

Lecture 15 - Measurement & Entanglement

Brittany's
Notes ↓

1. Measurement from Entanglement Perspective

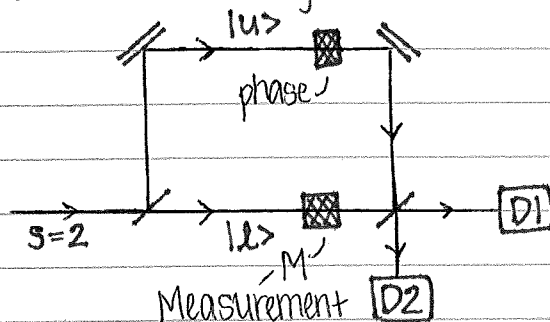
2. Density Matrices

3. Reduced Density Matrix

1. Measurement from the Entanglement Perspective

→ Back to interferometer scheme →

Particle with magnetic moment (spin 2)



Interference $P(D1) \propto \cos(\varphi)$

M = measurement

→ off: interference

→ on: Classical (no φ dependence)

Quantum description for M:

(in terms of entanglement)

$$M = \text{spin-1/2 } (|+x\rangle = \frac{|+z\rangle + |-z\rangle}{\sqrt{2}}) \text{ state}$$

$$= \text{pointer detector } \begin{cases} |+x\rangle \rightarrow |u\rangle \\ |-x\rangle \rightarrow |l\rangle \end{cases}$$

Hamiltonian -

$$\text{add } \mathcal{H}_{\text{int}} = g V \otimes Z$$

$\uparrow$ constant $\quad \uparrow S_z$ spin operator for M
 $\uparrow |l\rangle\langle l|$, spin of particle

$$U(t) = e^{-i\mathcal{H}_{\text{int}}t/\hbar}$$

→ choosing $gt/\hbar = \pi$ (maximally entangled state) →

$$= (|u\rangle\langle u|) \otimes \mathbb{I} + (|l\rangle\langle l|) \otimes (|+z\rangle\langle +z| - |-z\rangle\langle -z|)$$

$$U(t) \left(\frac{|u\rangle e^{i\varphi} + |l\rangle}{\sqrt{2}} \right) \otimes |+x\rangle$$

before time evolution

$$= [|u\rangle e^{i\varphi} \otimes |+x\rangle + |l\rangle \otimes \frac{1}{\sqrt{2}}(|+z\rangle - |-z\rangle)] \frac{1}{\sqrt{2}} \quad (\text{by linearity of the tensor product})$$

$$= \frac{[e^{i\varphi} |u\rangle \otimes |+x\rangle + |l\rangle \otimes |-x\rangle]}{\sqrt{2}}$$

if $t \neq \frac{\pi}{g}$:

$| -x \rangle \rightarrow \alpha | -x \rangle + \beta | +x \rangle$, "Partial or Weak Measurement"

* before, $|u\rangle \rightarrow | +x \rangle$ and $|l\rangle \rightarrow | -x \rangle$, but in the case of weak measurement $| +x \rangle$ & $| -x \rangle$ are not completely separated

$$\Rightarrow |\Psi\rangle \xrightarrow[\text{splitter}]{\text{beam}} \frac{1}{2} \left[e^{i\varphi} \left(\frac{|u\rangle + |l\rangle}{\sqrt{2}} \right) \otimes | +x \rangle + \left(\frac{|u\rangle - |l\rangle}{\sqrt{2}} \right) \otimes | -x \rangle \right]$$

↳ Measure at D1 ↘

$$\frac{1}{2} [e^{i\varphi} |u\rangle \otimes | +x \rangle + |u\rangle \otimes | -x \rangle] = \frac{1}{2} (e^{i\varphi} | +x \rangle + | -x \rangle) \otimes |u\rangle$$

Probability → NORM = $\frac{1}{4}(1+1) = \frac{1}{2}$

Without measurement:

$$t=0 \Rightarrow \beta=1$$

$$\rightarrow \frac{1}{4} |e^{i\varphi} | +x \rangle + | +x \rangle|^2 = \left(1 + \frac{\cos(\varphi)}{2} \right)$$

(interference)

2. Density Matrices

Classical vs. Quantum Uncertainty

Classical randomness $\begin{cases} \rightarrow | +z \rangle \text{ or } | -z \rangle \\ \rightarrow | +x \rangle \text{ or } | -x \rangle \end{cases}$

To represent this classical randomness, we can use a density matrix

For a set of states $\{ |\Psi_\alpha\rangle \}$ with classical probability p_α

$$\rho = \sum_{\alpha} p_{\alpha} |\Psi_{\alpha}\rangle \langle \Psi_{\alpha}|$$

then for some observable, $\hat{A}$

↳ expectation value

$$\langle \hat{A} \rangle = \sum_{\alpha} p_{\alpha} \langle \Psi_{\alpha} | \hat{A} | \Psi_{\alpha} \rangle$$

Conveniently: $\text{Tr}(\rho \hat{A}) = \langle \hat{A} \rangle$

$$\rho = |\Psi\rangle \langle \Psi| \text{ such that } \text{Tr}\{\rho \hat{P}\} = \langle \Psi | \hat{P} | \Psi \rangle$$

$| +x \rangle$ - Quantum uncertainty $\rightarrow \rho = | +; x \rangle \langle +; x |$

$$= \frac{1}{2} [| +z \rangle \langle +z | + | -z \rangle \langle -z | + | +z \rangle \langle -z | + | -z \rangle \langle +z |]$$

$$| +z \rangle, | -z \rangle \rightarrow \rho = \frac{1}{2} [| +z \rangle \langle +z | + | -z \rangle \langle -z |] = \frac{1}{2}$$

$$\begin{matrix} +\frac{1}{2} & -\frac{1}{2} & \dots \end{matrix} = \frac{1}{2} [| +x \rangle \langle +x | + | -x \rangle \langle -x |]$$

→

Traditional Measurement

(last problem, HW5)

start in $|\psi\rangle \rightarrow$ measure S_z $\begin{cases} \rightarrow |+\rangle \text{ with probability } |\langle +|\psi\rangle|^2 \\ \rightarrow |-\rangle \text{ with probability } |\langle -|\psi\rangle|^2 \end{cases}$

$$SO \Rightarrow \rho = \sum_{s=\pm} |S_z\rangle \langle S_z| (|\langle S_z|\psi\rangle|^2)$$

* Making a density matrix $\leftrightarrow$ entanglement

3. Reduced Density Matrix

A, B = Composite system; two Hilbert spaces

$$\mathcal{H}^A \otimes \mathcal{H}^B$$

$\{|a\rangle\}, \{|b\rangle\}$ set of composite eigenkets

$$|\psi\rangle = \sum_{a,b} C_{ab} |a\rangle \otimes |b\rangle$$

$\hookrightarrow$ most general wavefunction; composite pure state between A & B

for observable $P_A = P_A \otimes \mathbb{1} \Rightarrow \langle \psi | P_A | \psi \rangle$

$$\langle P_A \rangle = \sum_{a,b,a',b'} C_{ab}^* C_{a'b'} \langle a,b | P_A \otimes \mathbb{1} | a',b' \rangle$$

$$= \dots = \sum_{aa'} \underbrace{\sum_b C_{ab}^* C_{a'b}}_{\rho_{aa'}^{(A)} = \text{Reduced density matrix}}$$

$$= \text{Tr}(\rho^{\text{red}} P_A)$$

$$\hookrightarrow \text{Only if } \rho^{\text{red}} = \sum_b C_{ab}^* C_{a'b} |a'\rangle \langle a|$$