

HWK #4 Solutions

5.8

Eqn for A_ϕ

$$\nabla^2 A_\phi - \frac{1}{r^2 \sin^2 \theta} A_\phi = -\mu_0 \bar{J}_\phi(r, \theta)$$

⇒ calculate the response to a unit current

$$\mathbf{j} = \frac{\delta(r-r')}{r} \delta(\theta-\theta')$$

$$\int_{-a}^a \int_0^\pi dr r d\theta j = 1$$

$$\nabla^2 G_{\ell} - \frac{1}{r^2 \sin^2 \theta} G_{\ell} = -\mu_0 \mathbf{j}$$

$$G_{\ell} = \sum_{\ell} P_{\ell}(\cos \theta) R_{\ell}(r)$$

$$\sum_{\ell} \frac{1}{r^2} \left(\frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} R_{\ell} - \ell(\ell+1) R_{\ell} \right) P_{\ell}' = -\mu_0 \mathbf{j}$$

⇒ eliminate the sum over ℓ by multiplying by $P_{\ell}'(\cos \theta)$ and integrating over $d \cos \theta$

$$\frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} R_{\ell} - \ell(\ell+1) R_{\ell} = -\mu_0 r \int_{-1}^1 d \cos \theta \frac{\delta(r-r')}{r} \delta(\theta-\theta')$$

$$\int_{-1}^1 d \cos \theta (P_{\ell}')^2 = \frac{2}{2\ell+1} \frac{(\ell+1)!}{(\ell-1)!} \quad \textcircled{\times} P_{\ell}'(\cos \theta) \frac{2\ell+1}{2} \frac{1}{\ell(\ell+1)}$$

$$= -\mu_0 r \delta(r-r') P_{\ell}'(\cos \theta') \sin \theta$$

$$\textcircled{\times} \frac{2\ell+1}{2} \frac{1}{\ell(\ell+1)}$$

Let $S_l = -\mu_0 \frac{r^l}{2} P_l'(\cos\theta') \sin\theta' \frac{2l+1}{2l(l+1)}$

~~$\frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} R_e - l(l+1) R_e = S_l \delta(r-r')$~~

$r \neq r'$

$R_e = \left(\frac{r}{r'}\right)^l c_l \quad r < r'$
 $= \left(\frac{r'}{r}\right)^{l+1} c_l \quad r > r'$
 \Rightarrow continuity at $r=r'$

$r \approx r'$

$\frac{\partial^2}{\partial r^2} R_e \approx \frac{1}{r^2} S_l \delta(r-r')$

$\frac{\partial R_e}{\partial r} \Big|_{r'-\epsilon}^{r'+\epsilon} = \frac{1}{r'^2} S_l$

$(-2l-1) c_l \frac{1}{r'} = \frac{1}{r'^2} S_l$

$c_l = -\frac{S_l}{2l+1} \frac{1}{r'}$

$G_{ll} = \frac{\mu_0}{2} \sin\theta' \sum_l \frac{1}{l(l+1)} P_l'(\cos\theta') P_l'(\cos\theta)$

$\otimes \frac{r^l}{r'^{l+1}} r'$

interior

$A_{cl} = \int_0^\pi \int_{-\pi}^\pi \sin\theta' d\theta' G_{ll} J_{cl}(r', \theta')$

$$A_{\ell} = \frac{\mu_0}{2} \frac{1}{r} \int \underbrace{dr' d\cos\theta' r'^2}_{\frac{1}{2\pi} \sum_n dx'_n} \frac{1}{r'^{\ell+1}} P_{\ell}'(\cos\theta') P_{\ell}'(\cos\theta) J_{\ell}$$

$$A_{\ell} = -\frac{\mu_0}{4\pi} \sum_n m_e r^{\ell} P_{\ell}'(\cos\theta)$$

$$m_e = -\int dx'_n \frac{1}{r'^{\ell+1}} P_{\ell}'(\cos\theta') \frac{1}{r'^{\ell+1}} J_{\ell}(r', \theta')$$

outside

$$A_{\ell} = -\frac{\mu_0}{4\pi} \frac{M_{\ell}}{r^{\ell+1}} P_{\ell}'(\cos\theta)$$

$$M_{\ell} = -\frac{1}{r^{\ell+1}} \int dx'_n r'^{\ell} P_{\ell}'(\cos\theta') J_{\ell}(r', \theta')$$

(4)

(5.10)

$$a) \quad \nabla^2 A_{\alpha} - \frac{1}{e^2} A_{\alpha} = -\mu_0 J_{\alpha}(e, z) \quad (1)$$

$$J_{\alpha} = I \delta(e-a) \delta(z)$$

$$\text{note } \int_0^{\infty} \int_{-\infty}^{\infty} dz de J_{\alpha} = I$$

$$\text{Let } A_{\alpha} = \int_0^{\infty} dk \cos kz R_k(e)$$

\Rightarrow even in z and oscillatory

\Rightarrow non-oscillatory in e

\Rightarrow continuous since ∞ system

insert into Eq (1) \Rightarrow

$$\int_0^{\infty} dk \cos kz \left[\frac{1}{e} \frac{\partial}{\partial e} e \frac{\partial}{\partial e} - k^2 - \frac{1}{e^2} \right] R_k = -\mu_0 J_{\alpha}$$

\Rightarrow eliminate k integral

\Rightarrow operate with $\int dz \cos k'z$

$$\int dz \cos(k'z) \cos kz = a \int \frac{dz}{4a} \left(e^{ik'z} + e^{-ik'z} \right) \left(e^{ikz} + e^{-ikz} \right)$$

$$= \frac{\pi}{2} \left[2 \delta(k-k') + 2 \delta(k+k') \right] = \pi \delta(k-k')$$

$$\pi \left[\frac{1}{e} \frac{\partial}{\partial e} e \frac{\partial}{\partial e} - k^2 - \frac{1}{e^2} \right] R_k = \mu_0 I \delta(e-a)$$

$$\left[e \frac{\partial}{\partial e} e \frac{\partial}{\partial e} - k^2 e^2 - 1 \right] R_k = -\frac{\mu_0}{\pi} a^2 I \delta(e-a)$$

(5)

$$\underline{e \neq a}$$

$$\left(e \frac{\partial}{\partial e} e \frac{\partial}{\partial e} - k^2 e^2 - 1 \right) R_k = 0$$

\Rightarrow modified Bessel Eqn of order 1.

$$R_k \sim I_1(ke), K_1(ke)$$

$$R_k = c_k I_1(ke_<) K_1(ke_>)$$

$\Rightarrow I_1$ bounded as $ke \rightarrow 0$

K_1 bounded as $ke \rightarrow \infty$

$e_< = \text{smaller of } e, a$

$e_> = \text{larger of } e, a$

$$\underline{e = a}$$

$$\frac{\partial^2}{\partial e^2} R_k = -\frac{\mu_0}{u} \delta(e-a) I$$

$$\frac{\partial}{\partial e} R_k \Big|_{e=a} = -\frac{\mu_0}{u} I$$

$$k c_k [K_1'(ka) I_1(ka) - I_1'(ka) K_1(ka)] = -\frac{\mu_0}{u} I$$

$$c_k = -\frac{\mu_0}{u} I a \frac{1}{W(I_1, K_1)} \frac{1}{ka}$$

$$= \frac{\mu_0}{u} I a$$

$$A_\phi = \frac{\mu_0 I a}{u} \int_0^\infty dk \cos kz I_1(ke_<) K_1(ke_>)$$

$$c) \vec{B} = \nabla \times (A_{\phi} \hat{e})$$

$$\hat{e} = e \frac{\hat{e}}{e} = e \nabla \phi$$

$$\vec{B} = \nabla \times (A_{\phi} e \nabla \phi) = \nabla (e A_{\phi}) \times \nabla \phi$$

$$B_z = \frac{1}{e} \frac{\partial}{\partial e} (e A_{\phi})$$

$$B_{\phi} = 0$$

$$B_{\rho} = -\frac{1}{e} \frac{\partial}{\partial z} e A_{\phi} = -\frac{\partial}{\partial z} A_{\phi}$$

Limit $z=0$ and $\rho = a - \epsilon$ with ϵ small

\Rightarrow should find $B_z \sim \frac{1}{\epsilon} \Rightarrow$ like infinite wire

$$B_{\rho} = 0$$

Since $B_{\rho} = -\frac{\partial}{\partial z} A_{\phi} \sim \sin(kz) \rightarrow 0$

$$\Rightarrow B_{\rho} = 0$$

$$B_z = \frac{1}{e} \frac{\partial}{\partial e} (e A_{\phi}) = \frac{\mu_0 I a}{\pi} \int_0^{\infty} dk \frac{1}{e} \frac{\partial}{\partial e} [e I_1(k e)] k_1(k a)$$

but $\frac{d}{dx} [x I_1(x)] = x I_0(x)$

$$\frac{d}{dx} (x I_1) = x I_0 = k e I_0$$

$$= \frac{\mu_0 I a}{\pi} \int_0^{\infty} dk k I_0(k e) k_1(k a)$$

Note that the integral diverges as $\rho \rightarrow a$ because for large k ,

$$I_0 \sim \frac{e^{ke}}{(ke)^{1/2}} \text{ and } K_1 \sim \frac{e^{-ka}}{(ka)^{1/2}}$$

$$\text{so } k I_0(k\rho) K_1 \sim e^{-k(a-\rho)}$$

\Rightarrow integrand dominated by large k

\Rightarrow expand I_0, K_1 for large ~~the~~ argument

$$I_0 = \frac{e^x}{\sqrt{2\pi x}}, \quad K_1 = \sqrt{\frac{\pi}{2x}} e^{-x}$$

$$B_z = \frac{\mu_0 I a}{\pi} \int_0^\infty dk k \frac{e^{ke}}{\sqrt{2\pi ke}} \frac{\sqrt{\pi}}{\sqrt{2ka}} e^{-ka}$$

$$= \frac{\mu_0 I a}{2\pi a} \int_0^\infty dk e^{-k(a-\rho)}$$

$$= \frac{\mu_0 I}{2\pi} \frac{1}{a-\rho} \quad \Rightarrow \text{like infinite wire}$$