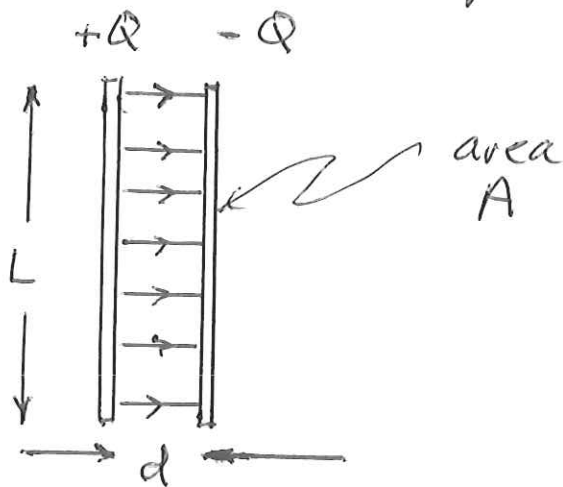


Capacitors and dielectrics

A capacitor is a device consisting of two conductors separated by an insulating material. A capacitor can store charge and therefore the energy associated with that charge.

Parallel plate capacitor

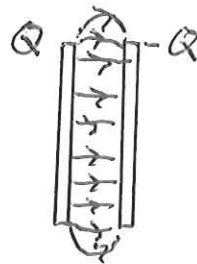


Assume that $d \ll L$

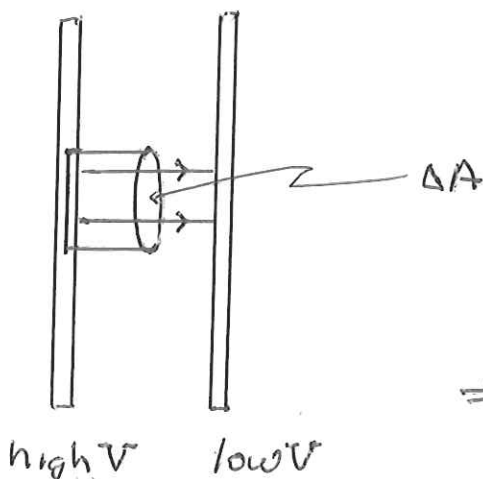
$\Rightarrow$ can calculate E by assuming that the plates are effectively of infinite extent

$\Rightarrow$ fringe field neglected

$$\sigma = \frac{Q}{A} = \frac{\text{charge}}{\text{area}}$$



Use Gauss' Law to find E :



Choose Gaussian surface with one side in conductor and other between the plates.

$$\oint \vec{E} \cdot d\vec{A} = E \Delta A = \frac{\sigma \Delta A}{\epsilon_0}$$

$\Rightarrow$ no E flux through sides or in conductor

$$E = \frac{\sigma}{\epsilon_0} = \frac{Q}{A \epsilon_0}$$

Potential drop:

$$\Delta V = Ed = \frac{Qd}{A\epsilon_0}$$

Define "capacitance" as the ability to store charge

$$C \equiv \frac{Q}{\Delta V} = \frac{Q A \epsilon_0}{Qd} = \frac{A \epsilon_0}{d}$$

$\Rightarrow C$ is a function of the geometry of the capacitor and not the charge

units: C is measured in Farads

$$\text{Farad} = \frac{C}{J/C} = \frac{C^2}{J}$$

A Farad is an unrealistically large capacitance

$\Rightarrow$ Take $d = 1 \text{ cm}$ and $C = 1 \text{ Farad}$

$$A = \frac{1 \text{ farad } d}{\epsilon_0} = \frac{1 \text{ farad } 10^{-2} \text{ m } \text{Nm}^2}{8.85 \times 10^{-12} \text{ C}^2}$$

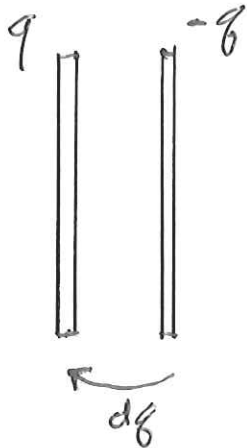
$$\approx 10^9 \frac{\cancel{\text{C}^2}}{\cancel{\text{Nm}}} \text{m} \frac{\text{Nm}^2}{\cancel{\text{C}^2}} = 10^9 \text{ m}^2$$

More typically measure capacitance in

$$\text{micro farads} = 10^{-6} \text{ farad} = 1 \mu\text{F}$$

$$\text{pico farads} = 10^{-12} \text{ farad} = 1 \text{ pF}$$

Energy stored in a capacitor



Consider a capacitor which at some time has a charge q and a voltage

$$V = \frac{q d}{\epsilon_0 A} = \frac{q}{C}$$

Move a small charge dq from $-q$ to q . The work required is

$$dW = V dq = \frac{1}{C} q dq$$

To find the total work required to assemble the charge from $q=0$ to $q=Q$, integrate dW

$$W = \int dW = \int_0^Q \frac{q dq}{C} = \frac{q^2}{2C} \Big|_0^Q = \frac{Q^2}{2C}$$

$\Rightarrow$ energy stored in capacitor

$$U = \frac{Q^2}{2C} = \frac{1}{2} C V^2 = \frac{1}{2} Q V$$

example :

Consider a capacitor with $d = 1 \text{ mm}$ and $A = 1 \text{ cm}^2$ charged to 100 volts. How much energy is stored?

$$C = \frac{A \epsilon_0}{d} = \frac{8.85 \times 10^{-12} \frac{\text{C}^2}{\text{Nm}^2} \cdot \frac{10^{-4} \text{ m}^2}{10^{-3} \text{ m}}}{10^{-3} \text{ m}} = 8.85 \times 10^{-13} \text{ F}$$

$$U = \frac{1}{2} C V^2 = \frac{1}{2} 8.85 \times 10^{-13} \frac{\text{C}^2}{\text{Nm}} \cdot 10^4 \frac{\text{J}^2}{\text{C}^2} = 4.4 \times 10^{-9} \text{ J}$$

Energy in the electric field

The energy stored in the capacitor to infer the energy in an electric field

$$\begin{aligned}
 U &= \frac{1}{2} C V^2 = \frac{1}{2} \frac{A \epsilon_0}{d} \cancel{V^2}^2 E^2 d^2 \\
 &= \frac{1}{2} \epsilon_0 (A d) E^2 = \underbrace{(A d)}_{\text{Volume}} \frac{1}{2} \epsilon_0 E^2
 \end{aligned}$$

$$u = \frac{U}{\text{Vol}} = \frac{1}{2} \epsilon_0 E^2$$

= energy density of the local electric field

$\Rightarrow$ this is a generic result

Calculating capacitance

Calculate the electric field

$\Rightarrow$ Gauss' Law or another approach

$\Rightarrow$ Calculate the potential difference between conductors or between conductor and infinity

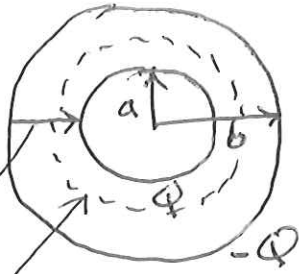
$$V_2 - V_1 = - \int_{x_1}^{x_2} \underline{E} \cdot d\underline{x}$$

example: Cylindrical capacitor

Length L

inner conductor ~~at~~ radius a

outer conductor radius b



$$\oint \vec{E} \cdot d\vec{A} = \frac{Q}{\epsilon_0}$$

$\vec{E}, d\vec{A}$ are radial

end caps drop out since

$\vec{E} \cdot d\vec{A} \approx 0$ there

$$E 2\pi r L = \frac{Q}{\epsilon_0}$$

$$E = \frac{Q}{2\pi r L \epsilon_0}$$

$$V = V_a - V_b = - \int_b^a E dr = - \frac{Q}{2\pi L \epsilon_0} \int_b^a \frac{dr}{r}$$

$$= - \frac{Q}{2\pi L \epsilon_0} \ln\left(\frac{a}{b}\right)$$

$$V = \frac{Q}{2\pi L \epsilon_0} \ln\left(\frac{b}{a}\right) \quad \text{note: } \ln\left(\frac{b}{a}\right) = -\ln\left(\frac{a}{b}\right)$$

$$C = \frac{Q}{V} = \frac{2\pi L \epsilon_0}{\ln(b/a)}$$

example: capacitance of an isolated sphere

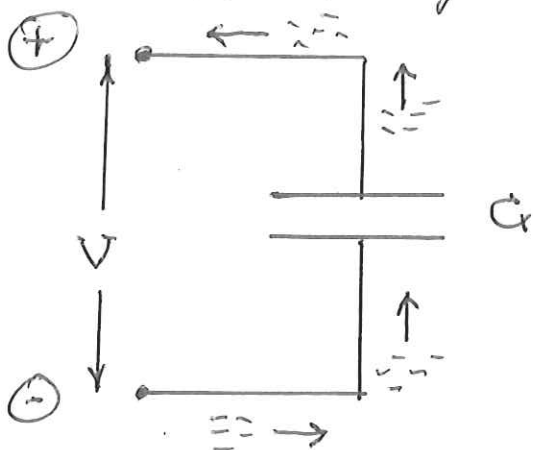


$$V = \frac{Q}{4\pi \epsilon_0 R} \Rightarrow C = 4\pi \epsilon_0 R$$

$\Rightarrow$ negative plate at ∞ .

Charging a capacitor

Apply a potential V across wires connected to the capacitor plates (battery, ...)



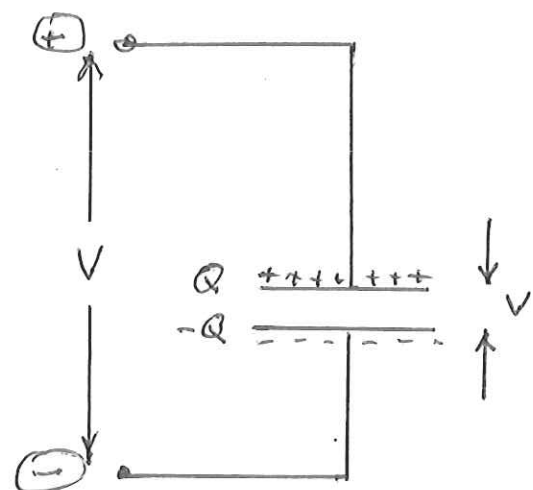
Electrons in the ^{upper} plate flow toward the (+).

Electron flow away from the (-) to the lower plate

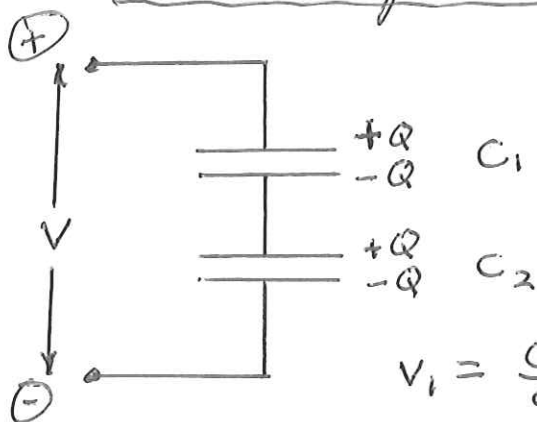
⇒ positive charge left on upper plate and negative on lower plate

⇒ charge flows until potential across the capacitor matches the imposed potential

$$\Rightarrow Q = CV$$



Series capacitors



Two capacitors in series

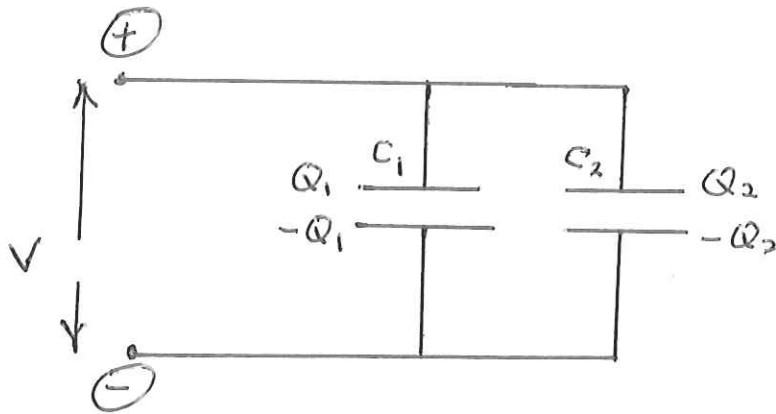
⇒ charge on top of C_2 and bottom of C_1 are equal in magnitude but opposite sign ⇒ charge conservation

$$V_1 = \frac{Q}{C_1}, V_2 = \frac{Q}{C_2}$$

$$V = V_1 + V_2 = \frac{Q}{C_1} + \frac{Q}{C_2} = Q \left(\frac{1}{C_1} + \frac{1}{C_2} \right)$$

$$\frac{1}{C} = \frac{1}{C_1} + \frac{1}{C_2}$$

Capacitors in parallel



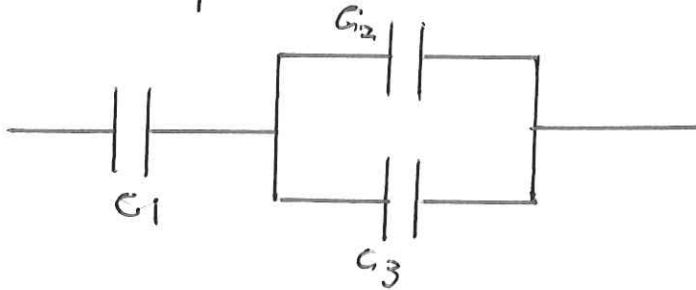
$$Q_1 = C_1 V$$

$$Q_2 = C_2 V$$

$$Q = Q_1 + Q_2 = V(C_1 + C_2)$$

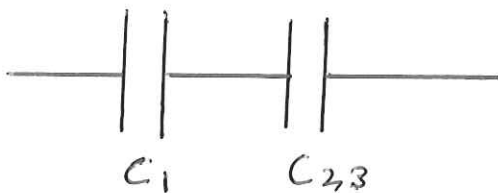
$$C = C_1 + C_2$$

More complicated combinations



C_2 and C_3 are in parallel $\Rightarrow$ replace by single capacitor

$$C_{23} = C_2 + C_3$$



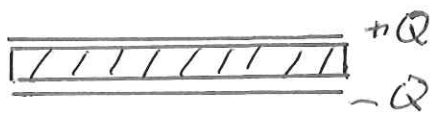
$\Rightarrow$ combine C_1 and C_{23} which are in parallel

$$\frac{1}{C} = \frac{1}{C_1} + \frac{1}{C_{23}}$$

Polarization of dielectrics

A dielectric is a non-conduction material (insulator) that is placed between the plates of a capacitor $\Rightarrow$ mechanically holds the plates apart

$$V_0 = \frac{Q}{C_0}$$

When the dielectric is inserted, the potential V across the capacitor is reduced

$$V = \frac{V_0}{K}, \quad K = \text{dielectric}$$

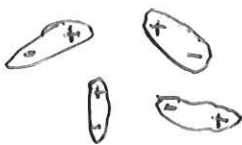
with $K > 1$

Thus, the capacitance increases

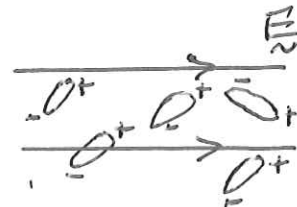
$$C = \frac{Q}{V} = \frac{Q}{V_0/K} = \frac{Q}{V_0} K > \frac{Q}{V_0} = C_0$$

$$C = K C_0 \quad \Rightarrow \text{why?}$$

① polar materials have molecules in which the charge is not evenly distributed



molecules align



$\Rightarrow$ positive portions shift in direction of E and negative portions in the opposite direction

$\Rightarrow$ recall motion of a dipole in E .

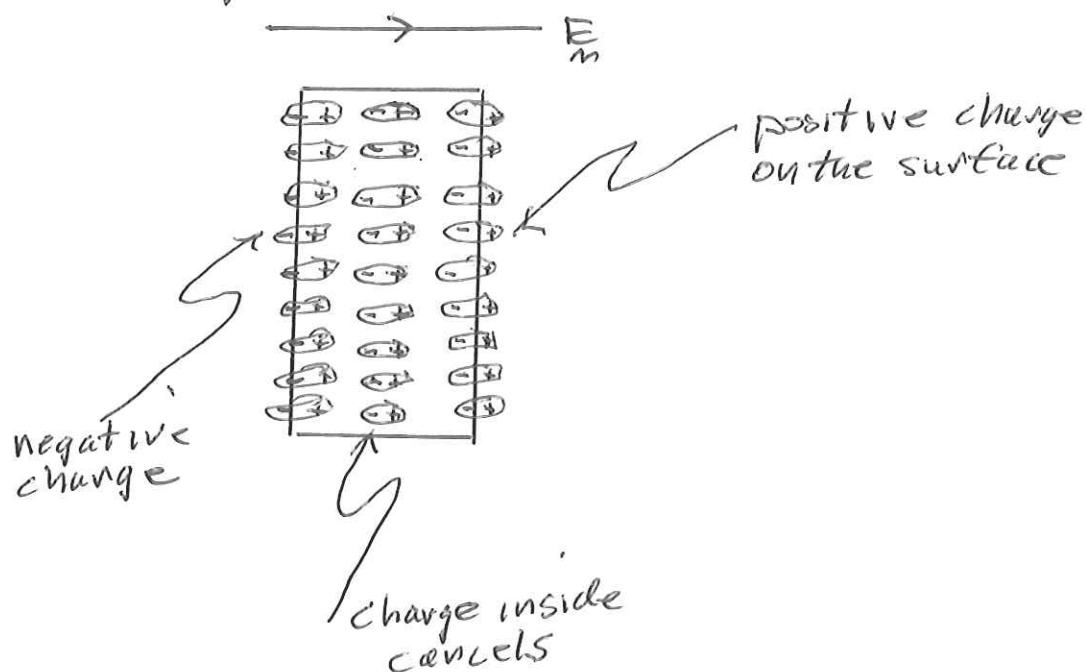
② non polar molecules have an even charge distribution



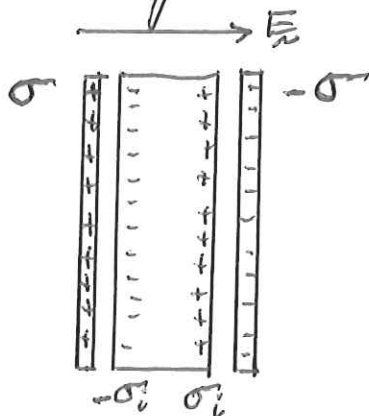
electrons shift in $-E_m$ direction
plus charges shift along E_m

$\Rightarrow$ the dielectric is polarized by E_m

The response of the dielectric



In the capacitor



σ_i is the induced charge on the surface of the dielectric

$$\Rightarrow \text{total } \frac{\text{charge}}{\text{area}} = \sigma - \sigma_i'$$

$$E = \frac{\sigma - \sigma_i'}{\epsilon_0} \quad \text{and} \quad E_0 = \frac{\sigma}{\epsilon_0}$$

(51)

$$\frac{E}{E_0} = \frac{Ed}{E_0 d} = \frac{V}{V_0} = \frac{1}{k}$$

$$\frac{E}{E_0} = \frac{\sigma - \sigma_i}{\epsilon_0} \frac{1}{\sigma/\epsilon_0} = \frac{\sigma - \sigma_i}{\sigma} = \frac{1}{k}$$

$$\sigma - \sigma_i = \frac{\sigma}{k} \Rightarrow \sigma_i = \sigma \left(1 - \frac{1}{k}\right)$$

Can also define a permittivity of the dielectric

$$E = \frac{\sigma - \sigma_i}{\epsilon_0} = \frac{\sigma}{k\epsilon_0} \equiv \frac{\sigma}{\epsilon}$$

$$\epsilon = k\epsilon_0 > \epsilon_0 \quad \text{with} \quad C = \frac{\epsilon A}{d}$$

$\epsilon = \text{permittivity}$

Dielectric constants of some materials

	k	dielectric strength
air	1.0006	$3 \times 10^6 \text{ V/m}$
oil	4	$12 \times 10^6 \text{ V/m}$
glass	5	$14 \times 10^6 \text{ V/m}$
water (liquid)	80	

The "dielectric strength" is the electric field above which the dielectric breaks down and a current flows.

Energy density with dielectrics

$U = \frac{1}{2} Q V \Rightarrow$ where V is reduced
compared with no dielectric

Since $Q = V C$

$$U = \frac{1}{2} C V^2 = \frac{1}{2} \frac{\epsilon A}{d} E^2 d^2 = \frac{1}{2} \epsilon E^2 (A d)$$

$u = \frac{1}{2} \epsilon E^2 =$ energy density with a dielectric

$\Rightarrow$ for a given E , it is greater than in vacuum

Gauss' Law with dielectrics

Gauss' Law is given by

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q}{\epsilon_0} \quad \text{where } Q \text{ is the total charge within the Gaussian surface}$$

$$Q = Q_f + Q_i$$

$Q_f =$ free charge

$Q_i =$ induced charge

A charge in a uniform dielectric is shielded by the induced charge

$$Q_f + Q_i = \frac{Q_f}{K}$$

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_f}{K \epsilon_0} = \frac{Q_f}{\epsilon}$$

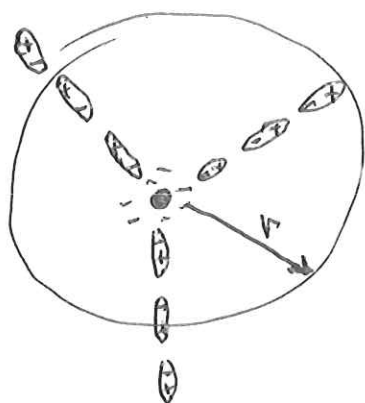
If the dielectric is not uniform, then must include ϵ in the calculation of the electric flux

$$\oint \epsilon \vec{E}_n \cdot d\vec{A} = Q_{\text{enc}}$$

$\vec{D} = \epsilon \vec{E}_n = \text{electric displacement}$

$\Rightarrow$ evaluate ϵ at the Gaussian surface

example: Charge Q in a uniform dielectric



The charge Q is shielded by the dielectric

$$\oint \epsilon \vec{E}_n \cdot d\vec{A} = Q$$

$$\epsilon_0 k E 4\pi r^2 = Q$$

$$E = \frac{Q}{k} \frac{1}{4\pi \epsilon_0 r^2}$$

$\Rightarrow$ reduced E

What if the dielectric is a sphere of radius R with Q at the center?

$\Rightarrow$ for $r < R$ E is unchanged from above

For $r > R$

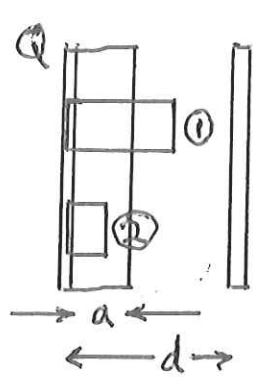


$$\oint \epsilon \vec{E}_n \cdot d\vec{A} = \epsilon_0 E 4\pi r^2 = Q$$

$$E = \frac{Q}{4\pi \epsilon_0 r^2} \Rightarrow \text{same as vacuum}$$

$\Rightarrow$ positive induced charge on surface at $r = R$

example: Capacitor with dielectric



The dielectric has a width a that is smaller than the plate separation d .

First use Gauss' Law in (1) to find $E = E_1$ in vacuum

$$\oint \epsilon E_1 dA = \epsilon_0 E_1 A = Q$$

$$E_1 = \frac{Q}{\epsilon_0 A}$$

$\Rightarrow$ dielectric has no effect. Why?

Now use Gauss' Law in (2) to find $E = E_2$ inside the dielectric

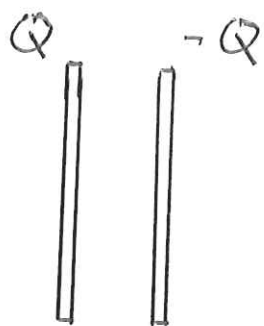
$$\oint \epsilon E_2 dA = \epsilon E_2 A = Q$$

$$E_2 = \frac{Q}{\epsilon A} = \frac{Q}{k\epsilon_0 A}$$

Total potential

$$V = E_2 a + E_1 (d-a) = \frac{Q}{A\epsilon_0} \left[\frac{a}{k} + d-a \right]$$

Forces of capacitor plates



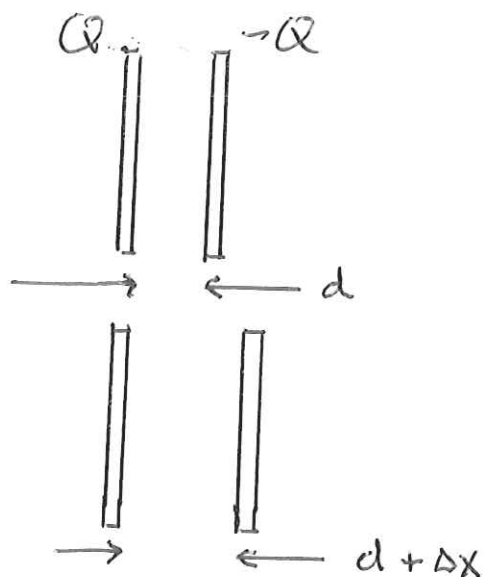
Can calculate the force on the plates (attractive) in two ways:

① Direct evaluation

$$F = qE$$

② Find the change in energy due to a small virtual displacement of a plate and equate to the work done

⇒ Do ② first



$$U_i = \frac{1}{2} \epsilon_0 E^2 (Ad)$$

= initial energy

$$U_f = \frac{1}{2} \epsilon_0 E^2 (d + \Delta x) A$$

⇒ note E is unchanged if Q is fixed

$$\text{work done} = F \Delta x$$

$$= U_f - U_i = \frac{1}{2} \epsilon_0 E^2 \Delta x A$$

$$F = \frac{1}{2} \epsilon_0 E^2 A$$

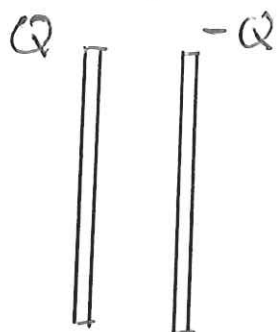
⇒ note the force between the plates is attractive

⇒ rewrite F in terms of E and Q

$$E = \frac{\sigma}{\epsilon_0} = \frac{Q}{A \epsilon_0}$$

$$F = \frac{1}{2} \epsilon_0 E A \left(\frac{Q}{A \epsilon_0} \right) = \frac{1}{2} Q E$$

Why is there a factor of $\frac{1}{2}$?



$\rightarrow E_Q$

$\leftarrow E$

$\Rightarrow$ a charge can not accelerate itself. E comes from both Q and $-Q$. What is E in the location of $-Q$ due only to Q ?

$$E_Q = \frac{\sigma}{2\epsilon_0} \Rightarrow E \text{ from a sheet of charge } \sigma$$

$$F = E_Q (-Q) = \frac{\sigma}{2\epsilon_0} Q \text{ (attractive)}$$

$$= \frac{1}{2} \epsilon_0 E^2 A$$

$\Rightarrow$ same as from energy argument.