Prove this using known normalization

Normalization constraint: $X_+^T X_+ = X_-^T X_- = 1$



Combination must be normalized

$$(\alpha X_+ + \beta X_-)^T (\alpha X_+ + \beta X_-) = |\alpha|^2 X_+^T X_+ + |\beta|^2 X_-^T X_- + (\alpha^* \beta X_+^T X_- + \alpha \beta^* X_+ X_-^T)$$

Where $|\alpha|^2 + |\beta|^2 = 1$

$$= |\alpha|^2 + |\beta|^2 + (\alpha^* \beta X_+^T X_- + \alpha \beta^* X_+ X_-^T)$$

1 constraint

$$= 1 + (\alpha^* \beta X_+^T X_- + \alpha \beta^* X_+ X_-^T) = 1$$

so this must = 0

$$\Rightarrow 2 \operatorname{Re}\{\alpha^* \beta X_+^T X_-\} = 0$$

if I can change phase...

$$\Rightarrow X_+^T X_- = 0$$