

Solutions

Phys410/F16/Hassam/Midterm 1
1 hr 30 min 1 sheet 2-sided

Box all important results

Problem 1 (20 points)

A projectile mass m is shot from the origin at speed v_0 , at an angle α with respect to the horizontal x -axis. There is gravity downward. There is also a funny friction – the atmosphere is anisotropic so that the friction acts only in the x -direction. In particular, if v_x is positive, the friction force is $-bv_x$; but there is no friction in the y -direction. The projectile is shot from a cliff so it does not encounter ground for a long time. g is in the negative y direction.

- (a) From differential equations, find explicit solutions for $x(t)$ and $y(t)$.
- (b) Find the x -range of the projectile, defined as the x -distance traversed as $t \rightarrow \infty$. Express the answer in terms of v_0 , α , m , b , g .

Problem 2 (30 points)

There are 2 masses m_1 and m_2 which attract each other by a central force, ie, the attractive force is along the distance vector joining the masses. The forces are equal and opposite, ie, $\mathbf{F}_{12} = -\mathbf{F}_{21}$. Let the mass positions be given by $\mathbf{r}_1(t)$ and $\mathbf{r}_2(t)$. The angular momentum of m_1 is $\mathbf{L}_1 = m_1 \mathbf{r}_1 \times \mathbf{v}_1$, $\mathbf{v}_1 = d\mathbf{r}_1/dt$. Likewise, for m_2 . The total angular momentum of the system is $\mathbf{L} = \mathbf{L}_1 + \mathbf{L}_2$.

- (a) Prove whether or not \mathbf{L} is a constant of the motion. Your proof must start with the line

$$d\mathbf{L}/dt = \dots\dots$$

and must show all the important steps. (Can be done in 6-10 steps). The proof depends on two crucial steps (apart from Newton's Laws and action = reaction), one of which is mathematical and the other one mathematical and dependent on the properties of the interparticle force. You must clearly highlight these 2 steps.

- (b) A grad student experiments theoretically with the forces. Instead of \mathbf{F}_{12} as defined above, the student assumes that $\mathbf{F}_{12} = b \mathbf{r}_2 |\mathbf{r}_1 - \mathbf{r}_2|$, ie, the force on 1 due to 2 is proportional to the position vector of 2 and proportional to the distance between 1 and 2. (b is a constant.) In a symmetric manner, $\mathbf{F}_{21} = b \mathbf{r}_1 |\mathbf{r}_2 - \mathbf{r}_1|$.

1. (Yes/No only) Is the new force a central force (in general, as conventionally defined)?
2. (Yes/No only) Is the new interparticle force equal and opposite?
3. Prove whether or not the total angular momentum a constant of the motion?
(you do not have to repeat the entire proof as long as you can demonstrate a key step or two).

Problem 3 (30 points)

A force field is given as $\mathbf{F} = (-x, y)$, in 2D Cartesian coordinates. A particle mass $m=1$ is placed at the origin and given a kick so that its initial velocity (at $t=0$) is v_0 in the x-direction, ie, $\mathbf{v}(0) = (v_0, 0)$.

(a) From differential equations, by explicit calculation starting from Newton's Laws $m\mathbf{a} = \mathbf{F}$, find $x(t)$ and $y(t)$.

(b) At some time T , the mass comes to a momentary stop. What is the earliest T ? At what (x, y) location does this happen?

(c) We can explore some facets of the above problem by using Energy as a constant of the motion. Show which property of \mathbf{F} allows you to do this. Given this fact, obtain the potential energy $U(x, y)$. Write down the expression for the Energy constant in terms of U , etc, in general Cartesian coordinates.

(d) For the explicit solution found in (a), show that the energy at $t=0$ is equal to the energy at $t=T$, at the location where the mass comes to a momentary stop.

(e) Make a rough sketch of $U(x, y)$ contours. (Use this to confirm for yourself your findings above.)

Problem 4 (20 points)

A particle mass m moving along the x-coordinate is subject to a force $F(x) = -F_0 \sinh(\alpha x)$, $\alpha = \text{constant}$.

(a) Find the equilibrium position of the particle.

(b) Show that the equation satisfied by $x(t)$ for *small* oscillations of this mass about its equilibrium position is of the undamped harmonic oscillator type with some angular frequency ω_0 . Define ω_0 in terms of F_0 , m , α .

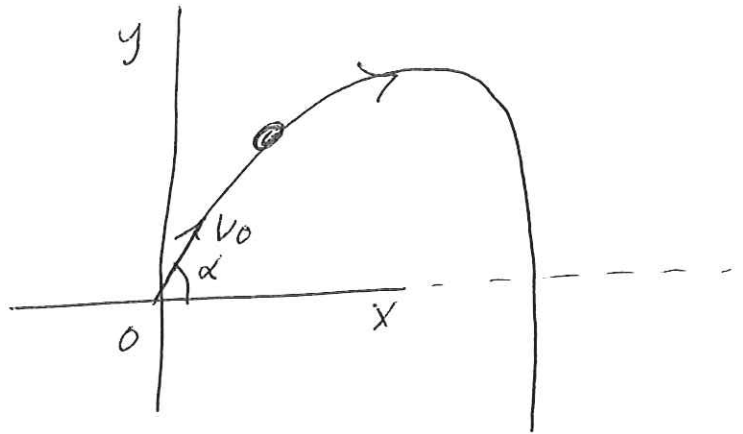
(c) The mass is subjected to an external driving force $f(t) = f_0 \cos(\omega t)$. Find an expression for the amplitude of small oscillations in response to the driving force if $\omega^2 = 0.99 \omega_0^2$. Compare this with the amplitude of oscillations for when $\omega \ll \omega_0$, or, effectively, when $\omega = 0$. Assume friction is negligible.

Projectile

$$m \ddot{x} = -bx$$

$$m \ddot{y} = -mg$$

$$\boxed{\begin{aligned} \ddot{x} &= -kx \\ \ddot{y} &= -g \end{aligned}} \quad k = b/m$$



$$\dot{x} = -kx + C \Rightarrow x = -kx + V_0 \cos \alpha$$

$$x_h = A e^{-kt}, \quad x_p = \frac{V_0 \cos \alpha}{k}$$

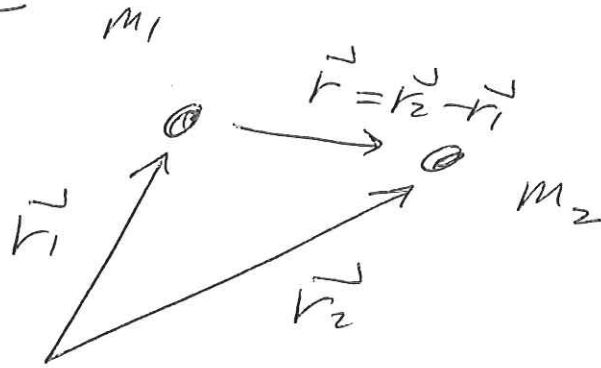
$$(a) \Rightarrow \boxed{x(t) = \frac{V_0 \cos \alpha}{k} (1 - e^{-kt})}$$

$$\dot{y} = V_0 \sin \alpha - gt$$

$$\boxed{y(t) = V_0 \sin \alpha t - \frac{gt^2}{2}}$$

$$(b) \underline{t \rightarrow \infty} \quad \boxed{x(t) \rightarrow \frac{V_0 \cos \alpha}{k} \equiv \text{Range}}$$

Angular Momentum



$$\vec{L} = \sum_k m_k \vec{r}_k \times \vec{v}_k, \quad \vec{v}_k = \dot{\vec{r}}_k$$

$$(a) \quad \vec{L}^0 = \sum_k m_k \left(\underbrace{\vec{r}_k \times \dot{\vec{r}}_k}_{=0} + \vec{r}_k \times \ddot{\vec{r}}_k \right) \leftarrow (\vec{r}_k \times \dot{\vec{r}}_k)^0 = \text{product rule}$$

$$= \sum_k m_k \vec{v}_k \times (\vec{F}_{k,e} / m_k)$$

$$= \vec{r}_1 \times \vec{F}_{12} + \vec{r}_2 \times \vec{F}_{21}$$

$$= \vec{r}_1 \times \vec{F}_{12} - \vec{r}_2 \times \vec{F}_{12}$$

$$= (\vec{r}_1 - \vec{r}_2) \times \vec{F}_{12} \quad \leftarrow \text{forces equal opposite}$$

$$= 0 \quad \leftarrow \text{because } \vec{r}_1 - \vec{r}_2 \parallel \vec{r} \parallel \vec{F}_{12} \text{ central force}$$

$$(b) \quad \vec{F}_{12} = b \vec{r}_2 |\vec{r}_1 - \vec{r}_2|$$

$$\vec{F}_{21} = b \vec{r}_1 |\vec{r}_1 - \vec{r}_2|$$

(b1) **No** $\vec{F}_{12} \parallel \vec{r}_2$. Central \vec{F}_{12} must be parallel to $(\vec{r}_1 - \vec{r}_2)$.

$$\vec{r}_2 \times (\vec{r}_1 - \vec{r}_2) = \vec{r}_2 \times \vec{r}_1 \neq 0 \text{ in general}$$

(b2) **No** If $A = R$, $F_{12} + F_{21} = 0$

$$\vec{F}_{12} + \vec{F}_{21} = b |\vec{r}_1 - \vec{r}_2| (\vec{r}_2 + \vec{r}_1) \neq 0 \text{ in general}$$

(b3) Based on the proof above,

Consider $\vec{r}_1 \times \vec{F}_{12} + \vec{r}_2 \times \vec{F}_{21}$

$$= (\vec{r}_1 \times \vec{r}_2) b |\vec{r}_1 - \vec{r}_2| + (\vec{r}_2 \times \vec{r}_1) b |\vec{r}_1 - \vec{r}_2|$$

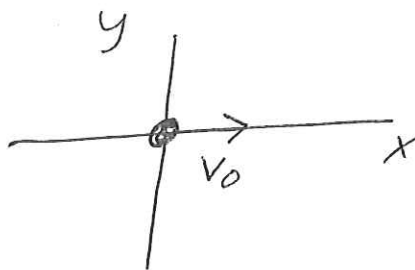
$$= b |\vec{r}_1 - \vec{r}_2| (\vec{r}_1 \times \vec{r}_2 + \vec{r}_2 \times \vec{r}_1)$$

$$= 0$$

$$\therefore \boxed{L = 0} \quad \text{Yes}$$

Fields

$$\vec{F} = (-x, y)$$



$$(a) \quad \ddot{x} = -x, \quad \ddot{y} = y$$

$$\Rightarrow \boxed{x(t) = v_0 \sin t} \quad \text{since } x(0) = 0 \\ \dot{x}(0) = v_0$$

$$y \in \left\{ \begin{matrix} e^t \\ e^{-t} \end{matrix} \right\} = A e^t + B e^{-t}$$

$$y(0) = 0, \quad \dot{y}(0) = 0$$

$$\Rightarrow \boxed{y(t) = 0}$$

$$(b) \quad \dot{x}(t) = v_0 \cos t$$

$$\dot{x}(T) = 0 \Rightarrow \cos T = 0 \Rightarrow \boxed{T = \pi/2}$$

$$\text{and } \boxed{x(T) = v_0} \text{ stopping distance}$$

$$(c) \quad \vec{\nabla} \times \vec{F} = \partial_x y - \partial_y (-x) = 0 \Rightarrow \vec{F} = -\vec{\nabla} U$$

$$\Rightarrow \boxed{\mathcal{E} = \frac{1}{2}(\dot{x}^2 + \dot{y}^2) + U(x, y) = \text{const}}$$

$$-dU = \vec{F} \cdot d\vec{r} = -x dx + y dy = \frac{1}{2} \left(-\frac{x^2}{2} + \frac{y^2}{2} \right)$$

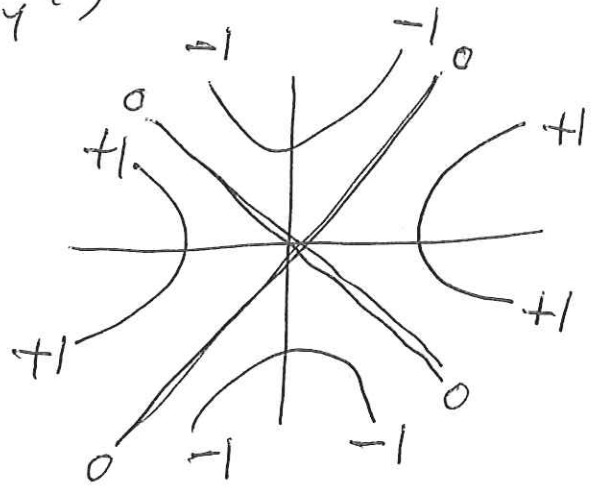
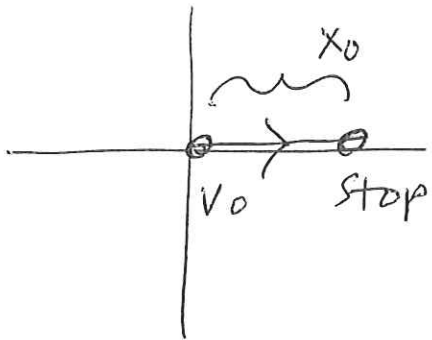
$$\Rightarrow \boxed{U = \frac{1}{2}(x^2 - y^2)}$$

$$(d) \quad \mathcal{E}(t=0) = \frac{1}{2} v_0^2 + U(0,0) = \frac{1}{2} v_0^2$$

$$\mathcal{E}(t=T) = U(v_0, 0) = \frac{1}{2} v_0^2$$

equal

~~U =~~ $U = \frac{1}{2} (x^2 - y^2)$



Mass goes "uphill" & momentarily stops cleanly, $\frac{1}{2} m v_0^2 = U(x_0, y=0)$
 $= \frac{1}{2} x_0^2$

$$\Rightarrow x_0 = v_0$$

Oscillations

$$m \ddot{x} = F(x) = -F_0 \sinh(\alpha x)$$

$$(a) \quad \ddot{x} = 0 \Rightarrow \sinh(\alpha x) = 0 \Rightarrow \boxed{x=0}$$

(b) let x be small, and assume stays small w.r.t. $x=0$

$$\Rightarrow \sinh(\alpha x) \approx \sinh(0) + \alpha \cosh(0) x \\ = \alpha x$$

$$\Rightarrow m \ddot{x} \approx -F_0 \alpha x$$

$$\Rightarrow \text{oscillations at } \boxed{\omega_0 = \sqrt{\frac{F_0 \alpha}{m}}}$$

$$(c) \quad m \ddot{x} = -F_0 \alpha x + f_0 \cos(\omega t)$$

$$\ddot{x} = -\omega_0^2 x + \frac{f_0}{m} \cos(\omega t)$$

$$\text{Try } x(t) = A \cos(\omega t)$$

$$\Rightarrow (\omega_0^2 - \omega^2) A = \frac{f_0}{m} \Rightarrow A(\omega)$$

$$\omega^2 = 0.99 \omega_0^2 \Rightarrow A = \frac{f_0/m}{\omega_0^2} 100 \Rightarrow \times 100$$

$$\omega^2 \rightarrow 0 \Rightarrow A = \frac{f_0/m}{\omega_0^2}$$