

## Motion in harmonic oscillator (HO) central force

Consider a single mass  $m$  in a central force field  $\mathbf{F} = -k\mathbf{r}$ , with respect to a fixed origin. Here  $\mathbf{r}(t)$  is the radius vector,  $\mathbf{r} = (x,y,z)$ , in Cartesian. In this problem, we want to understand the various orbits of the mass about the HO central force.

1. Check that  $\mathbf{F}$  is conservative.
2. Find the potential  $U(\mathbf{r})$ .
3. It follows that the energy is conserved. Write down the expression for this. Note that  $\mathbf{v}(t) = d\mathbf{r}/dt$ .
4.  $\mathbf{F}$  is a central force (in addition to being conservative). Therefore, show that the angular momentum  $\mathbf{L}$  is a constant, where  $\mathbf{L} = m \mathbf{r} \times \mathbf{v}$ . Your proof must use 2 important vector identities.
5. If  $\mathbf{L}$  is constant, prove that  $\mathbf{r}(t)$  must stay on a 2-dimensional plane. State the important vector identity needed to prove this. Thus, assume that  $\mathbf{r} = x(t)\mathbf{x}^\wedge + y(t)\mathbf{y}^\wedge$ .
6. Polar coordinates are clearly preferable. Refer to Taylor Sec 1.7 and check the following (you do not have to derive these – just collect and review them):  $\mathbf{r}(t) = r(t)\mathbf{r}^\wedge(t)$ ,  $\mathbf{v} = \mathbf{r}^\wedge dr/dt + \phi^\wedge r d\phi/dt$ . Also, obtain the expression for  $d^2\mathbf{r}/dt^2$ .
7. So, write down the 2 polar components of  $m\mathbf{a} = \mathbf{F}$ . From the  $\phi^\wedge$  component, show that the angular momentum magnitude,  $L$ , is also a constant of the motion (so far we only used the direction of  $\mathbf{L}$ ).
8. Eliminate in the radial equation all the  $\phi$  terms by using the angular momentum equation from 7. Thus, find a 2<sup>nd</sup> order nonlinear ODE for  $r(t)$ . This equation cannot be solved in an easy way.
9. But, we haven't used energy conservation yet. Write down the energy constant using polar coordinates (use 6 above). Again eliminate the  $\phi$  terms using the constancy of  $L$ , from 7. Check that this energy constant contains terms only in  $dr/dt$  and  $r$ .
10. As a check on energy conservation, start with your energy constant in 9 and differentiate  $E$  in  $t$ . This should be zero, ie,  $dE/dt = 0$ . Show that this results in the 2<sup>nd</sup> order ODE for  $r(t)$  in 8 above.
11. Take a square root in the equation from 9 and thus obtain a 1<sup>st</sup> order ODE for  $r(t)$ . (there will be a +/- sign from the square root). Check that this equation is separable and therefore solvable in principle.

12. Define an effective potential starting from your energy equation in 9. Let  $U_{\text{eff}}(r) = U(r)$  (of the central force) + angular momentum term in  $r$ . Make a sketch of  $U_{\text{eff}}(r)$  vs  $r$ . Draw a horizontal line for  $E$  on this sketch. Since  $(m/2)(dr/dt)^2 > 0$ , observe that we must have  $E > U_{\text{eff}}(r)$  for a solution to exist. Thus, show that  $r(t)$  must be bounded from above and below according to  $r_{\text{min}} < r(t) < r_{\text{max}}$ .
13. If  $U_{\text{eff}}(r)$  is at its minimum point, and  $E$  is adjusted so that  $E = U_{\text{min}}$ , observe that  $r(t)$  must be a constant, corresponding to circular orbit. Find  $r_0$ , the radius corresponding to  $U_{\text{min}}$ . Insert this value of  $r$  in the angular momentum equation for  $d\phi/dt$  found in 7 above to show that the angular frequency of the circular orbit,  $\omega_0$ , is a constant,  $\omega_0$ , and given in terms of  $L$ ,  $m$ , and  $r_0$ .
14. Using your knowledge of centripetal force, applied to the  $-kr$  force, find the frequency of circular orbit. Compare this to  $\omega_0$  found above. You will have to eliminate  $L$  from your expression using the angular momentum equation found in 7.
15. We now want to perform small oscillations of the mass *about the circular orbit*. To do this, start from the 2<sup>nd</sup> order ODE for  $r(t)$  found in 8. Let  $r(t) = r_0 + s(t)$ . Substitute this into the ODE and Taylor expand the RHS for small  $s$ . Use the definition of  $r_0$  to rewrite the RHS, so its only proportional to  $s$ . Thus, obtain a 2<sup>nd</sup> order ODE for  $s(t)$ . Solve this in general and find  $\omega$ , the frequency of small oscillations about  $r_0$ . Compare this frequency with  $\omega_0$ . Can you sketch the circular orbit + small oscillations added on?