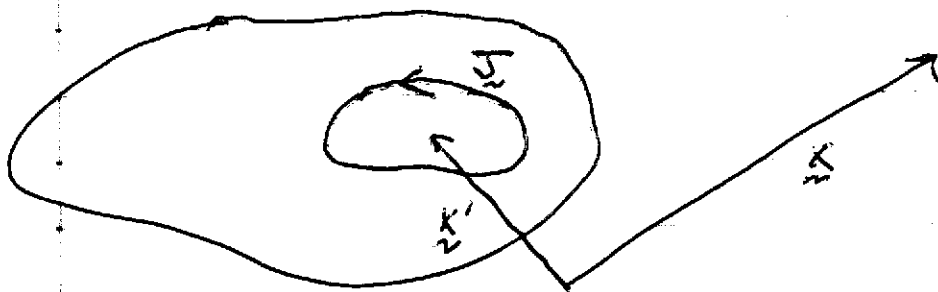


Magnetic Fields from localized current distributions

As in the case of localized charge distributions, we want to be able to calculate the magnetic field from localized current distributions. This will also make it possible to calculate how materials respond to magnetic fields.



$$\vec{A} = \frac{\mu_0}{4\pi} \int d\vec{x}' \vec{J}(\vec{x}') \frac{1}{|\vec{x} - \vec{x}'|}$$

For $|\vec{x}| \gg |\vec{x}'|$

$$|\vec{x} - \vec{x}'| = \left(x^2 + x'^2 - 2 \vec{x} \cdot \vec{x}' \right)^{1/2}$$

$$\approx x \left(1 + \frac{x'^2}{x^2} - \frac{2 \vec{x} \cdot \vec{x}'}{x^2} \right)^{1/2}$$

$$\approx x \left(1 - \frac{\vec{x} \cdot \vec{x}'}{x^2} + \dots \right)$$

$$\frac{1}{|\vec{x} - \vec{x}'|} \approx \frac{1}{x} + \frac{\vec{x} \cdot \vec{x}'}{x^3} + \dots$$

$$\vec{A} = \frac{\mu_0}{4\pi} \int d\vec{x}' \vec{J}(\vec{x}') \left[\frac{1}{x} + \frac{\vec{x} \cdot \vec{x}'}{x^3} + \dots \right]$$

$$\int d\vec{x}' \nabla' \cdot (\vec{J}(\vec{x}') \vec{x}') = 0 \quad \text{since } \vec{J} \rightarrow 0 \text{ at } \infty$$

$$= \int d\vec{x}' \left[\underbrace{(\nabla' \cdot \vec{J}(\vec{x}')) \vec{x}'}_{=0 \text{ since } \nabla' \cdot \vec{J} = 0 \text{ for } \frac{\partial}{\partial t} = 0} + \underbrace{\vec{J}(\vec{x}') \cdot \nabla' \vec{x}'}_{\vec{J}(\vec{x}') \cdot \vec{I}} \right]$$

$$\nabla' \cdot \vec{J} = 0$$

$$\text{for } \frac{\partial}{\partial t} = 0$$

$$\vec{J}(\vec{x}') \cdot \vec{I}$$

$\vec{I} = \text{unit tensor}$

$$\Rightarrow \int d\vec{x}' \vec{J}(\vec{x}') = 0$$

\Rightarrow no monopole contributions to \vec{A}

$$\vec{A} = \frac{\mu_0}{4\pi} \frac{\vec{x}}{x^3} \int d\vec{x}' \vec{x}' \vec{J}(\vec{x}')$$

$$\vec{x} \times (\vec{x}' \times \vec{J}) = \vec{x} \cdot \vec{J}(\vec{x}') \vec{x}' - \vec{x} \cdot \vec{x}' \vec{J}(\vec{x}')$$

$$\int d\vec{x}' \vec{J}(\vec{x}') \vec{x}' = \int d\vec{x}' (\vec{J} \cdot \nabla' \vec{x}') \vec{x}'$$

$$= - \int d\vec{x}' \vec{x}' \nabla' \cdot (\vec{J} \vec{x}') \quad \vec{I}$$

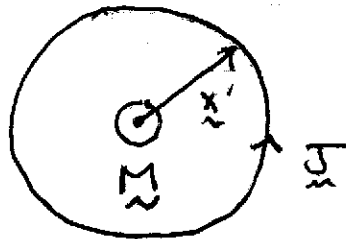
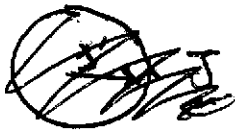
$$\begin{aligned}
 &= -\int d\vec{x}' \vec{x}' \cdot \vec{\nabla}' \cdot \vec{x}' = -\int d\vec{x}' \vec{x}' \cdot \vec{\nabla}' \cdot \vec{I} \\
 &= -\int d\vec{x}' \vec{x}' \cdot \vec{J}(\vec{x}')
 \end{aligned}$$

$$\Rightarrow \int d\vec{x}' \vec{x}' \times [\vec{x}' \times \vec{J}(\vec{x}')] = -2 \int d\vec{x}' \vec{x}' \cdot \vec{J}(\vec{x}')$$

$$\vec{A} = -\frac{1}{2} \frac{\mu_0}{4\pi} \frac{1}{x^3} \int d\vec{x}' \vec{x}' \times [\vec{x}' \times \vec{J}(\vec{x}')]$$

Define $\vec{M}(\vec{x}') = \frac{1}{2} \vec{x}' \times \vec{J}(\vec{x}')$

= magnetic moment density



$$\vec{m} = \int d\vec{x}' \vec{M}(\vec{x}') = \text{magnetic moment}$$

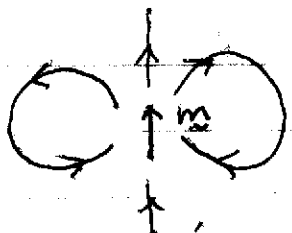
$$\sim L^3 \left(L \frac{\text{amps}}{L^2} \right) \sim \text{amps } L^2$$

$\sim I \text{ area}$

$$\vec{A} = -\frac{\mu_0}{4\pi} \frac{1}{x^3} \vec{x}' \times \vec{m}$$

$$\begin{aligned}
 \vec{B} &= \nabla \times \vec{A} = -\frac{\mu_0}{4\pi} \nabla \times \left(\frac{\vec{x} \times \vec{m}}{x^3} \right) \\
 &= -\frac{\mu_0}{4\pi} \left[\vec{m} \cdot \nabla \left(\frac{\vec{x}}{x^3} \right) - \vec{m} \cdot \underbrace{\nabla \cdot \left(\frac{\vec{x}}{x^3} \right)}_{-\nabla \cdot \left(\frac{1}{x} \right)} \right] \\
 &= -\frac{\mu_0}{4\pi} \vec{m} \cdot \nabla \left(\frac{\vec{x}}{x^3} \right) \quad \underbrace{-\nabla^2 \frac{1}{x} = 4\pi \delta(\vec{x})}_{=0} \\
 &= -\frac{\mu_0}{4\pi} \left[\frac{\vec{m} \cdot \vec{x}}{x^3} - \vec{x} \cdot \vec{m} \cdot \nabla \frac{1}{x^3} \right]
 \end{aligned}$$

$$\nabla \frac{1}{x^3} = -\frac{3}{2} \frac{1}{x^5} \cdot 2\vec{x} = -3 \frac{\vec{x}}{x^5}$$

$$\vec{B} = \frac{\mu_0}{4\pi} \left(\frac{3\vec{x}\vec{x} \cdot \vec{m} - \vec{m}}{x^3} \right)$$


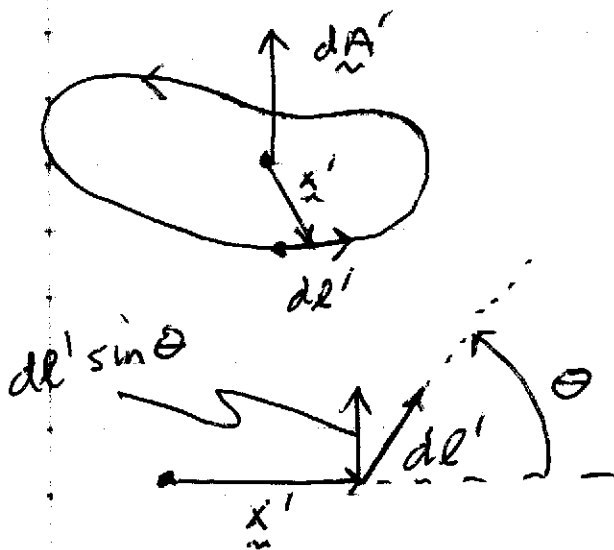
\Rightarrow same form as electric dipole

\Rightarrow magnetic field far from a localized current distribution is to lowest order dipolar.

Magnetic moment from a plane current loop

$$\underline{m} = \frac{1}{2} \int d\underline{x}' \underline{x}' \times \underline{J}(\underline{x}')$$

$$= \frac{1}{2} I \int \underline{x}' \times d\underline{\ell}'$$



$$\frac{1}{2} \underline{x}' \times d\underline{\ell}' = \underline{x}' \frac{dl' \sin \theta}{2}$$

$$= dA'$$

= area of segment

$$\underline{m} = I \int d\underline{A}' = I \underline{A}$$

A = area of loop

\Rightarrow direction of \underline{A} given by right hand rule.

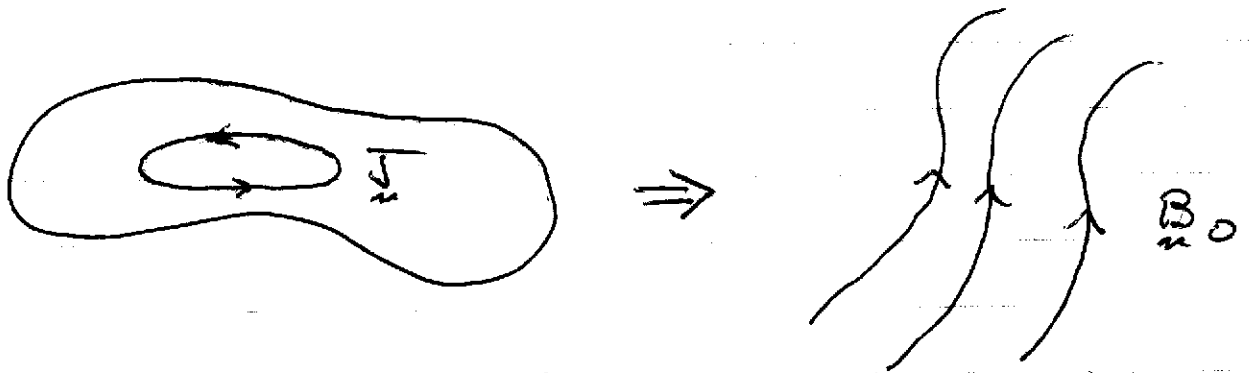
\Rightarrow generalizes result from circular current loop where found

$$m = I \pi a^2$$

\Rightarrow can extend to quadrupole moments

Magnetic forces

We previously calculated the energy of a charge distribution in an external electric field \vec{E}_0 and then calculated the forces using virtual displacements. We can, in principle, do the same for currents.



However, energy is required to maintain the currents I because magnetic flux cuts through the current loops.

⇒ need Faraday's law which we have not yet discussed.

⇒ Instead calculate the forces directly.

$$\vec{F} = \int d\vec{x}' \vec{J}(\vec{x}') \times \vec{B}(\vec{x}')$$

where \vec{B} is the external field

and not the field produced by \underline{J}

$\Rightarrow \underline{J}$ can not accelerate itself

Assume that \underline{B} is slowly varying over the spatial scale of \underline{J} .

$$\underline{B}(\underline{x}') \cong \underline{B}(\underline{x}) + (\underline{x}' - \underline{x}) \cdot \nabla \underline{B}(\underline{x}) + \dots$$

$$\underline{E} = \int d\underline{x}' \underline{J}(\underline{x}') \times \underline{B}(\underline{x}) + \int d\underline{x}' \underline{J}(\underline{x}') \times [(\underline{x}' - \underline{x}) \cdot \nabla \underline{B}]$$

$$= \int d\underline{x}' [\underline{J}(\underline{x}') \underline{x}' \cdot \nabla] \times \underline{B}(\underline{x})$$

$$\int d\underline{x}' \nabla \times [\underline{x}' \times \underline{J}(\underline{x}')] = \int d\underline{x}' [\underline{x}' \underline{J} \cdot \nabla - \underline{J} \underline{x}' \cdot \nabla]$$

\Rightarrow note that this is an operator eqn where ∇ does not act on \underline{x}' or $\underline{J}(\underline{x})$.

$$= -2 \int d\underline{x}' \underline{J}(\underline{x}') \underline{x}' \cdot \nabla$$

where $\underline{x}' \underline{J}(\underline{x}') \Rightarrow -\underline{J}(\underline{x}') \underline{x}'$ as before

$$\begin{aligned}
 \vec{F} &= -\frac{1}{2} \int d\vec{x}' \left(\nabla \times \left[\vec{x}' \times \vec{J}(\vec{x}') \right] \right) \times \vec{B}(\vec{x}) \\
 &= -(\nabla \times \vec{m}) \times \vec{B} = \nabla (\vec{m} \cdot \vec{B}) \\
 &\quad - \vec{m} \cdot \nabla \vec{B}
 \end{aligned}$$

$$\vec{F} = \nabla (\vec{m} \cdot \vec{B})$$

\Rightarrow where ∇ acts on $\vec{B}(\vec{x})$ only

\Rightarrow \vec{m} is the magnetic moment of the current distribution

Can write

$$U = -\vec{m} \cdot \vec{B}$$

as the energy of \vec{m} in an external field \vec{B} with

$$\vec{F} = -\nabla U$$

Torque on a current distribution

Consider the torque $d\vec{N}' = \vec{x}' \times d\vec{F}'$ on a volume element $d\vec{x}'$

$$\vec{N} = \int d\vec{N}' = \int d\vec{x}' \vec{x}' \times (\vec{J}(\vec{x}') \times \vec{B}(\vec{x}'))$$

To lowest order $\underline{B}(\underline{x}') = \underline{B}(0)$ where take the origin of \underline{J} to be at zero.

$$\begin{aligned}
N_{\underline{z}} &= \int d\underline{x}' \underline{x}' \times \left[\underline{J}(\underline{x}') \times \underline{B}(0) \right] \\
&= \int d\underline{x}' \left[\underline{B} \cdot \underline{x}' \underline{J}(\underline{x}') - \underline{B} \underline{x}' \cdot \underline{J}(\underline{x}') \right]
\end{aligned}$$

$$\begin{aligned}
\int d\underline{x}' \underline{x}' \cdot \underline{J}(\underline{x}') &= \int d\underline{x}' \underline{J}(\underline{x}') \cdot \nabla' \left(\frac{\underline{x}' \cdot \underline{x}'}{2} \right) \\
&= - \int d\underline{x}' \frac{x'^2}{2} \nabla' \cdot \underline{J}(\underline{x}') = 0
\end{aligned}$$

in steady state.

$$N = \underline{B} \cdot \int d\underline{x}' \underline{x}' \underline{J}(\underline{x}')$$

Have

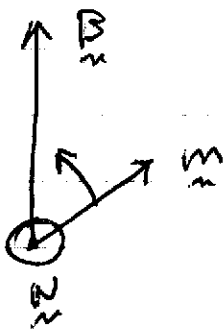
$$\begin{aligned}
\underline{B} \times \left[\underline{x}' \times \underline{J}(\underline{x}') \right] &= \underline{x}' \times \underline{J}(\underline{x}') \cdot \underline{B} \\
&\quad - \underline{J}(\underline{x}') \underline{x}' \cdot \underline{B} \\
&\Rightarrow - 2 \underline{J}(\underline{x}') \underline{x}' \cdot \underline{B}
\end{aligned}$$

where again $\underline{x}' \underline{J} \Rightarrow - \underline{J} \underline{x}'$ when integrated over \underline{x}' .

$$\Rightarrow \vec{B} \cdot \vec{x}' \vec{J}(\vec{x}') \Rightarrow -\frac{1}{2} \vec{B} \times [\vec{x}' \times \vec{J}(\vec{x}')]]$$

$$\vec{N} = -\frac{1}{2} \vec{B} \times \int d\vec{x}' [\vec{x}' \times \vec{J}(\vec{x}')] = -\vec{B} \times \vec{m}$$

$$\vec{N} = \vec{m} \times \vec{B}$$



\vec{m} twists to align with \vec{B}

\Rightarrow consistent with

$$U = -\vec{m} \cdot \vec{B} = -mB \cos \theta$$

$$N = -\frac{\partial U}{\partial \theta} = mB \sin \theta$$

