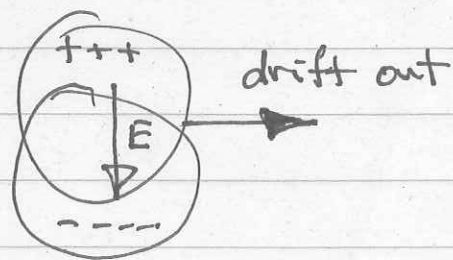
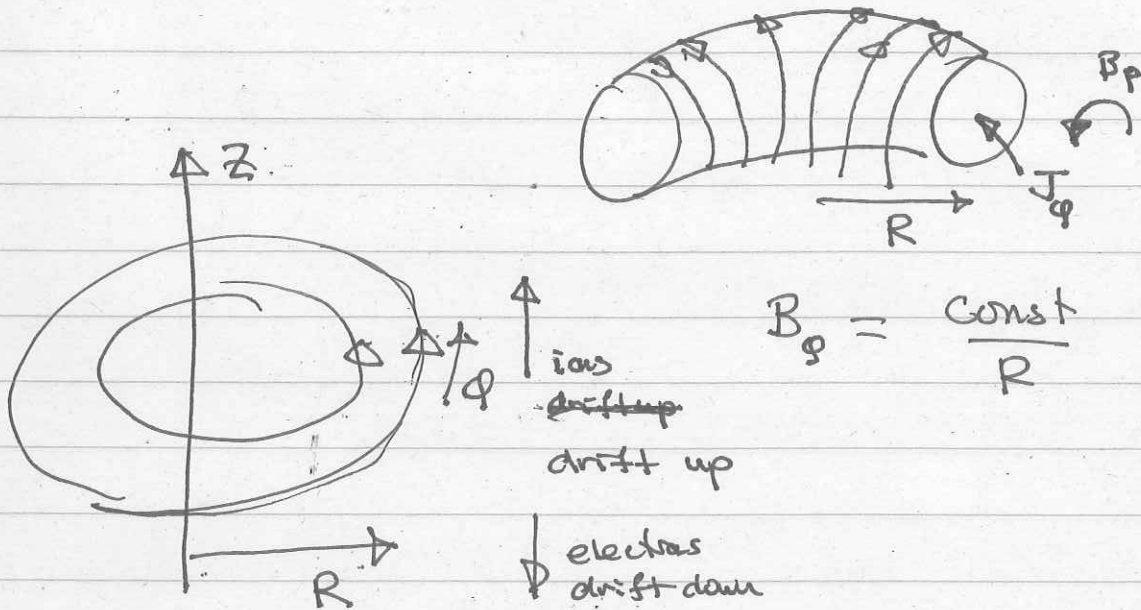
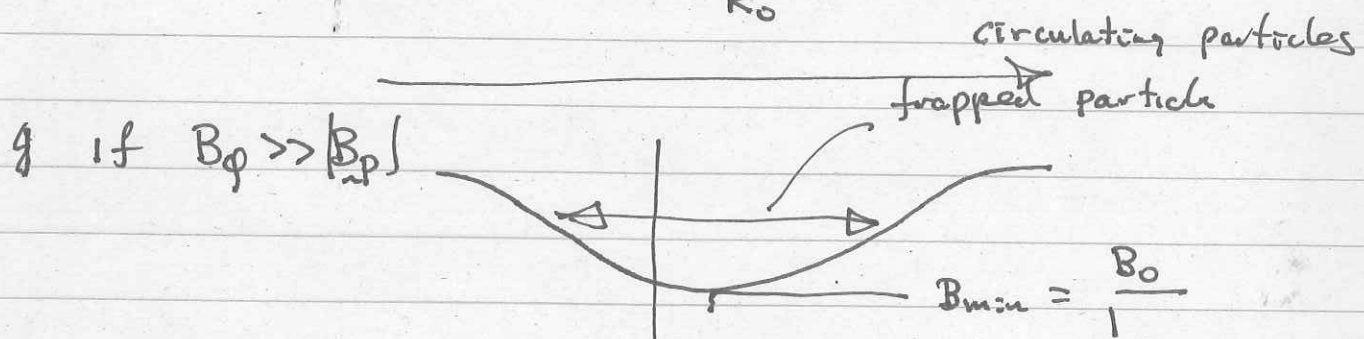
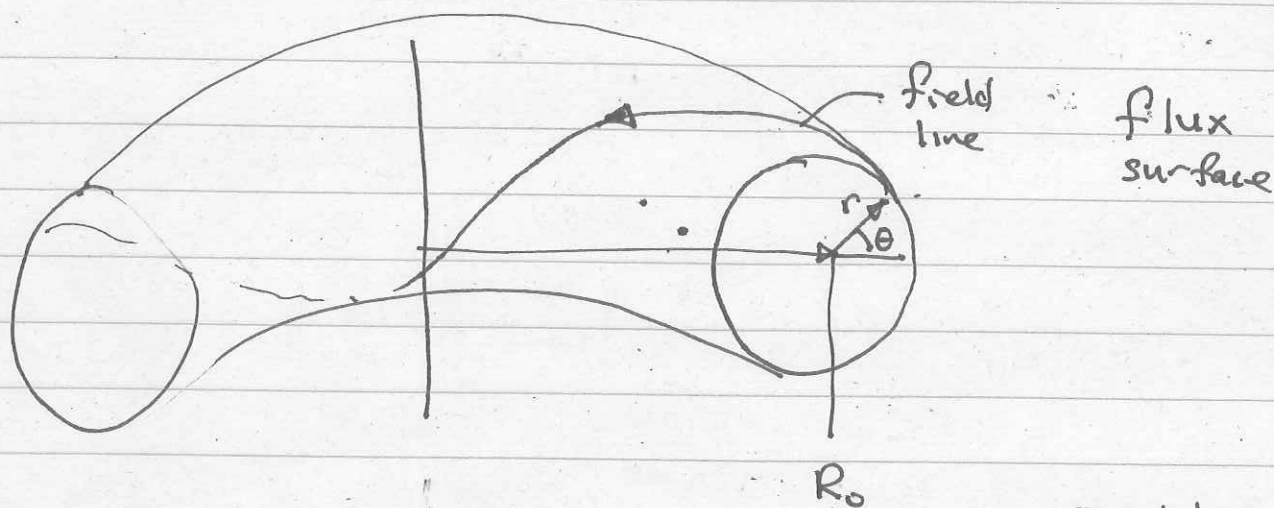


* Particle Orbits in Toroidal Geometry

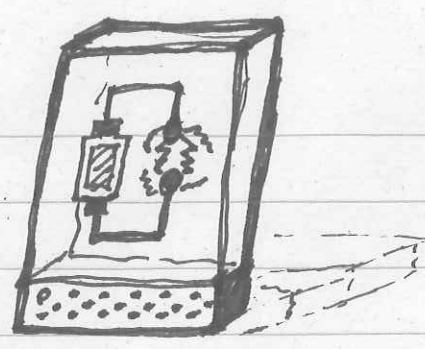


Particles
not confined

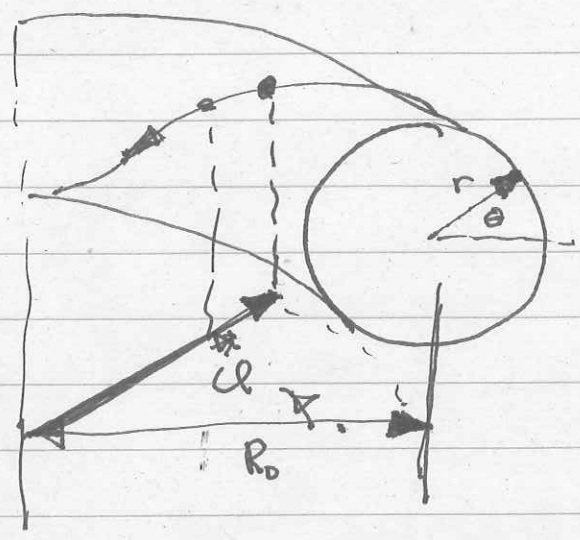
Add Poloidal field



$$|B| = \frac{B_0}{(1 + \frac{r}{R} \cos \theta)} \sim \frac{\text{const}}{R}$$



Safety
factor



Assum $\frac{r}{R_0} \ll 1$ Large aspect ratio
 $R = R_0 + r \cos \theta$

$$B_T \sim \frac{B_0 R_0}{R}$$

$$B_p(r)$$

Following a field line

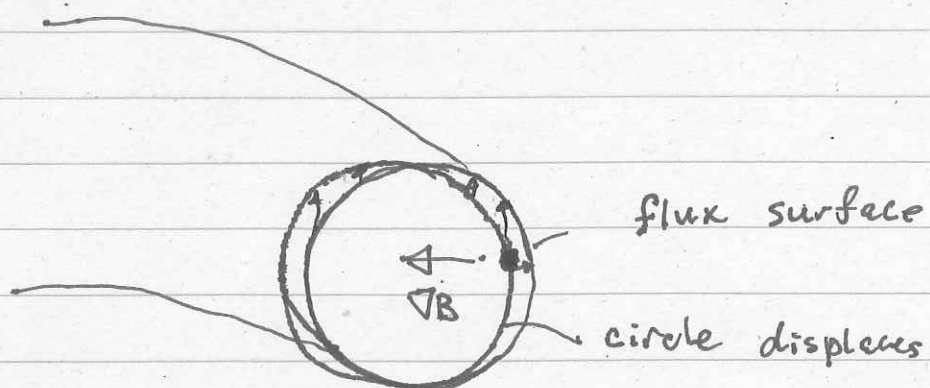
$$\frac{r d\theta}{B_p} = \frac{R d\phi}{B_T}$$

$$\frac{d\phi}{d\theta} = \frac{r B_T}{R B_p} = q(r) \quad \text{safety factor}$$

~~the~~ field line makes q toroidal loops for each poloidal loop.

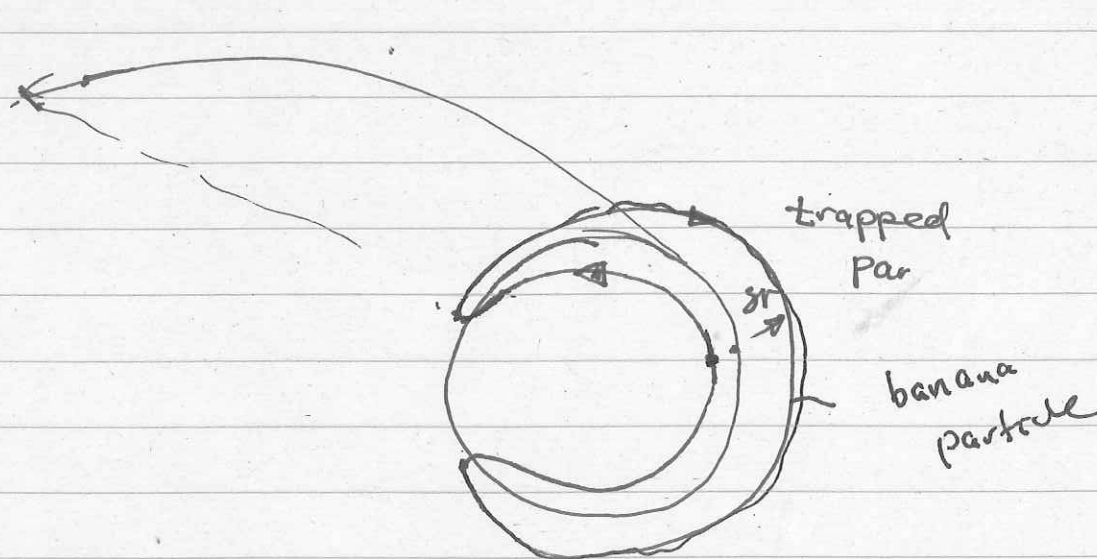
$$q(r) = m/n \quad \text{rational surface}$$

DRIFT ORBITS ARE closed



$$\vec{V} = \frac{c}{qB} \hat{b} \times (\mu B + m v_{\perp}^2) \left(\frac{\nabla B}{B} \right) + v_{\parallel} \hat{b}$$

$\swarrow$ mainly toroidal $\nwarrow$ mainly inward



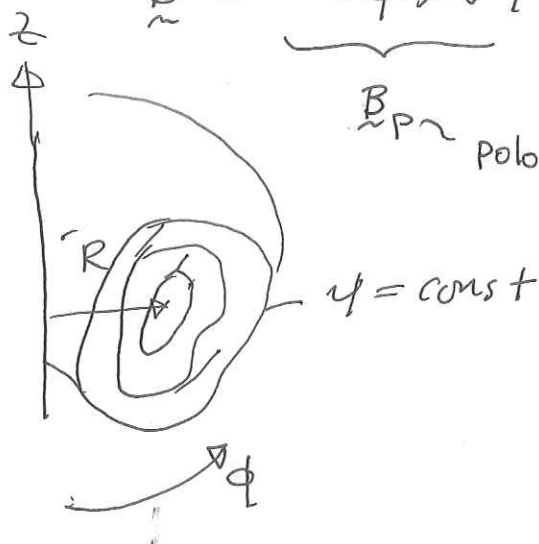
Variation in Radius

$$\delta r = \frac{V_{\perp} - \langle V_{\parallel} \rangle}{(q B_p / m c)}$$

$B_p \ll B$

2D Axisymmetric Field

$$\vec{B} = \underbrace{-\nabla\phi \times \nabla\psi}_{\vec{B}_P \text{ poloidal}} + \underbrace{I(\psi)\nabla\phi}_{\vec{B}_T \text{ toroidal}}$$



$$\vec{B} \cdot \nabla\psi = 0$$

field lines lie in surface of const ψ

Potential Flux Function

$$\nabla\psi \cdot \vec{V}_d = \nabla\psi \cdot \frac{c}{4\pi B} V_{||} \nabla \times (b_m V_{||})$$

$$= \frac{c V_{||}}{4\pi B} \nabla \cdot (b_m V_{||} \times \nabla\psi)$$

After some work

$$= - \frac{c I V_{||}}{4\pi B} b_m \cdot \nabla \left(\frac{V_{||}}{B} \right)$$

Variation in ψ

$$\frac{d\psi}{dt} = - \frac{c I}{4\pi} \frac{d}{dt} \left(\frac{V_{||}}{B} \right)$$

$$\psi = \psi(r_0) + \delta r \frac{\partial \psi}{\partial r}$$

$$\delta r = - \frac{c I}{4 \partial \psi / \partial r} \left(\frac{V_{||}}{B} \right)$$

$$\approx - \frac{c B_T m V_{||}}{B_P B}$$

For

$$V_{||} = \sqrt{\frac{2}{m}(\epsilon - \mu B)}$$

$$\approx \sqrt{\frac{2}{m}(\epsilon - \mu B_0 \frac{1 - \frac{r}{R_0} \cos \theta}{1 - \frac{r}{R_0} \cos \theta})}$$

For circulating

$$V_{||} - \langle V_{||} \rangle \sim \left(\frac{r}{R_0}\right) V_{||}$$

 ϵ

$$\delta r = \frac{V_{||} r}{R_0 B_p / m c} \sim \rho$$

same size as
Larmor radius
(not significant)

For trapped $\langle V_{||} \rangle = 0$

$$V_{||} \sim \left(\frac{r}{R_0}\right)^{1/2} \sqrt{\frac{2\epsilon}{m}}$$

 δr

$$\delta r \gg \rho$$

trapped particles
dominate ~~the~~ classical
transport of Heat/momentum

Neo-Classical Transport

To avoid transport

$$V_{||} \nabla \psi = 0$$

omnigenous