

Notes to accompany

Ch 4 - Part II

updated: 09/14/16

WE Thm for conservative \vec{F}

\Rightarrow Energy conservation

• If $W = \int_A^B d\vec{r} \cdot \vec{F}$

then WE Thm $\Rightarrow K = K_0 + W$

where $K = \frac{1}{2}mv^2$

• if \vec{F} is conservative, $\Rightarrow \vec{F} = -\vec{\nabla}U$ convention \swarrow

$\Rightarrow K_B = K_A - \int_A^B d\vec{r} \cdot \vec{\nabla}U$ \swarrow indep of path

$\Rightarrow K_B = K_A - \int_A^B dU$, because $d\vec{r} \cdot \vec{\nabla}U \equiv dU$
 $= K_A - U_B + U_A$

$\Rightarrow K_B + U_B = K_A + U_A$

$\Rightarrow \boxed{K + U \equiv \mathcal{E} = \text{const}}$

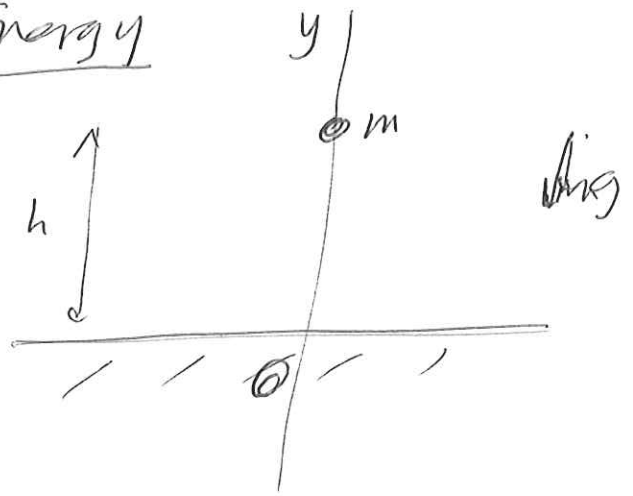
conservation of sum Kinetic + Potential Energy

Free fall in \vec{g} by Energy

w/2

Find K at ground

Let $K_h = 0$ @ h .



$$U = mgy$$

$$\Rightarrow K_h + mgh = K_0 + mg(0)$$

$$\Rightarrow \boxed{K_0 = mgh} \Rightarrow \boxed{V_0^2 = 2gh}$$

Direct Solution
(ie, not new)

$$v_y = -gt$$

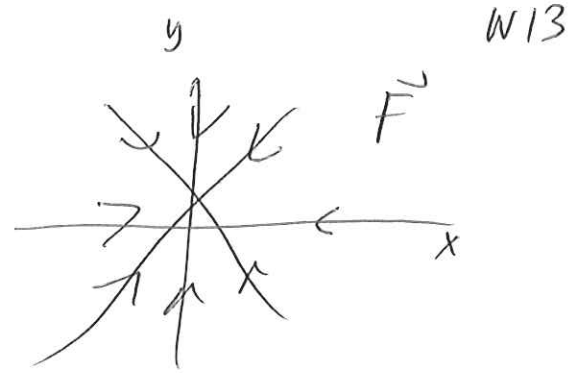
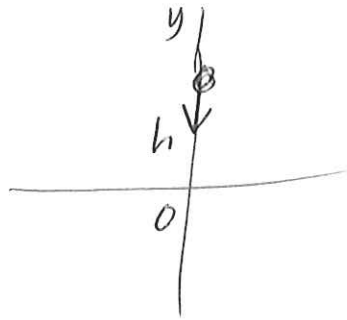
$$y = h - \frac{1}{2}gt^2$$

$$y = 0 \Rightarrow t^2 = 2h/g$$

$$\Rightarrow v_y(0)^2 = (gt)^2 = g^2 \frac{2h}{g}$$

$$\Rightarrow \boxed{V_0^2 = 2gh}$$

Free Fall in $\vec{F} = -k\vec{r}$



$$U = \frac{k}{2} (x^2 + y^2), \quad \vec{F} = -\vec{\nabla}U = -k\vec{r}$$

o Find K_0 if $K_h = 0$ and fall is directly along y axis

$$K + U = \text{const}$$

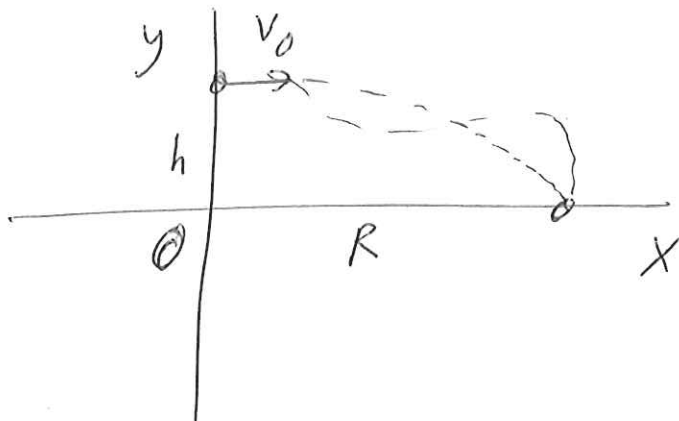
$$\Rightarrow U_h = K_0 + U_0$$

$$U_h = \frac{1}{2}kh^2 \quad U_0 = 0$$

$$\Rightarrow \boxed{K_0 = \frac{1}{2}kh^2}$$

"Projectile" in different \vec{F} fields

W14



Given mass starts at $(0, h)$ with $\vec{v}_0 = \hat{x}v_0$, ends at $(R, 0)$

Find $K(R, 0)$ for $\vec{F} = -\vec{\nabla}U$

for ∇U 3 cases: $U_1 = mgy$, $U_2 = \frac{1}{2}k(x^2 + y^2)$, $U_3 = kxy$

In all cases, $\frac{1}{2}mv_0^2 + U(0, h) = K_{\text{final}} + U(R, 0)$

Case 1 $K_f = K_0 + mgh$

Case 2 $K_f = K_0 + \frac{1}{2}kh^2 - \frac{1}{2}kR^2 = K_0 + \frac{1}{2}k(h^2 - R^2)$

only possible if $K_f \geq 0$. Not always "accessible".

Case 3 $K_f = K_0 + k(xy)_{(0, h)} - k(xy)_{(R, 0)} = K_0$

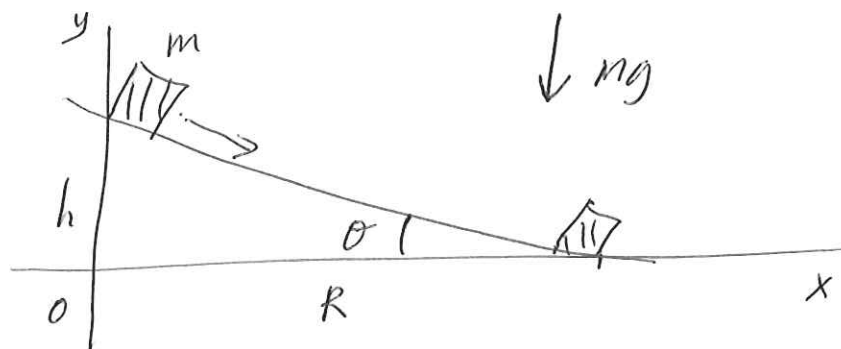
same "height", no change

Constrained Motion

W15

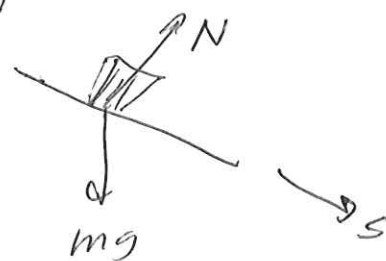
2D - Constrained motion in one-dimensional system

Ex (1)



- Mass slides down incline, no friction
- Mass stays on incline, $N > 0$
- N does no work

\Rightarrow potential energy only
from $m\vec{g}$



• If $v(t=0) = 0$ find v at bottom.

Note v along incline

• only single dimension, s , along incline.

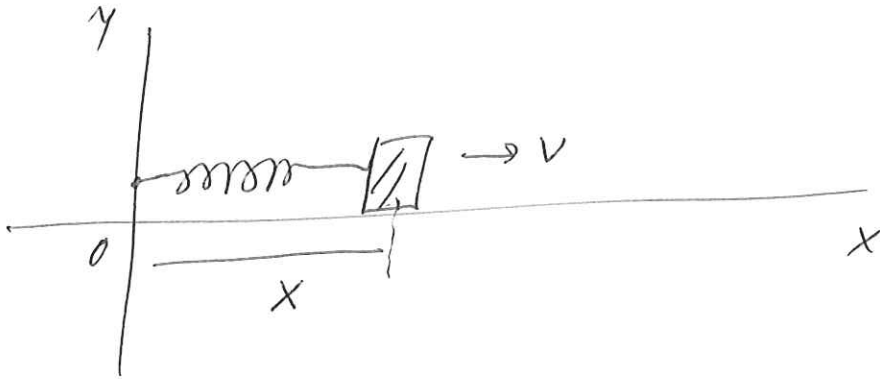
• First, $N = mg \cos \theta$

• $K + U_g = \text{const} \Rightarrow mgh = K_{\text{final}} = \frac{1}{2} m v_f^2$

$$\Rightarrow \boxed{v_f^2 = 2gh}$$

• Compare this to Taylor Problem 4.8, where N can go to zero.

2D constrained motion, another example w/6



• Spring force $F = -kx$

• Other forces

A free-body diagram of the mass. It shows a rectangular mass on a horizontal surface. Four force vectors are drawn: a normal force 'N' pointing vertically upwards, a gravitational force 'mg' pointing vertically downwards, a spring force '-kx' pointing horizontally to the left, and a friction force pointing horizontally to the right.

N and mg do no work
 \Rightarrow only spring force

• $F = -kx$. Is F conservative?

Yes, if it's one-dimensional and depends only on that dimension (but not on $v(t)$ or t)

• If conservative, then $F(x) = -dU/dx$

$$\Rightarrow U = -\int F dx = +\int kx dx = \frac{1}{2} kx^2 + C$$

$$\text{let } U = \frac{1}{2} kx^2$$

• Find amount of extension if $v(0) = v_0$

$$K + U = \text{const} \Rightarrow \frac{1}{2} m v_0^2 = \frac{1}{2} k x_f^2$$

$$\Rightarrow x_f^2 = m v_0^2 / k$$

Solution of $x(t)$ in 1-D problems

w/107

for conservative F

$$m \ddot{x} = F(x)$$

$$\vec{\nabla}_x F(x) \vec{x} = 0 \Rightarrow F = -dU/dx$$

\Rightarrow can find $U(x)$

Energy is a const of motion

$$\Rightarrow \mathcal{E} = \frac{1}{2} m \dot{x}^2 + U(x) = \text{const}$$

$$\text{i.e., } \frac{1}{2} m \dot{x}^2 + U(x) = \underbrace{\frac{1}{2} m \dot{x}^2(0) + U(x(0))}_{\text{call this } \mathcal{E} \text{ from i.o.c.}}$$

$$\Rightarrow \dot{x}^2 = \frac{2}{m} [\mathcal{E} - U(x)]$$

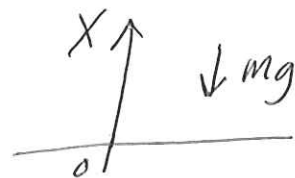
↑ ↑
known

$$\Rightarrow \frac{dx}{dt} = \pm \sqrt{\frac{2}{m} [\mathcal{E} - U(x)]}^{1/2}$$

Use
separation
variable

$$\frac{dx}{[\mathcal{E} - U(x)]^{1/2}} = \pm \sqrt{\frac{2}{m}} dt$$

2 signs



let $U = mgx$

$$v(t=0) = 0, \quad x(0) = 0$$

$$\Rightarrow \mathcal{E} = \frac{1}{2} mv^2 \Big|_0 + mgx \Big|_0 = 0$$

$$\int_0^x \frac{dx}{(-mgx)^{1/2}} = -\sqrt{\frac{2}{m}} \int_0^t dt$$

$dx/dt|_0 < 0$

A vertical coordinate system with the origin at the bottom. An upward-pointing arrow is labeled 'x' and a downward-pointing arrow is labeled 'v'.

↳ so pick negative sign

$$-2(-x)^{1/2} = -\sqrt{2g} t$$

$$(-x)^{1/2} = \sqrt{\frac{g}{2}} t$$

$$-x = \frac{g}{2} t^2$$

$$x = -\frac{1}{2} g t^2$$

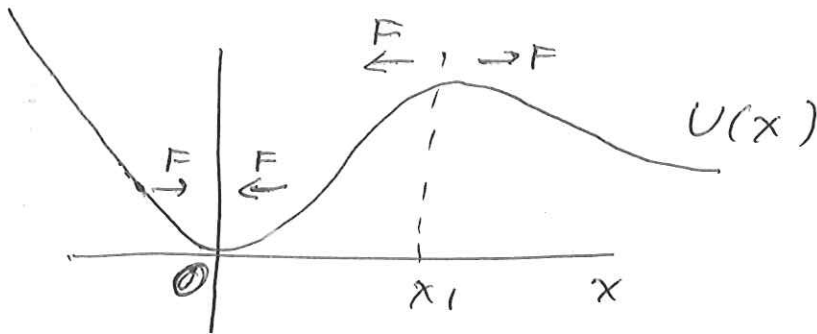
One-Dim potential

W19

$$m\ddot{x} = F(x)$$

$$m\ddot{x} = -dU/dx$$

$$\Sigma = \frac{1}{2}mv^2 + U(x) = \text{const}$$



Equil

$$dU/dx = 0 \Rightarrow x=0, x=x_1$$

stable / unstable

$$U'' > 0 \Rightarrow \text{stable}, U'' < 0 \Rightarrow \text{unst}$$

Small osc

$$x = x_0 + \tilde{x} \quad // 0$$

$$U = U(x_0 + \tilde{x}) = U_0 + U_0' \tilde{x} + U_0'' \frac{\tilde{x}^2}{2}$$

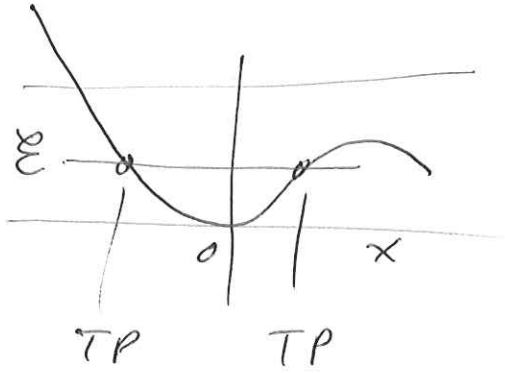
$$\Rightarrow U' = U_0'' \tilde{x}$$

$$\boxed{m\ddot{\tilde{x}} = -U_0'' \tilde{x}} \quad H=0.$$

$$U_0'' > 0 \text{ stable} \quad U_0'' < 0 \Rightarrow e^{\alpha t} \text{ UNST}$$

Turning points

$\mathcal{E} = \text{const} \Rightarrow$ ball on hill



$$\frac{1}{2} m \dot{x}^2 = \mathcal{E} - U(x)$$

$$\Leftrightarrow \mathcal{E} > U(x)$$

TP = turning points