

G.O. of Magnonic mode about X-point 41

$$\partial_t^2 \psi = v^2 \nabla^2 \psi, \quad v^2 = B^2 / \rho \mu_0$$

$$\vec{B} = \hat{z} \times \nabla (Hxy)$$

$$\vec{B} = [-x, y] H, \quad B^2 = r^2 H^2$$

$$\frac{x}{y} \Big| \frac{x}{y} \frac{B^2}{H}$$

WKB dir. $\Rightarrow \omega^2 = \frac{k^2 B^2}{\rho \mu_0} = \frac{k^2 H^2 r^2}{\rho \mu_0}$

$$\frac{H^2}{\rho \mu_0} = 1 \Rightarrow \boxed{\omega^2 = k^2 r^2} \quad (1)$$

2D $\Rightarrow \vec{r} = \{x, y\}$ 4
 $\vec{k} = \{k_x, k_y\}$ vars

$$\omega \omega_{\vec{k}} = k^{\vec{r}} r^2, \quad \omega \omega_{\vec{r}} = k^{\vec{k}} r^2$$

$$\Rightarrow \boxed{\omega_{\vec{r}} = k^{\vec{r}} r^2, \quad \omega_{\vec{k}} = -r^2 k^{\vec{k}}} \quad (2)$$

H eqns

obvious

(3) Constants of motion

$$\boxed{\omega = kr} \quad (3a)$$

= const

$$\omega = \omega(k_x, k_y, x, y)$$

No obvious ignorables in Cartesian space

Try $\vec{D} = k \circ \vec{r}$

$$\begin{aligned} \vec{D} \cdot \vec{r} &= k \circ \vec{r} \cdot \vec{r} + k \circ \vec{r} \cdot \vec{r} \\ &= -\frac{k^2 r^2}{\omega} + \frac{k^2 r^2}{\omega} = 0 \end{aligned}$$

$\therefore \boxed{k \circ \vec{r} = D = \text{const}}$ (3)

(4) Direct differentiation of $r(t)$

$$\begin{aligned} r^2 &= \vec{r} \circ \vec{r}; & r r' &= \vec{r} \circ \vec{r}' \\ &= \frac{\vec{r} \circ k \vec{r}}{\omega} &= \frac{D r^2}{\omega} \end{aligned}$$

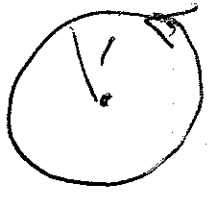
$\Rightarrow \boxed{r' = \frac{D}{\omega} r}$

let $\frac{D}{\omega} = \alpha \Rightarrow |\alpha| = \frac{|D|}{|\omega|} = \frac{|k \circ \vec{r}|}{|k r|} < 1$

$\therefore \boxed{r' = \alpha r}$ (4) any α , $|\alpha| < 1$

$$\Rightarrow r = r(0) e^{\alpha t}$$

$$r(0) = 1 \Rightarrow \boxed{r = e^{\alpha t}} \quad (5)$$

$$\alpha = 0 \Rightarrow r = 1 \rightarrow$$


$$\alpha > 0 \Rightarrow r \rightarrow e^{\alpha t}, \text{ increasing radius}$$

$$\alpha < 0 \Rightarrow r \rightarrow e^{-|\alpha|t}, \text{ radius} \rightarrow 0$$

(5) $\{x, k_x\}$ System

$$\omega x^0 = k_x r^2, \quad \omega k_x^0 = -x k^2$$

$r(t)$ is known, eliminate $\{k_y, y\}$

$$\text{use } \omega^2 = h^2 r^2 \Rightarrow \omega k_x^0 = \frac{-x \omega^2}{r^2}$$

$$\begin{aligned} \Rightarrow \omega x^{00} &= k_x^0 r^2 + 2 k_x r r^0 \\ &= \frac{-x \omega^2}{r^2} + 2 \frac{r r^0 \omega x^0}{r^2} \end{aligned}$$

$$\Rightarrow \mu \ddot{x} = -\mu x + 2 \frac{r^0}{r} \mu \dot{x}, \quad \frac{r^0}{r} = \alpha$$

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$$\Rightarrow \boxed{\ddot{x} - 2\alpha \dot{x} + x = 0} \quad (6)$$

(6) Try $x(t) \Rightarrow e^{\gamma t}$

$$\Rightarrow \gamma^2 - 2\alpha\gamma + 1 = 0$$

$$\Rightarrow \boxed{\gamma = \alpha \pm i\sqrt{1-\alpha^2}} \quad |\alpha| < 1$$

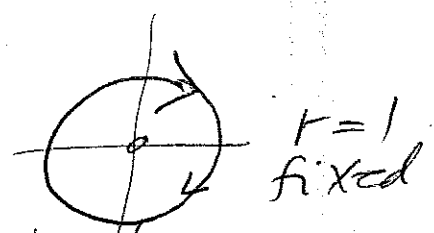
$$\Rightarrow \boxed{x(t) = e^{\alpha t} e^{\pm i\sqrt{1-\alpha^2} t} \quad (7)$$

$$\text{Re}(x) \Rightarrow \boxed{\text{Re} x(t) = r(t) \cos[\sqrt{1-\alpha^2} t]}$$

$$= r \cos \theta$$

$$\Rightarrow \theta = (\sqrt{1-\alpha^2})t$$

$$\alpha = 0 \Rightarrow \theta = t \Rightarrow r = 1 \Rightarrow$$



$1 > \alpha > 0 \Rightarrow \theta = \sqrt{1-\alpha^2} t \Rightarrow$ spiral out

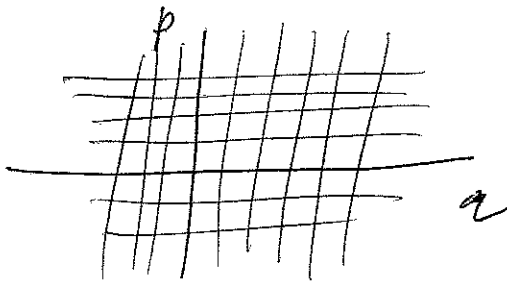
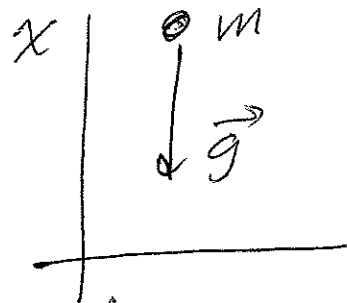
H-J Theory for g

$$T = \frac{1}{2} m \dot{x}^2$$

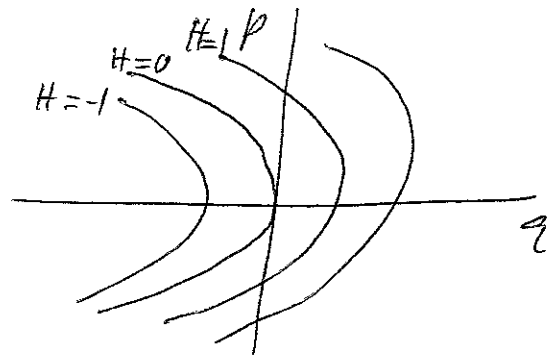
$$V = mgx \quad p = \frac{\partial L}{\partial \dot{x}} = m\dot{x}$$

$$H = \frac{1}{2} \frac{p^2}{m} + mgx$$

$$\text{let } m=1, g=1 \Rightarrow \boxed{H = \frac{1}{2} p^2 + q}$$



$\{p, q\}$ coords



H coord.

We want $\{p, q\} \rightarrow \{P, Q\} \ni P = \alpha = H$

$$\text{Now } K=0 \Rightarrow H + \frac{\partial S}{\partial t} = 0$$

$$\text{where } p = \frac{\partial S}{\partial q}, \quad \alpha = \frac{\partial S}{\partial P}, \quad S(q, \alpha, t)$$

$$\Rightarrow \frac{1}{2} \left(\frac{\partial S}{\partial q} \right)^2 + q + \frac{\partial S}{\partial t} = 0$$

$$\text{let } S(q, \alpha, t) = W(q, \alpha) - \alpha t$$

$$\Rightarrow \frac{1}{2} \left(\frac{\partial W}{\partial q^2} \right) = (\alpha - q)$$

$$\frac{\partial W}{\partial q} = \sqrt{2(\alpha - q)}$$

We pick the + sign, revisit if there is a problem

$$\Rightarrow \beta = \frac{\partial S}{\partial \alpha} = \frac{\partial W}{\partial \alpha} - t$$

$$\begin{aligned} \Rightarrow \beta + t &= \frac{\partial}{\partial \alpha} \int^q k(q') \sqrt{2(\alpha - q')} \\ &= \int^q dq' \frac{\partial}{\partial \alpha} \sqrt{2(\alpha - q')} \end{aligned}$$

$$= \int^q dq' (-) \frac{1}{2} \frac{\partial}{\partial \alpha} \sqrt{2(\alpha - q')}$$

$$(\beta + t) = -\sqrt{2(\alpha - q)}$$

$$\Rightarrow (\beta + t)^2 = 2(\alpha - q)$$

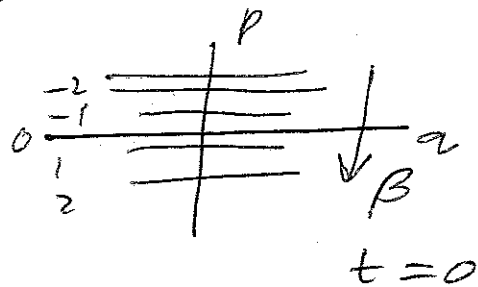
$$\Rightarrow \boxed{q = \alpha - \frac{1}{2} (\beta + t)^2} \Leftarrow q(\alpha, \beta, t)$$

$$p = \frac{\partial S}{\partial q} = \frac{\partial W}{\partial q} = \sqrt{2(\alpha - q)} = -\beta - t$$

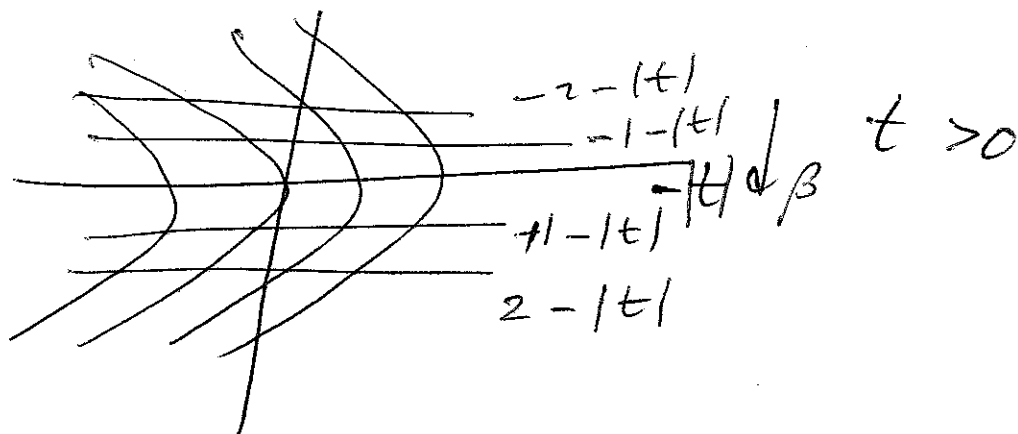
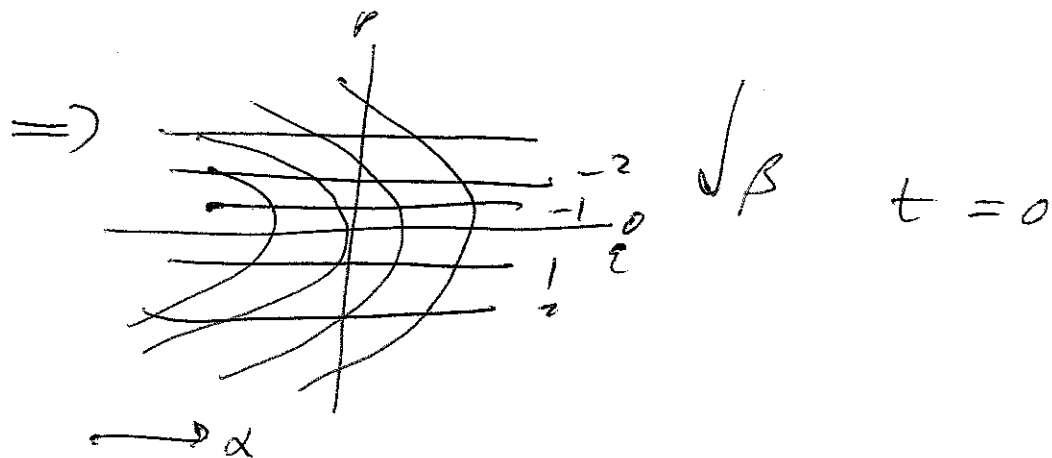
$$\boxed{p = -\beta - t} \Leftarrow p(\alpha, \beta, t)$$

Now find $\alpha(p, q, t)$, $\beta(p, q, t)$

$$\Rightarrow \beta = -p - t$$



$$\alpha = q + \frac{1}{2}p^2$$



β coordinates translate downward