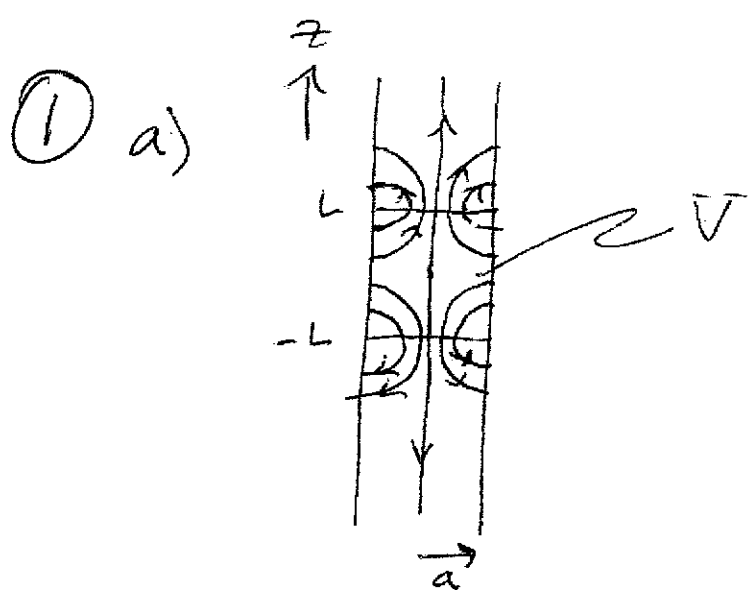
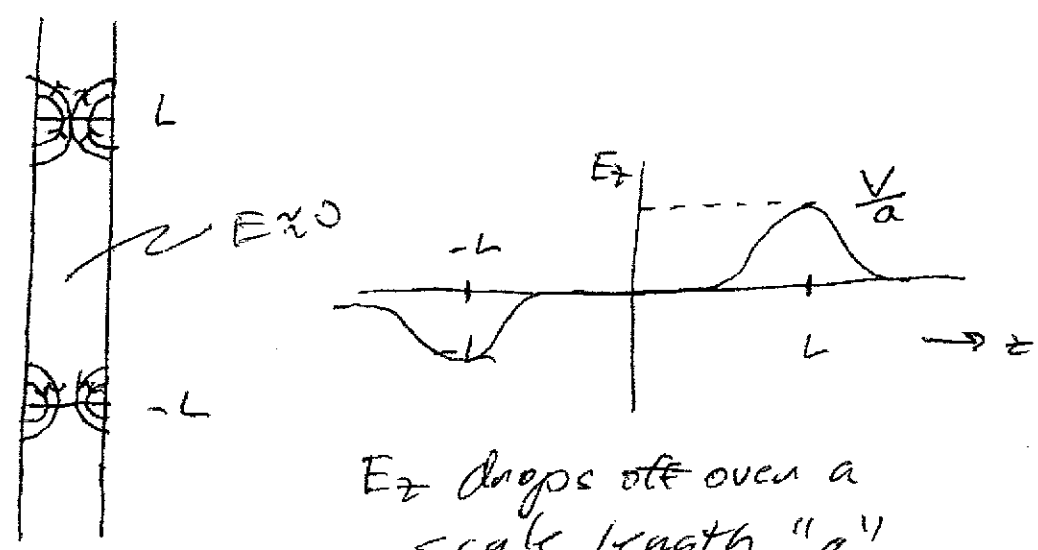


Physics 606 Midterm #1 Solutions



b) $a \ll L$



E_z drops off over a scale length "a" for $z > L$

$$E_z \sim \frac{V}{a} \quad \text{max value for } \ell = 0$$

2

c) Solve Laplace's eqn with $\frac{\partial}{\partial \rho} = 0$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} \Phi + \frac{\partial^2}{\partial z^2} \Phi = 0$$

\Rightarrow oscillatory in z so can match $\Phi(\rho=a, z)$

$$\Phi \sim e^{\pm ikz}$$

$\Rightarrow \cos kz$ since even in z .

$$\Phi \sim \int_0^{\infty} dk \cos kz R_k(\rho) g_k$$

\Rightarrow continuous in k since infinite system

Equation for $R_k(\rho)$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} R_k - k^2 R_k = 0$$

\Rightarrow modified Bessel equation with $\nu=0$

$$R_k \sim I_0(k\rho), K_0(k\rho)$$

$\Rightarrow K_0 \rightarrow \infty$ at $\rho=0$

$$\Phi = \int_0^{\infty} dk \cos kz I_0(k\rho) g_k$$

BC at $e=a$

$$\Phi(e=a, z) = \int_0^\infty dk \cos kz I_0(ka) g_k$$

\Rightarrow multiply by $\cos k'z$ and integrate over z

$$V \int_{-L}^L dz \cos k'z = \int_0^\infty dk I_0(ka) g_k I$$

$$I = \int_{-L}^L dz \cos kz \cos k'z$$

$$= 2\pi \int_{-\infty}^{\infty} \frac{dt}{2\pi} \left(\frac{e^{ikt} - e^{-ikt}}{tc} \right) \left(\frac{e^{ik't} - e^{-ik't}}{tc} \right)$$

$$= 2\pi \frac{1}{4} [2\delta(k+ k') + 2\delta(k-k')] \quad \text{note: } k, k' > 0$$

$$V \frac{\sin k'L}{k'} \Big|_{-L}^L = \pi I_0(ka) g_k$$

$$V 2 \frac{\sin(k'L)}{k'} = \pi I_0(ka) g_k$$

$$\Phi = \frac{2V}{\pi} \int_0^\infty dk \frac{\sin kL}{k} \cos kz \frac{I_0(ka)}{I_0(ka)}$$

~~②) charge density $\rho(r, z) = \frac{2\sigma_0 \cos kz}{2\pi r}$
 since $\sigma_0 \int_0^{2\pi} d\phi \int_0^a dr \int_{-\infty}^{\infty} dz = \frac{\sigma_0 \int_0^{2\pi} d\phi \int_0^a dr}{2\pi r} = \sigma_0 a$~~

(2) a)

$$\nabla^2 \Phi = -\frac{\rho}{\epsilon_0} \delta(x - \frac{a}{2}) \delta(y - \frac{a}{2}) \delta(z)$$

Use oscillatory function in x, y directions

$$\Phi = \sum_{\substack{m, n \\ \text{odd}}} \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{a}\right) g_{mn}(z)$$

Because of symmetry around $a/2$, m, n are odd integers.

Inserting into Poisson's equation,

$$\sum_{m, n} \left(\frac{\partial^2}{\partial z^2} - K^2 \right) g_{mn} \sin(K_x x) \sin(K_y y) = -\frac{\rho}{\epsilon_0} \delta(x) \delta(y) \delta(z)$$

$$\text{with } K_x = \frac{n\pi}{a}, K_y = \frac{m\pi}{a}, K = (K_x^2 + K_y^2)^{\frac{1}{2}}$$

Eliminate sum over m, n by multiplying by $\sin\left(\frac{n'\pi x}{a}\right) \sin\left(\frac{m'\pi y}{a}\right)$ and integrating over x, y

$$\frac{a^2}{4} \left(\frac{\partial^2}{\partial z^2} - K^2 \right) g_{mn} = -\frac{\rho}{\epsilon_0} \underbrace{\sin\left(\frac{n'\pi}{2}\right)}_{(-1)^{\frac{n'-1}{2}}} \underbrace{\sin\left(\frac{m'\pi}{2}\right)}_{(-1)^{\frac{m'-1}{2}}} \delta(z)$$

$$\left(\frac{\partial^2}{\partial z^2} - k^2\right) g_{mn} = A_{mn} \delta(z)$$

$$A_{mn} = -\frac{2}{\epsilon_0} \frac{4}{a^2} (-1)^{\frac{n-1}{2}} (-1)^{\frac{m-1}{2}}$$

z > 0

$$g_{mn} = c_{mn} e^{-kz}$$

z < 0

$$g_{mn} = c_{mn} e^{kz}$$

z ≈ 0

$$\frac{\partial^2}{\partial z^2} g_{mn} = A_{mn} \delta(z)$$

⇒ integrating

$$\int_{-e}^e \frac{\partial^2}{\partial z^2} g_{mn} dz = A_{mn}$$

$$g_{mn} \Big|_{-e}^e = 0$$

$$(-k - k) c_{mn} = A_{mn}$$

$$c_{mn} = -\frac{1}{2k} A_{mn}$$

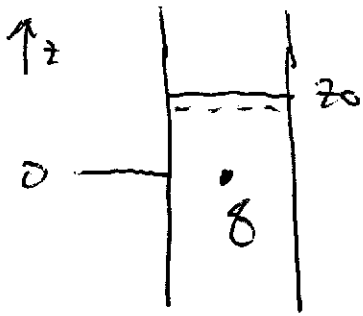
$$\underline{\underline{g}} = \sum_{\substack{m,n \\ \text{odd}}} \frac{16}{\epsilon_0} \frac{2}{ka^2} (-1)^{\frac{n-1}{2}} (-1)^{\frac{m-1}{2}} e^{-k|z|} \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi x}{a}\right)$$

(6)

For large positive $z \Rightarrow$ only smallest k contributes $\Rightarrow m=1, n=1$ $k = \sqrt{2} \left(\frac{\pi}{a} \right)$

$$\begin{aligned} \Phi &\approx \frac{q}{\epsilon_0} \frac{z}{\left[\sqrt{2} \left(\frac{\pi}{a} \right) a \right]^2} e^{-\frac{\sqrt{2} \pi z}{a}} \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{\pi y}{a}\right) \\ &= \frac{\sqrt{2} q}{\epsilon_0 \pi a} e^{-\frac{\sqrt{2} \pi z}{a}} \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{\pi y}{a}\right) \end{aligned}$$

b) Dielectric added in region $z > z_0 > 0$.



For $q > 0$, a negative charge will be induced on the surface of the dielectric at z_0

\Rightarrow attractive force with q

\Rightarrow force on q in positive z direction

c) Change in Φ given by $\delta\Phi$

$$\delta\Phi = \sum_{\substack{m,n \\ \text{odd}}} \delta q_{mn}(z) \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{a}\right)$$

\Rightarrow continuous across $z=0$ since Φ already produces jump due to q .

\Rightarrow will have jump in $\frac{\partial \Phi}{\partial z}$ across z_0

At surface of dielectric tangential \vec{E} is continuous.

$$\Rightarrow \left(\frac{\partial}{\partial x} \Phi + \frac{\partial}{\partial x} \delta \Phi \right) \Big|_{z_0-\epsilon}^{z_0+\epsilon} = 0$$

\Rightarrow but $\frac{\partial}{\partial x} \Phi$ continuous at z_0

$$\Rightarrow \frac{\partial}{\partial x} \delta \Phi \Big|_{z_0-\epsilon}^{z_0+\epsilon} = 0$$

$$\Rightarrow \delta q_{mn}(z) \Big|_{z_0-\epsilon}^{z_0+\epsilon} = 0 \Rightarrow \delta q_{mn} = \delta C_{mn} e^{-k|z-z_0|}$$

At surface of dielectric normal \vec{D} is continuous

$$\epsilon \left(\frac{\partial}{\partial z} \Phi + \frac{\partial}{\partial z} \delta \Phi \right) \Big|_{z_0-\epsilon}^{z_0+\epsilon} = 0$$

$$\epsilon \left(\frac{\partial}{\partial z} \Phi + \frac{\partial}{\partial z} \delta \Phi \right) \Big|_{z_0+\epsilon} = \epsilon_0 \left(\frac{\partial}{\partial z} \Phi + \frac{\partial}{\partial z} \delta \Phi \right) \Big|_{z_0-\epsilon}$$

$$\epsilon \left(\frac{\partial}{\partial z} q_{mn} + \frac{\partial}{\partial z} \delta q_{mn} \right) \Big|_{z_0+\epsilon} = \epsilon_0 \left(\frac{\partial}{\partial z} q_{mn} + \frac{\partial}{\partial z} \delta q_{mn} \right) \Big|_{z_0-\epsilon}$$

d) Once $\delta \Phi$ is known, calculate $\delta E_z = -\frac{\partial}{\partial z} \delta \Phi$
 $F_z = \rho \delta E_z$