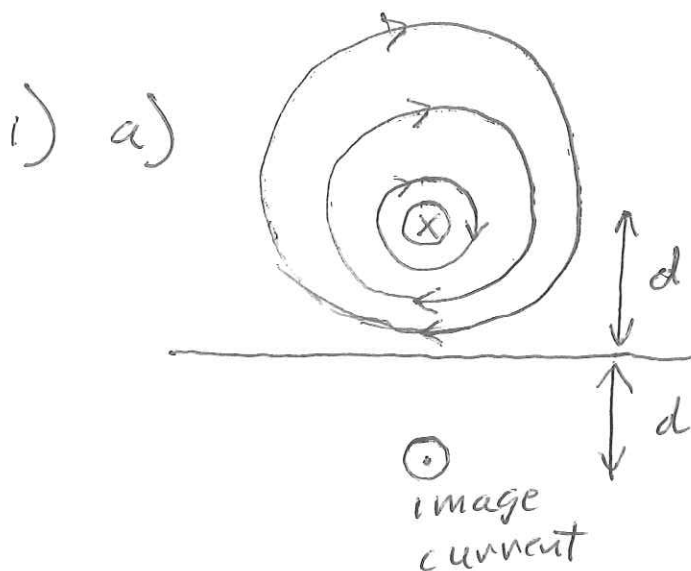


①

Physics 606 Final Exam Solutions



But I due to
image current

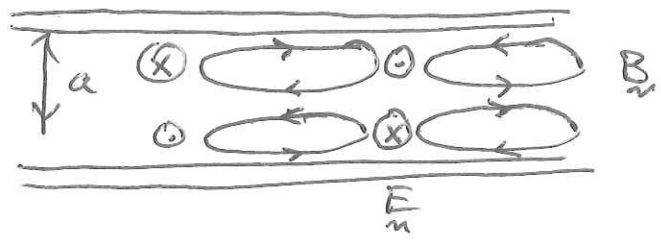
$$B = \frac{\mu_0 I}{2\pi(2d)}$$

$$\text{force/length} = IB = \lambda g$$

$$\boxed{\frac{\mu_0 I^2}{4\pi d} = \lambda g}$$

TE mode $\Rightarrow E_z = 0, B_z \neq 0$

b)



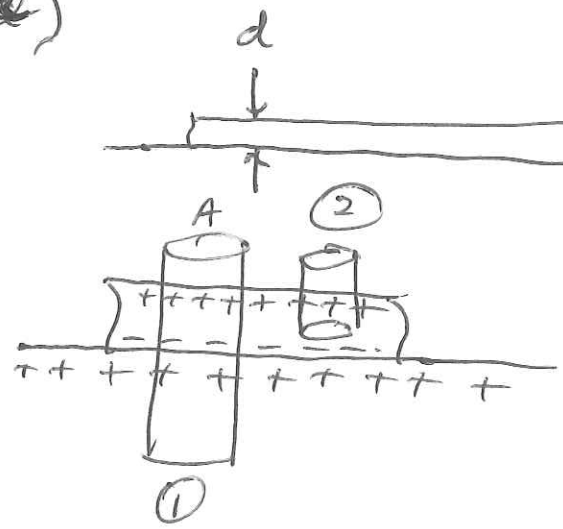
$$\left(\nabla^2 + \frac{\omega^2}{c^2}\right) B_z = 0$$

$$\nabla^2 = \nabla_t^2 - k_z^2 \approx -\frac{1}{a^2} - k_z^2 = -\frac{\omega^2}{c^2}$$

$$\Rightarrow k_z^2 = \frac{\omega^2}{c^2} - \frac{1}{a^2} > 0$$

$$\Rightarrow \omega^2 > \frac{c^2}{a^2}$$

c)



dielectric is polarized
 E_0 outside dielectric use
 Gaussian surface ①

$$2AE_0 = \frac{\sigma}{\epsilon_0} A \quad E_0 = \frac{\sigma}{2\epsilon_0}$$

Use Gaussian surface ② to calculate
 E_i inside dielectric. Since no free charge

$$D_0 = D_i = \epsilon_0 E_0 = \epsilon_i E_i$$

$$\Rightarrow E_i = \frac{\epsilon_0}{\epsilon_i} E_0 < E_0$$

But $\nabla \cdot \vec{E} = \frac{\rho_{tot}}{\epsilon_0}$ For ② $\rho_{tot} = \rho_p = \text{pol charge}$

$$E_0 - E_i = \frac{1}{\epsilon_0} \sigma_{pol} = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{\epsilon_0}{\epsilon_i}\right)$$

$$\sigma_{\text{pol}} = \frac{\epsilon}{2} \left(1 - \frac{\epsilon_0}{\epsilon_i}\right) \Rightarrow \text{top of dielectric} \quad (3)$$

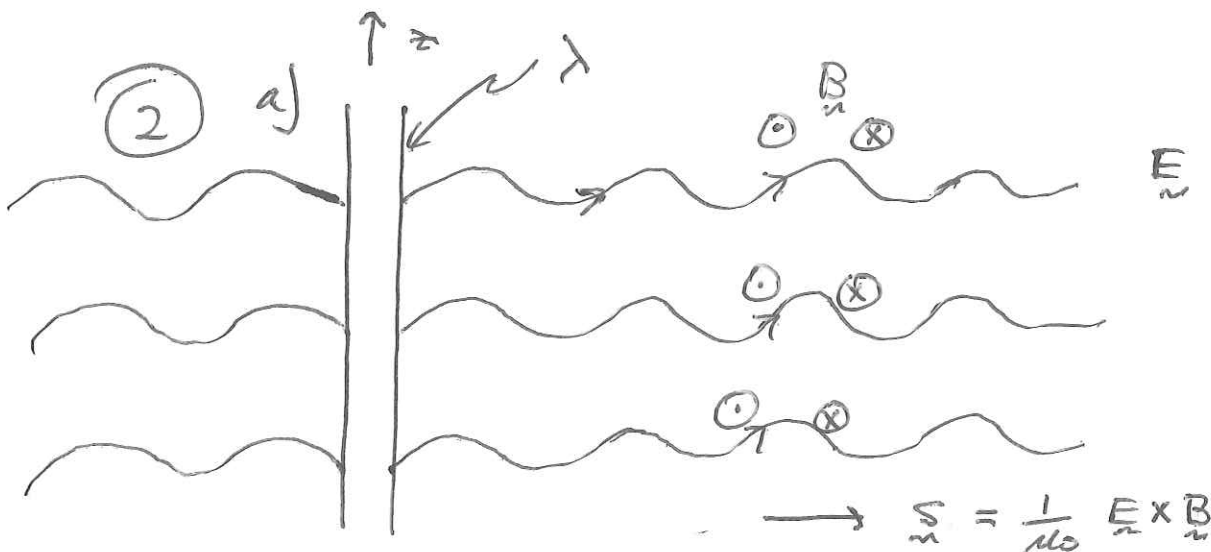
$\Rightarrow -\sigma_{\text{pol}} \Rightarrow \text{bottom of dielectric}$

Calculate force on σ_{pol} and $-\sigma_{\text{pol}}$ due to σ .

~~F~~

\Rightarrow E due to σ is uniform
so forces on σ_{pol} and $-\sigma_{\text{pol}}$
cancel

\Rightarrow no net force



\vec{E} is along z

\vec{B} is in ϕ direction

\vec{S} is radial

b) \Rightarrow electromagnetic wave

$$E_z = -c B_\phi$$

total power from radiation field is
independent of ϵ

$$\Rightarrow 2\pi r S \sim \epsilon E^2 \sim \text{const}$$

$$E, B \sim \frac{1}{r^{1/2}}$$

(4)

$$c) \quad \nabla \times \underline{B} = \mu_0 \underline{J} + \epsilon_0 \frac{\partial}{\partial t} \underline{E} \quad \mu_0$$

$$\underline{B} = \nabla \times \underline{A} \quad \underline{J} = J_z \hat{z}$$

$$\underline{E} = -\frac{\partial}{\partial t} \underline{A} - \nabla \phi \Rightarrow E_z = -\frac{\partial}{\partial t} A_z$$

$$\left[\nabla \times (\nabla \times \underline{A}) \right]_z = \left[\nabla \nabla \cdot \underline{A} - \nabla^2 \underline{A} \right]_z$$

$$= \frac{\partial}{\partial z} \frac{\partial}{\partial z} A_z - \nabla^2 A_z$$

$$\frac{\partial}{\partial z} = 0$$

$$-\nabla^2 A_z = \mu_0 J_z - \cancel{\mu_0} \ddot{A}_z \frac{1}{c^2}$$

$$\nabla^2 A_z + \frac{\omega^2}{c^2} A_z = \mu_0 J_z$$

$$J_z = \underline{J} \cdot \underline{v}_z = -i\omega d \lambda \delta(r-a) \frac{1}{2\pi a}$$

$$\frac{1}{x} \int dx \delta(x) = \lambda = \frac{1}{l} \int_0^{2\pi} d\phi \int_0^a \rho \rho$$

$$\Rightarrow \rho = \frac{\lambda}{2\pi a} \delta(r-a)$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} A_z \right) + \frac{\omega^2}{c^2} A_z = \mu_0 \left(\frac{-i\omega d \lambda}{2\pi a} \delta(r-a) \right)$$

$$\boxed{A_z'' + \frac{1}{r} A_z' + \frac{\omega^2}{c^2} A_z = S_0 \delta(r-a)}$$

$$\Rightarrow \text{Bessel Eqn with } \nu=0, k=\frac{\omega}{c}$$

$$S_0 = \frac{-i\omega d \lambda \mu_0}{2\pi a}$$

d) $\rho > a$

$$A_z'' + \frac{1}{\rho} A_z' + \frac{\omega^2}{c^2} A_z = 0$$

B.C. \Rightarrow outgoing wave

$$A_z \sim H_0^{(1)} \sim \frac{e^{ike}}{e^{i1/2}} e^{-i\omega t} \Rightarrow \text{outgoing}$$

~~$A_z = \dots$~~ $A_z = C^+ H_0^{(1)}\left(\frac{\omega}{c} \rho\right)$

$\rho < a$

\Rightarrow bounded at $\rho=0$

$\Rightarrow J_0(k\rho)$

$$A_z = C^- J_0\left(\frac{\omega}{c} \rho\right)$$

Integration eqn for A_z across $\rho=a$

$$\frac{\Delta A_z}{\Delta \rho} \Big|_{a-\epsilon}^{a+\epsilon} = S_0$$

$$A_z \Big|_{a-\epsilon}^{a+\epsilon} = 0 \Rightarrow C^+ = C^0 J_0\left(\frac{\omega}{c} a\right)$$
$$C^- = H_0^{(1)}\left(\frac{\omega}{c} a\right) C^0$$

$$C^0 \frac{\omega}{c} \left[H_0^{(1)'}\left(\frac{\omega}{c} a\right) J_0\left(\frac{\omega}{c} a\right) - H_0^{(1)}\left(\frac{\omega}{c} a\right) J_0'\left(\frac{\omega}{c} a\right) \right]$$
$$= S_0$$

A

$\underline{r > a}$ $A_z = c^0 J_0\left(\frac{\omega}{c}a\right) H_0^{(1)}\left(\frac{\omega}{c}r\right)$

$\underline{r < a}$ $A_z = c^0 J_0\left(\frac{\omega}{c}r\right) H_0^{(1)}\left(\frac{\omega}{c}a\right)$

$$c^0 = \frac{-i \omega d \lambda \mu_0 c}{2\pi a \omega} \left[\underbrace{H_0^{(1)'} J_0 - H_0^{(1)} J_0'}_{i(N_0' J_0 - N_0 J_0')} \right]_{r=a}^{-1}$$

$$i \frac{2}{\pi x} = i \frac{2c}{\pi \omega a}$$

e) $c^0 = -\frac{\lambda d \mu_0 \omega}{4}$
 $\vec{B} = \nabla \times \vec{A}$

$B_\phi = -\frac{\partial}{\partial r} A_z$

$A_z \sim c^0 J_0\left(\frac{\omega}{c}a\right) \sqrt{\frac{2}{\pi \frac{\omega}{c}e}} e^{i\frac{\omega}{c}e - i\frac{\pi}{4}}$

$B_\phi = -i \frac{\omega}{c} c^0 J_0\left(\frac{\omega}{c}a\right) \sqrt{\frac{2}{\pi \frac{\omega}{c}e}} e^{i\frac{\omega}{c}e - i\frac{\pi}{4}}$

$E_z = -c B_\phi$

f) $S_e = -\frac{1}{2} \frac{1}{\mu_0} \text{Re}(E_z B_\phi^*)$

$= \frac{1}{2\mu_0 c} \frac{\omega^2}{c^2} |c^0|^2 J_0^2\left(\frac{\omega}{c}a\right) \frac{2}{\pi \frac{\omega}{c}e}$

total power

$P \sim S_e 2\pi e \Rightarrow \text{indep of } e$

3

$$a) \quad \vec{J}'(x') = K \delta(x') \hat{y}$$

$$\nabla \times \vec{B}' = \mu_0 K \delta(x') \hat{y}$$

$$-\frac{\partial}{\partial x} B_z' = \mu_0 K \delta(x')$$

$$\boxed{B_z' = -\frac{\mu_0 K}{2} H(x')}$$

define

$$H = 1 \quad x' > 0$$

$$= -1 \quad x' < 0$$

$$e' = 0 \quad \text{and} \quad \frac{\partial}{\partial t} = 0 \quad \Rightarrow \quad \boxed{\vec{E}' = 0}$$

b) Lab frame S \Rightarrow (ce, \vec{J}) are a 4-vector

$$ce = \gamma (ce' + \beta \vec{J} x')$$

$$\cancel{J_x} = \cancel{J_x'} \quad J_y = J_y'$$

$$\boxed{J_y = K \delta(\gamma(x - \beta ct))}$$

 ~~$x = \beta ct$~~

$$x' = \gamma(x - \beta ct)$$

$$\boxed{e = 0}$$

$$c) \quad \frac{1}{c} E_y = \gamma \left[\cancel{\frac{1}{c} E_y'} - \beta \times B_z' \right]_y$$

$$= \gamma \beta B_z'$$

$$\boxed{E_y = -\frac{c \gamma \beta \mu_0 K}{2} H(\gamma(x - \beta ct))}$$

$$B_z = \gamma \left(B_z' + \beta \times \frac{1}{c} E_{y'} \right)$$

$B_z = -\gamma \frac{\mu_0 K}{2} H(\gamma(x - \beta ct))$	$E_y = v B_z$
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d) $\nabla \times \underline{B} = \mu_0 \underline{J} + \frac{1}{c^2} \frac{\partial \underline{E}}{\partial t}$

$$-\frac{\partial}{\partial x} B_z = \mu_0 J_y + \frac{1}{c^2} \frac{\partial}{\partial t} E_y$$

$$\nabla \times \underline{E} + \frac{\partial \underline{B}}{\partial t} = 0$$

$$\frac{\partial}{\partial x} E_y + \frac{\partial}{\partial t} B_z = 0$$

$$J_y = J_x (x - vt)$$

$$\Rightarrow \frac{\partial}{\partial t} = -v \frac{\partial}{\partial x}$$

$$\frac{\partial}{\partial x} E_y - v \frac{\partial}{\partial x} B_z = 0$$

$E_y = v B_z$

$B_z = -\frac{\mu_0 \gamma K}{2} H(\gamma(x - vt))$

$$-\frac{\partial}{\partial x} B_z = \mu_0 J_y + \frac{1}{c^2} \left(-v \frac{\partial}{\partial x} \right) v B_z$$

$$\frac{\partial}{\partial x} B_z = -\mu_0 \gamma^2 J_y$$

$$\Delta B_z = -\mu_0 \gamma^2 K \int dx H(\gamma(x - vt))$$