

Green's Function Solutions of Poisson's Eqn.

We have shown previously that the solution of Poisson's eqn.

$$\nabla^2 \phi = -\frac{\rho}{\epsilon_0}$$

in free space is given by

$$\phi(\underline{x}) = \frac{1}{4\pi\epsilon_0} \int d\underline{x}' \frac{\rho(\underline{x}')}{|\underline{x}-\underline{x}'|}$$

This is because $\frac{1}{|\underline{x}-\underline{x}'|}$ is the Green's

function of ~~Laplace~~ Poisson's eqn with a δ -function source in free space,

$$\nabla^2 \frac{1}{|\underline{x}-\underline{x}'|} = -4\pi \delta(\underline{x}-\underline{x}')$$

More generally, we can have a system with conductors where ϕ or $\vec{E} \cdot \hat{n}$ are specified rather than simply

$\phi \rightarrow 0$ as $x \rightarrow \infty$. Can we construct a Green's function to enable us to solve such problems?

$$\phi(\underline{x}) = \int d\underline{x}' \delta(\underline{x} - \underline{x}') \phi(\underline{x}')$$

$$\text{Let } \nabla'^2 G(\underline{x}, \underline{x}') = -4\pi \delta(\underline{x} - \underline{x}')$$

G is the response at \underline{x}' from a source at \underline{x} .

$$\phi(\underline{x}) = -\frac{1}{4\pi} \int d\underline{x}' [\nabla'^2 G(\underline{x}, \underline{x}')] \phi(\underline{x}')$$

But

$$\begin{aligned} \nabla' \cdot (\phi(\underline{x}') \nabla' G) &= \nabla' \phi \cdot \nabla' G + \phi(\underline{x}') \nabla'^2 G \end{aligned}$$

$$\Rightarrow = -\frac{1}{4\pi} \int d\underline{x}' \underbrace{\nabla' \cdot (\phi(\underline{x}') \nabla' G)}_{\text{use divergence theorem}} + \frac{1}{4\pi} \int d\underline{x}' \underbrace{\nabla' \phi \cdot \nabla' G}_{\text{again write as divergence}} + \phi(\underline{x}') \frac{1}{\epsilon_0}$$

$$\begin{aligned} &= -\frac{1}{4\pi} \int ds' \phi(\underline{x}') \hat{n}' \cdot \nabla' G \\ &+ \frac{1}{4\pi} \int d\underline{x}' \left[\underbrace{\nabla' \cdot G \nabla' \phi}_{\text{divergence theorem}} - \underbrace{G \nabla'^2 \phi}_{-\phi(\underline{x}') \frac{1}{\epsilon_0}} \right] \end{aligned}$$

where $\hat{n}' = \hat{u}' \cdot \nabla'$

$$\phi(\underline{x}) = \frac{1}{4\pi\epsilon_0} \int_V d\underline{x}' G(\underline{x}, \underline{x}') \rho(\underline{x}') + \frac{1}{4\pi} \int_S ds' \left[G(\underline{x}, \underline{x}') \frac{\partial \phi}{\partial n'} - \phi(\underline{x}') \frac{\partial G}{\partial n'} \right]$$

This is completely general. In infinite system where $S \rightarrow \infty$ and $G = 1/|\underline{x} - \underline{x}'|$ same as earlier result.

Dirichlet BC's specify ϕ on S :
Choose $G(\underline{x}, \underline{x}') = 0$ on S so that

$$\phi(\underline{x}) = \frac{1}{4\pi\epsilon_0} \int_V d\underline{x}' G(\underline{x}, \underline{x}') \rho(\underline{x}') - \frac{1}{4\pi} \int_S ds' \phi(\underline{x}') \frac{\partial G}{\partial n'}$$

Neumann BC's specify $\partial \phi / \partial n'$ on S :
Would like to choose $\partial G / \partial n' = 0$.
However, this is not possible since

$$\begin{aligned} \int_S ds' \frac{\partial G}{\partial n'} &= \int_V d\underline{x}' \nabla' \cdot (\nabla' G) \\ &= -4\pi \int_V d\underline{x}' \delta(\underline{x} - \underline{x}') = -4\pi \end{aligned}$$

Best option is to choose $\frac{\partial G}{\partial n'} = -\frac{4\pi}{S}$

where $S = \int_S ds'$

Then,

$$\phi(x) = \frac{1}{4\pi\epsilon_0} \int dx' G(x, x') \rho(x')$$

$$+ \frac{1}{4\pi} \int ds' G(x, x') \frac{\partial \phi}{\partial n'}$$

$$+ \underbrace{\frac{1}{S} \int ds' \phi(x')}_{\langle \phi \rangle_S}$$

$\langle \phi \rangle_S =$ average over S of ϕ

$\langle \phi \rangle_S$ typically zero if some part of surface is at ∞ .

Conclusion: Have reduced the problem of solving $\nabla^2 \phi = -\rho/\epsilon_0$ to solving

$$\nabla^2 G = -4\pi \delta(x - x')$$

with simple BC's.

Green's function for a sphere with Dirichlet BCs

$G(\underline{x}, \underline{x}')$ satisfies

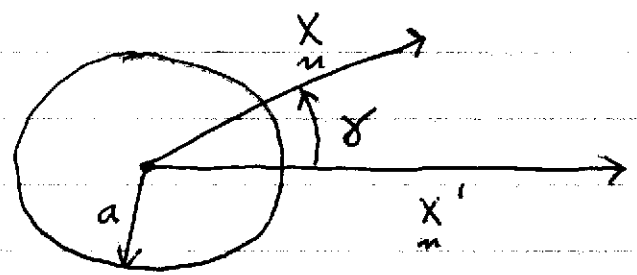
$$\nabla^2 G(\underline{x}, \underline{x}') = -4\pi \delta(\underline{x} - \underline{x}')$$

For Dirichlet BC want $G(\underline{x}, \underline{x}') = 0$ for \underline{x} on the surface. We previously calculated Q for a point charge outside a grounded sphere. For that problem

$$\nabla^2 Q(\underline{x}) = -\frac{\sigma}{\epsilon_0} \delta(\underline{x} - \underline{x}')$$

Let $\underline{x} \rightarrow \underline{x}'$ and $\frac{\sigma}{4\pi\epsilon_0} = 1$. This yields

$$G(\underline{x}, \underline{x}') = \frac{1}{|\underline{x} - \underline{x}'|} - \frac{a}{|\underline{x}'|} \frac{1}{\left| \underline{x} - \frac{a^2}{|\underline{x}'|^2} \underline{x}' \right|}$$



$$|\underline{x} - \underline{x}'| = \left[(\underline{x} - \underline{x}') \cdot (\underline{x} - \underline{x}') \right]^{\frac{1}{2}} = \left[x^2 + x'^2 - 2xx' \cos \theta \right]^{\frac{1}{2}}$$

$$G = \frac{1}{[x^2 + x'^2 - 2xx' \cos \gamma]^{1/2}} - \frac{a}{[x^2 + x'^2 + a^4 - 2xx'a^2 \cos \gamma]^{1/2}}$$

\Rightarrow note that $G(\underline{x}, \underline{x}')$ is symmetric when $\underline{x} \leftrightarrow \underline{x}'$

To calculate $\mathcal{Q}(\underline{x})$ also need $\frac{\partial G}{\partial x'} \Big|_{x'=a} = - \frac{\partial G}{\partial x'} \Big|_{x'=a}$

$$\frac{\partial G}{\partial x'} = -\frac{1}{2} \cdot 2 \frac{(x' - x \cos \gamma)}{[x' - x \cos \gamma]^3} + \frac{1}{2} a \cdot 2 \frac{(x' x^2 - x a^2 \cos \gamma)}{(x^2 + x'^2 + a^4 - 2xx'a^2 \cos \gamma)^{3/2}}$$

$$\frac{\partial G}{\partial x'} \Big|_{x'=a} = - \frac{a - x \cos \gamma}{|\underline{x} - a \hat{n}'|^3} + \frac{a^2}{a^3} \frac{(x^2 - x a \cos \gamma)}{(x^2 + a^2 - 2x a \cos \gamma)^{3/2}}$$

$$\underbrace{\hspace{10em}}_{|\underline{x} - \hat{n}' a|^3}$$

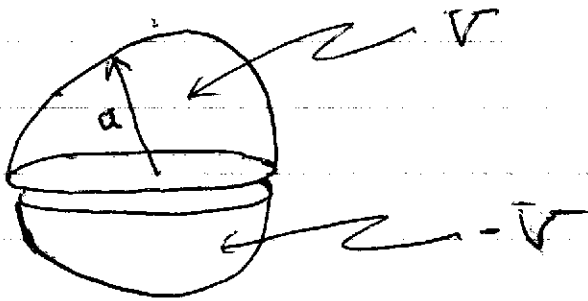
$$= \frac{1}{a |\underline{x} - a \hat{n}'|^3} \left[-a^2 + a x \cos \gamma + x^2 - x a \cos \gamma \right]$$

Note that \hat{n}' is actually inward since the volume V is outside the sphere

$$\frac{\partial G}{\partial n'} = - \frac{\partial G}{\partial x'} \Big|_{x'=a} = - \frac{(x^2 - a^2)}{a |\underline{x} - a \hat{n}'|^3}$$

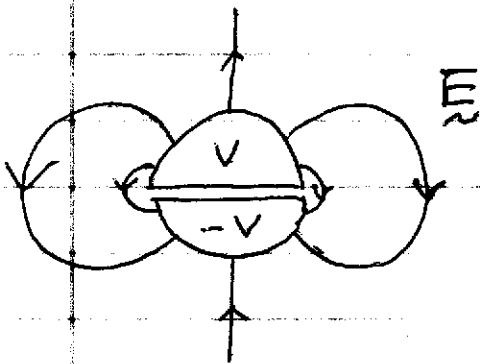
$$\mathcal{Q}(\underline{x}) = \int_V d\underline{x}' \rho(\underline{x}') G(\underline{x}, \underline{x}') - \frac{1}{4\pi} \int_S ds' \sigma(\underline{x}') \frac{\partial G}{\partial n'}$$

Hemisphere's at different potentials



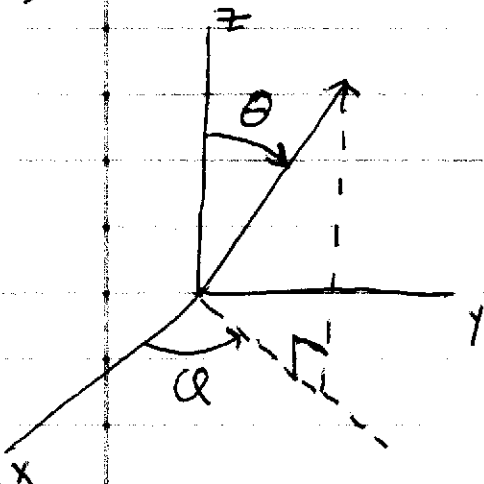
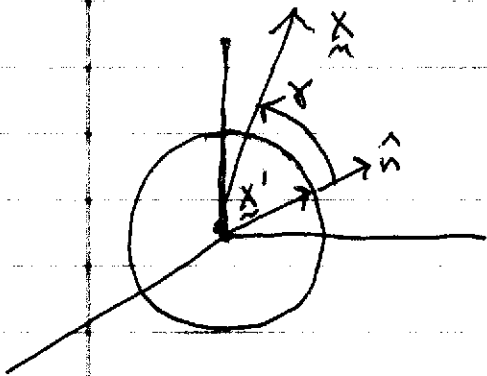
No change outside

$$\phi(\underline{x}) = 0 \text{ for } |\underline{x}| > a$$



$$\phi(\underline{x}) = - \int \frac{ds'}{4\pi} \phi(\underline{x}') \frac{\partial G}{\partial n'}$$

$$\frac{\partial G}{\partial n'} \Big|_a = - \frac{(x^2 - a^2)}{a |\underline{x} - a\hat{n}|^3}$$



$$\frac{\underline{x}}{|\underline{x}|} = \hat{u} \sin\theta \cos\phi + \hat{j} \sin\theta \sin\phi + \hat{k} \cos\theta$$

⇒ similar for $\hat{n}(\theta', \phi')$

ϕ here is azimuthal angle

$$\frac{\vec{x} \cdot \vec{n}}{|\vec{x}|} = \cos \gamma = \sin \theta \cos \varphi \sin \theta' \cos \varphi' + \sin \theta \sin \varphi \sin \theta' \sin \varphi' + \cos \theta \cos \theta'$$

identity: $\cos \varphi \cos \varphi' + \sin \varphi \sin \varphi' = \cos(\varphi - \varphi')$

$$\cos \gamma = \sin \theta \sin \theta' \cos(\varphi - \varphi') + \cos \theta \cos \theta'$$

⇒ integrate over surface of sphere with $ds' = a^2 d\cos\theta' d\varphi'$

$$\cos\theta' \in (-1, 1) \text{ and } \varphi' \in (0, 2\pi)$$

⇒ Can't do surface integral exactly

⇒ consider case with $\theta = 0$

⇒ along the axis of the sphere

$$\Rightarrow \cos \gamma = \cos \theta'$$

$$\Rightarrow |\vec{x}| = z$$

$$\frac{\partial G}{\partial n'} = - \frac{(z^2 - a^2)}{a(z^2 + a^2 - 2za \cos \theta')^{3/2}}$$

$$Q = -\frac{1}{4\pi} a^2 \bar{V} 2\pi \left[\int_0^1 d\cos\theta' - \int_{-1}^0 d\cos\theta' \right] \frac{dG}{\sqrt{r'}}$$

$$= + \frac{a^2 \bar{V}}{2} \frac{(z^2 - a^2)}{-2za} \left[\int_0^1 d\cos\theta' - \int_{-1}^0 d\cos\theta' \right] \frac{-2za}{(z^2 + a^2 - 2za\cos\theta')}$$

$$= \frac{V(z^2 - a^2)}{-4z} \left[\frac{-2}{(z^2 + a^2 - 2za\cos\theta')^{1/2}} \right] \left[\begin{matrix} 1 & -1 \\ 0 & -1 \end{matrix} \right]$$

$$= \frac{V}{2z} (z^2 - a^2) \left[\frac{1}{(z^2 + a^2 - 2za)^{1/2}} - \frac{1}{(z^2 + a^2)^{1/2}} + \frac{1}{(z^2 + a^2 + 2za)^{1/2}} \right]$$

$$= \frac{V}{2z} (z^2 - a^2) \left[\frac{1}{z-a} - \frac{2}{(z^2 + a^2)^{1/2}} + \frac{1}{z+a} \right]$$

$$= \frac{V}{2z} (z^2 - a^2) \left[\frac{2z}{z^2 - a^2} - \frac{2}{(z^2 + a^2)^{1/2}} \right]$$

$$Q = V \left[1 - \frac{z^2 - a^2}{z(z^2 + a^2)^{1/2}} \right] \text{ for } z > a.$$

As $z \rightarrow a \Rightarrow Q \rightarrow V$

~~As~~ For $z \gg a$

$$Q \approx V \left[1 - \frac{z^2(1 - \frac{a^2}{z^2})}{z^2} \left(1 - \frac{1}{2} \frac{a^2}{z^2} \right) \right]$$

$$= V \left[1 - 1 + \frac{a^2}{z^2} + \frac{1}{2} \frac{a^2}{z^2} \right] = \frac{3}{2} V \frac{a^2}{z^2}$$

\Rightarrow Falls off faster than $\frac{1}{z}$. Why?

\Rightarrow Consider large $|x| \equiv r \gg a$
with arbitrary angle θ

$$\begin{aligned} \frac{\partial G}{\partial n'} &= - \frac{(r^2 - a^2)}{a (r^2 + a^2 - 2ar \cos \gamma)^{3/2}} \\ &= - \frac{r^2 (1 - \frac{a^2}{r^2})}{a r^3 (1 + \frac{a^2}{r^2} - \frac{2a}{r} \cos \gamma)^{3/2}} \\ &\approx - \frac{1}{ar} (1 + 3 \frac{a}{r} \cos \gamma) \end{aligned}$$

$$\phi = - \frac{a^2 V}{4\pi} \left(-\frac{1}{ar}\right) \int_0^{2\pi} d\phi' \left[\int_0^1 d\cos\theta' - \int_{-1}^0 d\cos\theta' \right]$$

\otimes $\left[1 + 3 \frac{a}{r} (\underbrace{\sin\theta \sin\theta' \cos(\phi - \phi')}_{\text{zero when averaged over } \theta'} + \underbrace{\cos\theta \cos\theta'}_{\text{periodic in } \phi' \Rightarrow \text{zero average}}) \right]$

$$= \frac{3a^2 V}{2r^2} \cos\theta \left[\int_0^1 d\cos\theta' - \int_{-1}^0 d\cos\theta' \right] \cos\theta'$$

$\underbrace{\hspace{10em}}_{2 \int_0^1 d\cos\theta'}$

$$\phi = \frac{3}{2} \frac{a^2}{r^2} V \cos\theta \Rightarrow \text{dipole field}$$