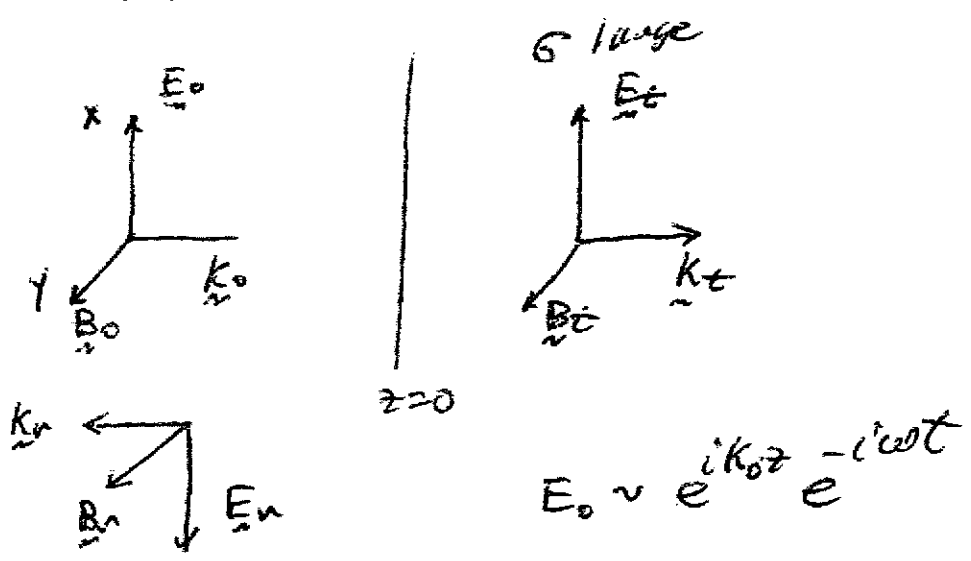


Phys. 606

Exam 2 Solutions

1

a)



b) calculate dispersion relation for conductor

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \approx \mu_0 \sigma \vec{E}$$

$$\nabla \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$$

$$\frac{\partial}{\partial x} = 0$$

$$\nabla \times \vec{B} = -\nabla \times (\nabla \times \vec{E}) = -\nabla(\nabla \cdot \vec{E}) + \nabla^2 \vec{E} = \mu_0 \sigma \vec{E}$$

~~$$k^2 = \mu_0 \sigma \omega$$~~

$$+k^2 = \mu_0 \sigma (i\omega)$$

$$k = \sqrt{\frac{\mu_0 \sigma \omega}{2}} (1+i) \equiv \frac{1}{\delta} (1+i)$$

vacuum

$$\nabla \times \vec{B} = \nabla^2 \vec{E} = \mu_0 \epsilon_0 \omega^2 \vec{E}$$

$$k^2 = \mu_0 \epsilon_0 \omega^2$$

(2)

c) Since $\mu = \mu_0$ on both sides

$$\hat{n} \times \underline{B} \Big| = 0$$

$$\hat{n} \times \underline{E} \Big| = 0$$

d) ~~no~~
continuity of E_x

$$\textcircled{1} E_0 - E_r = E_t$$

continuity of B_y

$$\textcircled{2} B_0 + B_r = B_t$$

relate B_y to E_y

in vacuum: $\nabla \times \underline{E} + \frac{\partial \underline{B}}{\partial t} = 0$

$$ik E_x - i\omega B_y = 0$$

$$E_x = c B_y$$

in conductor:

$$\nabla \times \underline{E} = -\omega \underline{B}$$

from $\textcircled{1}$ and $\textcircled{2}$

$$\cancel{k} B_0 - \cancel{k} B_r = \frac{\omega B_t \epsilon_0}{(1+i)c} = \frac{k_0 \epsilon_0 B_t}{1+i}$$

$$B_0 + B_r = B_t$$

$$2 B_0 = B_t \left(1 + \frac{k_0 \epsilon_0}{1+i} \right)$$

$$B_t = \frac{2B_0}{1 + \frac{k_0 \delta}{1+i}} \approx \boxed{2B_0}$$

$$\# E_t = \frac{\omega 2B_0 \delta c}{(1+i)c}$$

$$\boxed{E_t = \frac{2k_0 \delta c B_0}{1+i}}$$

$$\boxed{B_r = B_t - B_0 \approx B_0}$$

$$\boxed{E_r \approx c B_0}$$

2) Poynting flux

conductor:

$$S_z^t = \text{Re} \frac{E_x B_y^*}{2\mu_0} = \text{Re} \frac{\omega \delta}{1+i} \frac{(B_0)^2 (1-i)}{2\mu_0 (1-i)}$$

$$= \frac{\omega \delta}{\mu_0 c} B_0^2 c = k_0 \delta \frac{B_0^2 c}{\mu_0}$$

vacuum:

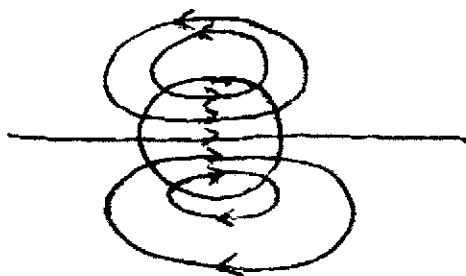
$$S_z^0 = \frac{1}{2\mu_0} c B_0^2$$

$$\boxed{\frac{S_z^t}{S_z^0} = 2k_0 \delta}$$

Fraction transmitted

⇒ dissipated by Joule heating

(4)

(2) a) sketch \vec{B} b) Since $\nabla \times \vec{H} = 0$

$$\Rightarrow \vec{H} = -\nabla \varphi_m$$

$$\vec{B} = \mu_0 (\vec{H} + \vec{M}) \Rightarrow \vec{M} \text{ specified}$$

$$\nabla \cdot \vec{B} = 0 \Rightarrow \nabla^2 \varphi_m = \nabla \cdot \vec{M}$$

$$\nabla \cdot \vec{M} = \int_{\Sigma} M_0 H(a-r)$$

$$= -\frac{\partial r}{\partial x} M_0 \delta(a-r)$$

$$r^2 = x^2 + y^2$$

$$\frac{\partial r}{\partial x} = \frac{x}{r} \Rightarrow \frac{\partial r}{\partial x} = \cos \varphi$$

$$\nabla^2 \varphi_m = -M_0 \cos \varphi \delta(a-r)$$

$$\text{where } \nabla^2 = \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2}$$

Basis functions

$$\sim \sin(n\varphi), \cos(n\varphi), e^{\pm in}$$

$$n = 1, 2, \dots$$

Since $\nabla \cdot \mathbf{M} \sim$ has even symmetry in φ ,
 Q_m is also even in φ

For $e > a$,

$$Q_m^> = \sum_n b_n^> \left(\frac{a}{e}\right)^n \cos n\varphi$$

For $e < a$

$$Q_m^< = \sum_n b_n^< \left(\frac{e}{a}\right)^n \cos n\varphi$$

Only $n=1$ yields non-zero $b_n^> <$ since
 $\nabla \cdot \mathbf{M} \sim \cos \varphi$

$$\Rightarrow Q_m^> = b_n^> \left(\frac{a}{e}\right) \cos \varphi$$

$$Q_m^< = b_n^< \left(\frac{e}{a}\right) \cos \varphi$$

c) Boundary conditions:

$$B_e \Big|_{a-e}^{a+e} = 0$$

$$\underbrace{H_e}_{\frac{1}{a} \frac{\partial Q_m}{\partial \varphi}} \Big|_{a-e}^{a+e} = 0 \Rightarrow Q_m \Big|_{a-e}^{a+e} = 0$$

$$\frac{1}{a} \frac{\partial Q_m}{\partial \varphi}$$

⑥

d) see previous page for $Q_m^>$
matching \Rightarrow

$$Q_m^>(a) = Q_m^<(a) \Rightarrow b_>^1 = b_<^1 \equiv b^1$$

$$B_e^< = \mu_0 \left(-\frac{\partial Q_m^<}{\partial r} + M_0 \cos \varphi \right)$$

$$B_e^> = \mu_0 \left(-\frac{\partial Q_m^>}{\partial r} \right)$$

$$-b_<^1 \frac{1}{a} + M_0 = b_>^1 \frac{1}{a}$$

$$2b^1 = M_0 a \Rightarrow$$

$$b^1 = \frac{M_0 a}{2}$$

$$Q_m^> = \frac{M_0 a^2}{2r} \cos \varphi$$

$$Q_m^< = \frac{M_0}{2} r \cos \varphi = \frac{M_0}{2} x$$

e) $B_z = \mu_0 \left(H_z + M_z \right)$

$a > r$

$$B_x = \mu_0 \left(-\frac{M_0}{2} + M_0 \right) = +\frac{\mu_0 M_0}{2}$$

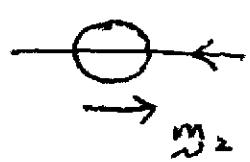
$$B_y = 0$$

$a < r$

$$B_e = -\mu_0 \frac{\partial Q_m^>}{\partial r} = \mu_0 \frac{M_0 a^2}{2r^2} \cos \varphi$$

$$B_e = -\frac{\mu_0}{r} \frac{\partial Q_m^>}{\partial \varphi} = \mu_0 \frac{M_0 a^2}{2r^2} \sin \varphi$$

(f)

(2)  $B_{1Q} \sim \frac{\mu_0 M_0 a^2}{2 r^2} \sim B_1$ from (1)

(7)

(1) $U_{21} = -m_2 \cdot B_1$
 $= m_2 B_{1Q} > 0$

$U_{21} = \frac{\text{energy}}{\text{length}}$ of (2) from (1)

$m_2 = M_0 \pi a^2 = \frac{\text{magnetic moment}}{\text{length}}$

$$F_{21} = - \frac{\partial U_{21}}{\partial r} = \frac{2 \mu_0 M_0 a^2}{2 r^3} M_0 \pi a^2$$
$$= \frac{2 \mu_0 \pi (M_0 a^2)^2}{2 r^3}$$

\Rightarrow repulsive