

Fourier Series and Fourier ~~and Laplace~~ Transforms

Boas Ch. 7

Periodic Functions

Many phenomena in physics are periodic in time and/or space.

⇒ spring-mass system ⇒ time

⇒ R-C circuit ⇒ time

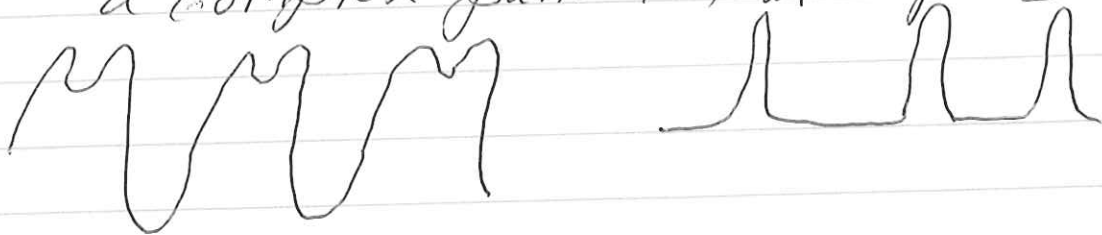
⇒ waves ⇒ space-time

$$P \approx P_0 \cos\left[\frac{2\pi}{\lambda}(x-vt)\right]$$

⇒ at a given time periodic in space

⇒ at a given location periodic in time

⇒ A periodic function consist of a complex pattern that repeats



⇒ need to develop techniques to describe complicated but periodic functions

⇒ consider a generic system that is periodic in space. Choose a spatial coordinate such that the periodicity length is 2π

$\cos\left(\frac{2\pi x}{\lambda}\right) \Rightarrow \text{let } \frac{2\pi x}{\lambda} \rightarrow x$

$\Rightarrow \cos x, \sin x$

⇒ periodic over 2π

⇒ same for time

$\cos \frac{2\pi}{T} t \Rightarrow \frac{2\pi}{T} t \rightarrow t$

$\Rightarrow \cos t, \sin t$

⇒ want to represent an arbitrary function that is periodic over 2π .

$\Rightarrow f(x)$

Sin and Cos Series

⇒ express $f(x)$ in terms of sines and cosines

$$f(x) = \frac{1}{2} a_0 + \sum_{n=1}^{\infty} a_n \cos(nx) + \sum_{n=1}^{\infty} b_n \sin(nx)$$

To calculate the a_n 's and b_n 's we use the orthogonality of

all of these functions over the interval $(0, 2\pi)$.

$$I_{sc} = \int_0^{2\pi} dx \sin(nx) \cos(mx) = 0 \quad \text{all } m, n$$

$$I_{ss} = \int_0^{2\pi} dx \sin(nx) \sin(mx) = 0 \quad m \neq n$$

$$I_{cc} = \int_0^{2\pi} dx \cos(nx) \cos(mx) = 0 \quad m \neq n$$

Consider
$$I_{cc} = \int_0^{2\pi} dx \left(\frac{e^{inx} - e^{-inx}}{2i} \right) \left(\frac{e^{imx} - e^{-imx}}{2i} \right)$$

\Rightarrow all of the form

$$\int_0^{2\pi} dx e^{iqx} = \frac{e^{iqx}}{iq} \Big|_0^{2\pi} = 0$$

for $m \neq n$, where q is an integer.

$\Rightarrow I_{cc} = 0$

\Rightarrow same for I_{ss}, I_{sc}

\Rightarrow multiply $f(x)$ by $\sin(n'x)$ and integrate over 2π

$$\int_0^{2\pi} dx f(x) \sin(n'x) = \sum_{n=1}^{\infty} b_n \int_0^{2\pi} dx \sin(nx) \sin(n')$$

$$= \sum_{n=1}^{\infty} b_n \frac{1}{2} \delta_{nn'} = b_{n'} \pi$$

since $\int_0^{2\pi} dx \sin^2(n'x) = \frac{1}{2} (2\pi)$.

$$\Rightarrow \boxed{b_n = \frac{1}{\pi} \int_0^{2\pi} dx f(x) \sin(nx)} \quad \text{for } n=1, 2, \dots$$

Similar for a_n

$$\boxed{a_n = \frac{1}{\pi} \int_0^{2\pi} dx f(x) \cos(nx)} \quad n=1, 2, \dots$$

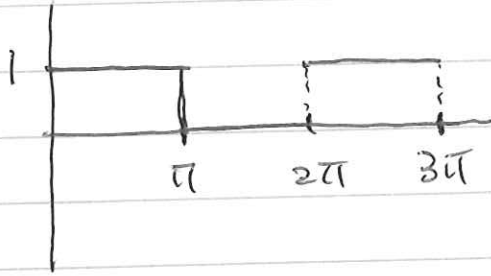
For a_0 integrate f over $(0, 2\pi)$

$$\int_0^{2\pi} dx f(x) = \frac{1}{2} a_0 2\pi = a_0 \pi$$

$$\boxed{a_0 = \frac{1}{\pi} \int_0^{2\pi} dx f(x)}$$

Same as for other a_n 's.

Example Consider the "square wave" function



$$a_0 = \frac{1}{\pi} \int_0^{2\pi} dx f(x)$$

$$= \frac{1}{\pi} \int_0^{\pi} dx 1 = 1$$

$$a_n = \frac{1}{\pi} \int_0^{\pi} dx \cos(nx) = \frac{1}{\pi} \frac{\sin nx}{n} \Big|_0^{\pi} = 0$$

for $n \neq 0$

$$b_n = \frac{1}{\pi} \int_0^{\pi} dx \sin(nx) = -\frac{1}{\pi} \frac{\cos nx}{n} \Big|_0^{\pi}$$

$$= -\frac{1}{\pi} \frac{\cos n\pi - 1}{n}$$

$$= -\frac{1}{\pi} \frac{(-1)^n - 1}{n} = \begin{cases} 0 & \text{never} \\ \frac{2}{n\pi} & \text{for } n \text{ odd} \end{cases}$$

$$f(x) = \frac{1}{2} + \frac{2}{\pi} \sum_{n \text{ odd}} \frac{1}{n} \sin(n\pi)$$

Under what conditions does the Fourier
Sin and Cos series converge to f ?

\Rightarrow are these functions a "complete set"?

Dirichlet's Theorem: If f has a finite # of
max and min values and discontinuities then the
series converges where f is continuous.

Convergence in the mean

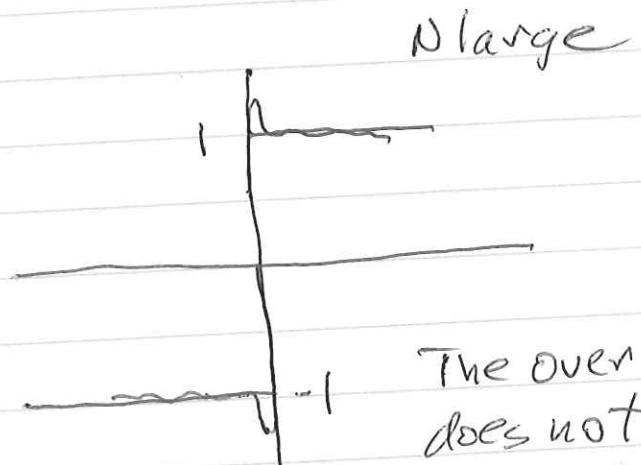
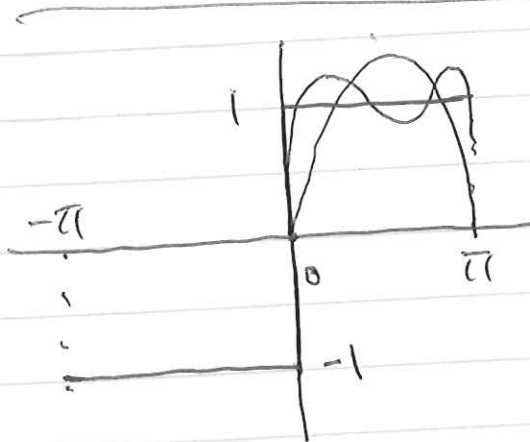
Introduce a weak convergence criterion

$$\lim_{N \rightarrow \infty} \frac{1}{2\pi} \int_0^{2\pi} dx (f(x) - f_N(x))^2 = 0$$

$$f_N = \frac{1}{2} a_0 + \sum_{n=1}^N (a_n \cos nx + b_n \sin nx)$$

$f(x)$ can deviate from f_N at a finite # of points

Gibbs Phenomenon



The overshoot does not go away for large N . Width ~~get~~ gets smaller.

Note: can often exploit the symmetry of the problem to simplify the evaluation of a_n, b_n

\Rightarrow in the above problem all of the a_n are zero.

Complex exponential Series

We know that the sines and cosines can be written as exponentials

$$\cos(nx) = \frac{1}{2} (e^{inx} + e^{-inx})$$

$$\sin(nx) = \frac{1}{2i} (e^{inx} - e^{-inx})$$

It therefore seems reasonable to represent periodic functions in terms of complex exponentials.

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx}$$

To find the c_n we first note that the average of e^{iqx}

over 2π is zero — we showed this ~~earlier~~ earlier. Thus, if we multiply the series by $e^{-in'x}$

and integrate over 2π , we find

$$\begin{aligned} \int_0^{2\pi} dx f(x) e^{-in'x} &= \sum_{n=-\infty}^{\infty} c_n \int_0^{2\pi} dx e^{i(n-n')x} \\ &= \sum_{n=-\infty}^{\infty} 2\pi c_n \delta_{nn'} = 2\pi c_{n'} \end{aligned}$$

Thus,

$$C_n = \frac{1}{2\pi} \int_0^{2\pi} dx f(x) e^{-inx}$$

For $f(x)$ real

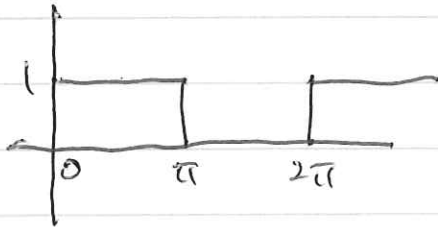
$$f^* = f(x) = \sum_n C_n^* e^{-inx} = \sum_n C_n e^{inx}$$

comparing the coefficients of e^{-inx}

$$\Rightarrow C_n^* = C_{-n}$$

for f real

example - square wave



$$C_n = \frac{1}{2\pi} \int_0^{\pi} dx e^{-inx} = \frac{1}{2\pi} \left[\frac{e^{-inx}}{-in} \right]_0^{\pi} = \frac{1}{2\pi} \frac{e^{-in\pi} - 1}{-in} = \begin{cases} 0 & \text{even } n \\ \frac{1}{\pi in} & \text{odd } n \end{cases}$$

$$C_0 = \frac{1}{2}$$

$$f(x) = \frac{1}{2} + \frac{1}{\pi i} \sum_{n=1,3,\dots} \frac{1}{n} e^{inx}$$

$$+ \frac{1}{\pi i} \sum_{n=-1,-3,\dots} \frac{1}{n} e^{inx}$$

$$= \frac{1}{2} + \frac{2}{\pi i} \sum_{n=1,3,\dots} \frac{1}{n} \frac{e^{inx} - e^{-inx}}{2i}$$

$$= \frac{1}{2} + \frac{2}{\pi} \sum_{n=1,3,\dots} \frac{1}{n} \sin(nx) \quad \text{as before}$$

Other intervals

Suppose $f(x)$ is periodic over the length $2l$. Choose as basis functions

$$\sin\left(\frac{\pi x n}{l}\right), \cos\left(\frac{\pi x n}{l}\right)$$

Again, these are orthogonal over the interval $(0, 2l)$ or $(-l, l)$.

Replace $\frac{1}{2\pi} \int_0^{2\pi} dx$ by $\frac{1}{2l} \int_0^{2l} dx$ or $\frac{1}{2l} \int_{-l}^l dx$

$$a_n = \frac{1}{l} \int_0^{2l} dx f(x) \cos\left(\frac{\pi x n}{l}\right)$$

etc.

Parseval's Theorem

Calculate the average of f^2 over the interval $(0, 2\pi)$

$$\frac{1}{2\pi} \int_0^{2\pi} dx f^2(x) = \frac{1}{2\pi} \int_0^{2\pi} dx \left[\frac{1}{2} a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx) \right]$$

Since each of the functions is orthogonal to the others unless the index n is the same

$$\Rightarrow \frac{1}{2\pi} \int_0^{2\pi} dx \left[\left(\frac{a_0}{2}\right)^2 + \sum_{n=1}^{\infty} (a_n^2 \cos^2 nx + b_n^2 \sin^2 nx) \right]$$

$$= \cancel{\frac{1}{2\pi} \int_0^{2\pi} dx} \left(\frac{a_0}{2} \right)^2 + \frac{1}{2} \sum_{n=1}^{\infty} (a_n^2 + b_n^2)$$

$$\boxed{\frac{1}{2\pi} \int_0^{2\pi} dx f^2 = \left(\frac{a_0}{2}\right)^2 + \frac{1}{2} \sum_{n=1}^{\infty} (a_n^2 + b_n^2)}$$

This is also called the completeness relation.

\Rightarrow not really useful for calculating $\langle f^2 \rangle$

\Rightarrow tells you if the sines and cosines form a complete set.

If some terms in the series are discarded \Rightarrow e.g. $n=1$ then will find

$$\frac{1}{2\pi} \int_0^{2\pi} dx f^2 > \left(\frac{a_0}{2}\right)^2 + \sum_{n=2}^{\infty} (a_n^2 + b_n^2)$$

The equality must be satisfied if the basis functions form a complete set.

For the exponential series

$$\frac{1}{2\pi} \int_0^{2\pi} dx |f|^2 = \sum_n |c_n|^2$$

Fourier Transform

We have explored how to represent periodic functions but many systems are not periodic \Rightarrow how do we treat such systems

\Rightarrow allow the periodicity length l to go to infinity

Start with a system of length $2l$,

$$f(x) = \sum_{-\infty}^{\infty} c_n e^{i \frac{n\pi x}{l}}$$

where calculate c_n by multiplying by ~~$e^{-i \frac{n'\pi x}{l}}$~~ $e^{-i \frac{n'\pi x}{l}}$

and integrating over $(-l, l)$

$$\int_{-l}^l dx e^{-i \frac{n'\pi x}{l}} f(x) = 2l c_{n'}$$

change n' to n and substitute above.

$$f(x) = \sum_{-\infty}^{\infty} c_n e^{i \frac{n\pi x}{l}} = \frac{1}{2l} \sum_{-\infty}^{\infty} \int_{-l}^l dx' e^{-i \frac{n\pi x'}{l}} f(x')$$

where use x' in the integral to not confuse it with $e^{i \frac{n\pi x}{l}}$

Define $k_n = \frac{n\pi}{l}$ and note the increment of k_n is $k_{n+1} - k_n = \frac{\pi}{l}(n+1 - n) = \frac{\pi}{l} \equiv \Delta k_n$

$$f(x) = \sum_{k_n=-\infty}^{\infty} e^{ik_n x} \Delta k_n \frac{1}{2\pi} \int_{-l}^l dx' e^{-ik_n x'} f(x')$$

Let

$$F(k) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dx' e^{-ikx'} f(x')$$

~~what~~

Over what scale does $F(k)$ vary?

example $f(x') = \frac{1}{x'^2 + a^2}$

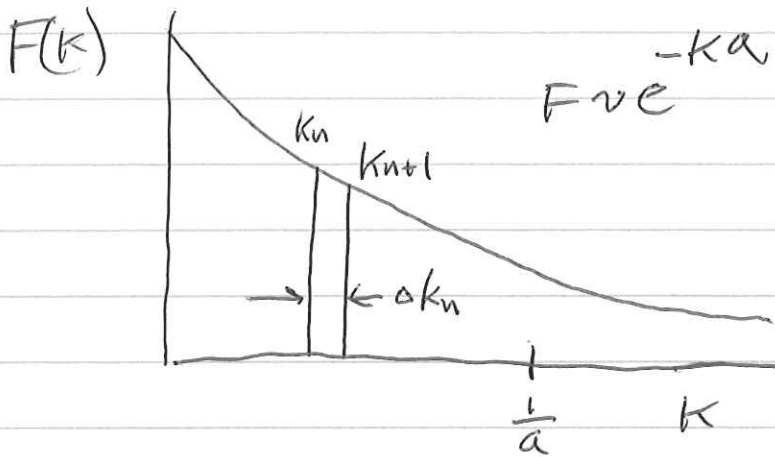
$$F(k) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dx' \frac{e^{-ikx'}}{x'^2 + a^2}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} dz \frac{e^{-ikz}}{z^2 + a^2}$$

$$= \frac{-2\pi i}{2\pi} \frac{e^{-ikz}}{z - ia} \Big|_{z = -ia}$$

$$= \frac{-i e^{-ka}}{-2ia} = \frac{1}{2a} e^{-ka}$$

~~$e^{-ikx'}$~~
 e^{-ikz} bounded in LHP. close contour and residue at $z = -ia$



$F(k)$ falls off with a characteristic scale $\sim \frac{1}{a}$

The point is that $\Delta k_n \sim \frac{\pi}{l} \ll \frac{1}{a}$

\Rightarrow this means that

$$\sum_{k_n=-\infty}^{\infty} \Delta k_n \Rightarrow \int_{-\infty}^{\infty} dk_n$$

as long as $l \gg a$ where "a" is the characteristic scale of $f(x)$. Thus

$$f(x) = \int_{-\infty}^{\infty} dk e^{ikx} F(k)$$

$$F(k) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dx' e^{-ikx'} f(x')$$

$F(k)$ is the Fourier transform of $f(x)$ and the integral over k is the inverse transform, which yields back $f(x)$.

\Rightarrow the Fourier representation is very useful for solving differential equations.

example Represent $f(x) = \begin{cases} 1, & -1 < x < 1 \\ 0, & |x| > 1 \end{cases}$
as a Fourier integral.

$$F(k) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dx' e^{-ikx'} f(x')$$

$$= \frac{1}{2\pi} \int_{-1}^1 dx' e^{-ikx'} = \frac{1}{2\pi} \left. \frac{e^{-ikx'}}{-ik} \right|_{-1}^1$$

$$= \frac{1}{2\pi} \frac{e^{-ik} - e^{ik}}{-ik} = \frac{\sin k}{\pi k}$$

So

$$f(x) = \int_{-\infty}^{\infty} dk e^{ikx} \frac{\sin k}{\pi k}$$

or

$$f(x) = \int_{-\infty}^{\infty} dk \frac{\sin k}{\pi k} (\cos kx + i \sin kx)$$

$$= \int_{-\infty}^{\infty} dk \frac{\sin k \cos kx}{\pi k}$$

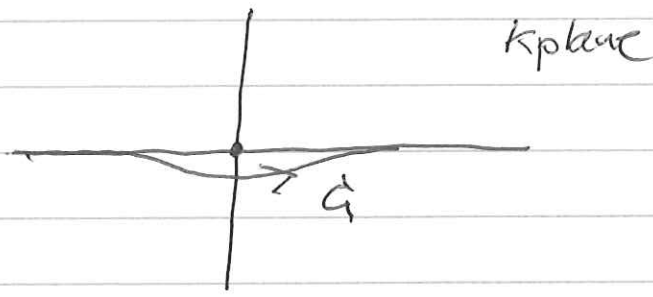
since the $\sin(kx)$ term is an odd function while $\sin(k)/k$ is even

Now complete the inverse transform to see that we obtain $f(x)$

$$f(x) = \int_{-\infty}^{\infty} dk e^{ikx} \left(\frac{e^{ik} - e^{-ik}}{2i\pi k} \right)$$

note: no singularity at $k=0$
 \Rightarrow move contour integral away from $k=0$

$$f(x) = \int_{\Gamma} dk \left(\frac{e^{ik(x+1)} - e^{ik(x-1)}}{2i\pi k} \right)$$



For $x > 1$, close in UHP for both exponentials
 $\Rightarrow x+1$ and $x-1$ positive
 \Rightarrow residue at $k=0$ is zero since two exponentials cancel.

For $x < -1$, close in LHP since $x+1$ and $x-1$ both negative
 \Rightarrow no enclosed pole
 \Rightarrow zero

For $-1 < x < 1$, $x+1 > 0$ and $x-1 < 0$.
 Close $e^{ik(x+1)}$ piece in UHP and $e^{ik(x-1)}$ in LHP. Only $e^{ik(x+1)}$ has a nonzero residue

$$f(x) = \frac{2\pi i}{2\pi i} e^{ik(x+1)} \Big|_{k=0} = 1$$

example Dirac delta function $\delta(x)$

$\delta(x)$ has ~~the~~ ^{the} property that its amplitude goes to ∞ but its width goes to zero such that its area is unity

$$\int_{-\infty}^{\infty} dx \delta(x) = 1$$

What is the Fourier representation of $\delta(x)$?

$$\hookrightarrow F(k) = \int_{-\infty}^{\infty} dx' e^{-ikx'} \delta(x') \frac{1}{2\pi}$$

$$= \frac{1}{2\pi}$$

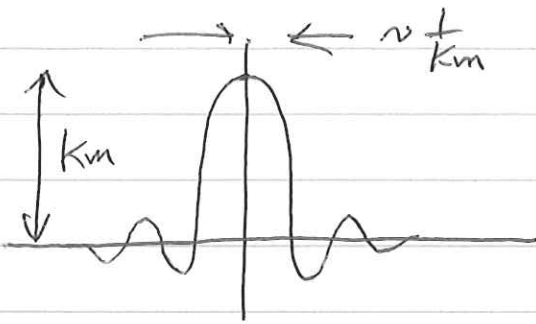
$$\Rightarrow \boxed{\delta(x) = \int_{-\infty}^{\infty} dk e^{ikx} \frac{1}{2\pi}}$$

Can think of this as

$$\delta(x) = \lim_{K_m \rightarrow \infty} \frac{1}{2\pi} \int_{-K_m}^{K_m} dk e^{ikx}$$

$$= \lim_{K_m \rightarrow \infty} \frac{1}{2\pi} \frac{e^{iK_mx} - e^{-iK_mx}}{ix}$$

$$= \lim_{K_m \rightarrow \infty} \frac{1}{\pi} \frac{\sin K_mx}{x}$$



For $km \ll \pi$,

$$\delta \sim km$$

Parseval's Theorem for Fourier Integrals

$$\int_{-\infty}^{\infty} dx |f(x)|^2 = \int_{-\infty}^{\infty} dx \underbrace{\int_{-\infty}^{\infty} dk e^{ikx} F(k)}_{f(x)} \underbrace{\int_{-\infty}^{\infty} dk' e^{-ik'x} F^*(k')}_{f^*(x)}$$

$$= \int_{-\infty}^{\infty} dk F(k) \int_{-\infty}^{\infty} dk' F^*(k') \underbrace{\int_{-\infty}^{\infty} dx e^{i(k-k')x}}_{2\pi \delta(k-k')}$$

$$= 2\pi \int_{-\infty}^{\infty} dk F(k) F(k)^*$$

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} dx |f(x)|^2 = \int_{-\infty}^{\infty} dk |F(k)|^2$$