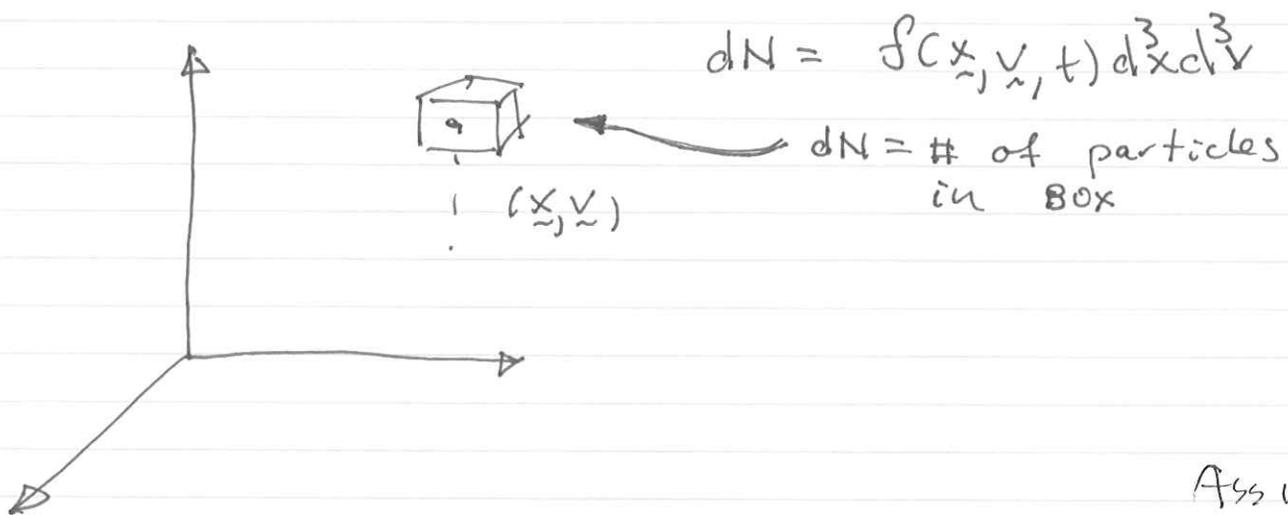


# The Vlasov Equation (collisionless Boltzmann Equation)

Goal: obtain an evolution equation for the distribution function

$$f(\underline{x}, \underline{v}, t)$$

Recall the definition of  $f(\underline{x}, \underline{v}, t)$  <sup>3+1 args</sup>



Answer:

force/m

$$\frac{\partial f}{\partial t} + \underline{v} \cdot \frac{\partial f}{\partial \underline{x}} + \frac{q}{m} \left( \underline{E} + \frac{\underline{v} \times \underline{B}}{c} \right) \cdot \frac{\partial f}{\partial \underline{v}} = 0$$

~~collisionless~~

Comments:

## Newton's Laws &amp; Classical Mechanics

$\vec{x}_n, \vec{v}_n$  are dependent variables

problem is solved when we know

$\vec{x}(t), \vec{v}(t)$  particle description  
 $\uparrow$  independent variable

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Vlasov description

$\vec{x}_n, \vec{v}_n, t$  are all independent variables

problem is solved when we know

$f(\vec{x}_n, \vec{v}_n, t)$

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Fluid equations

$\vec{x}, t$  are independent variables

Problem is solved when we know

$\vec{u}(\vec{x}, t)$	velocity	} dependent variables
$n(\vec{x}, t)$	density	
$T(\vec{x}, t)$	Temperature	

Before we start we should ask the basic question. Does knowledge of  $f(\underline{x}, \underline{v}, t)$  completely describe the system?

Answer: No

Lets say our system consists of  $N$  particles. Then the complete state of the system is given by  $6N$  variables, Three components of position and velocity for each particle. Certainly a lot more information than we can handle.

Even if we adopt a statistical description our pdf would have many ~~arguments~~ arguments

pdf 6N+1 args

$$dP = F_N(\underline{x}_1, \underline{v}_1, \underline{x}_2, \underline{v}_2, \dots, \underline{x}_N, \underline{v}_N, t) d^3x_1 d^3v_1 \dots d^3x_N d^3v_N$$

$dP =$  probability of finding particle #1 in box 1  
and " #2  
and " #N in box N

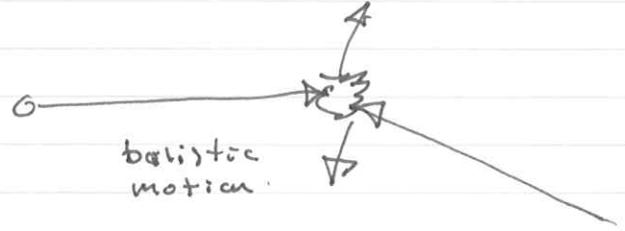
Here we are assuming particles are distinguishable.

Evolution equation for  $F$  is known as the Liouville Equation.

So how do we get away with characterizing the system in terms of ~~a function of~~ only 6 arguments?

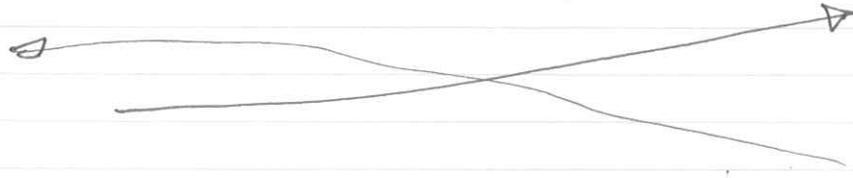
Answer, particles are "weakly" interacting. When does this apply?

dilute gas



infrequent collisions by particles not likely encounter each other again soon

Plasma



Single inter particle reactions have small effect given particle interacts with many at once

Mainly, particles are uncorrelated.

If particles are uncorrelated

$$F_N(\underline{x}_1, \underline{v}_1, \underline{x}_2, \underline{v}_2, \dots, \underline{x}_N, \underline{v}_N) = F_1(\underline{x}_1, \underline{v}_1) F_1(\underline{x}_2, \underline{v}_2) \dots F_1(\underline{x}_N, \underline{v}_N)$$

joint pdf is product of <sup>single particle</sup> pdfs

our distribut  
total # of particles

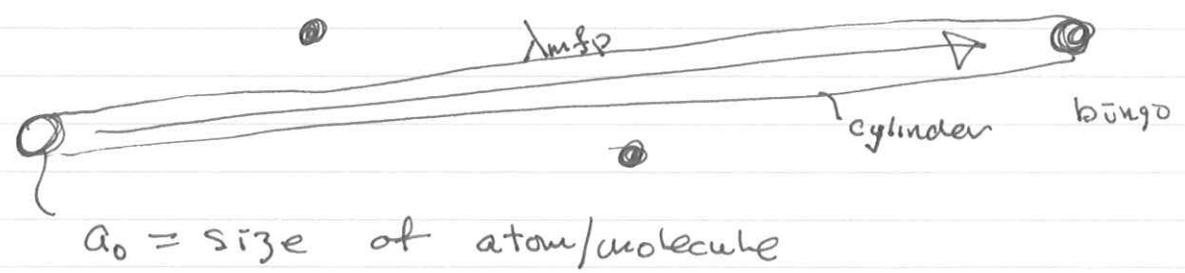
$$f(\underline{x}, \underline{v}) = N F_1(\underline{x}, \underline{v})$$

why

$$\int d^3x d^3v F_1 = 1$$

$$\int d^3x d^3v f = N$$

~~For~~ Justification for a dilute gas



$\lambda_{mfp} =$  how far it goes before collision  
cross-section for collision

$$\text{Volume of cylinder} = a_0^2 \lambda_{mfp}$$

$$n a_0^2 \lambda_{mfp} = 1$$

$$\lambda_{mfp} = \frac{1}{n a_0^2}$$

normalize  $\lambda_{mfp}$  to typical particle spacing  $r_s = n^{-1/3}$

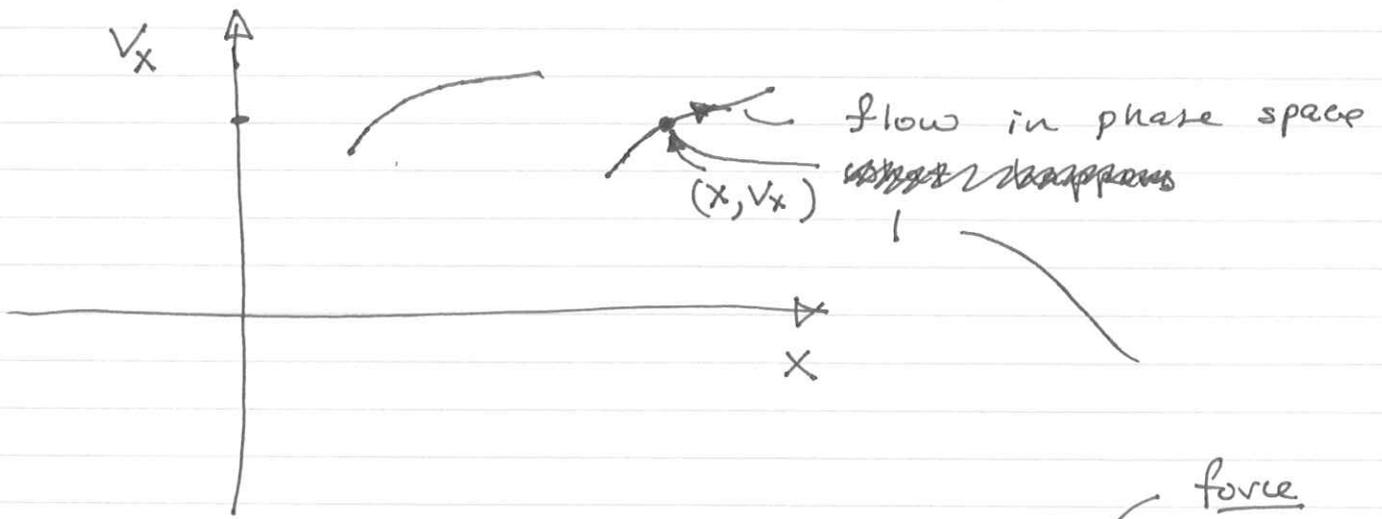
$$\frac{\lambda_{mfp}}{r_s} = \frac{r_s^2}{a_0^2} \frac{1}{r_s^3 n} = \frac{r_s^2}{a_0^2} \gg 1$$

Particle has traveled "long distance" before next collision at which point correlation with past collision is lost.

~~What is~~ Simplification

What is  $\underline{U}$  6D velocity in phase space

Focus on  $x-v_x$  phase space



$$\left. \frac{dx}{dt} \right|_{\text{particelle}} = v_x$$

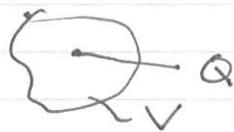
$$\frac{dv_x}{dt} = \frac{F_x}{m}$$

$v_x$  &  $\frac{F_x}{m}$  are ~~the~~  $x$  components of  $\underline{U}$

$v_y$   $\frac{F_y}{m}$  }  $y$  &  $z$  components  
 $v_z$   $\frac{F_z}{m}$  }

Anologous situation, charge conservation  
in 3D

$$\frac{dQ_V}{dt} = - \int_S d\vec{s} \cdot \vec{J}$$



$$Q_V = \int_V q n d^3x \quad \begin{array}{l} \text{number density} \\ \text{"fluid velocity"} \end{array}$$

$$\vec{J} = q n \vec{u}(x, t)$$

$$\frac{d}{dt} \int_V d^3x n = - \int_S d\vec{s} \cdot \vec{u} n$$

use ~~the~~ Divergence theorem

$$\int_S d\vec{s} \cdot \vec{u} n = \int_V d^3x \nabla \cdot (\vec{u} n)$$

~~TRUE FOR ANY~~ statis

$$\int_V d^3x \left( \frac{\partial n}{\partial t} + \nabla \cdot (\vec{u} n) \right) = 0$$

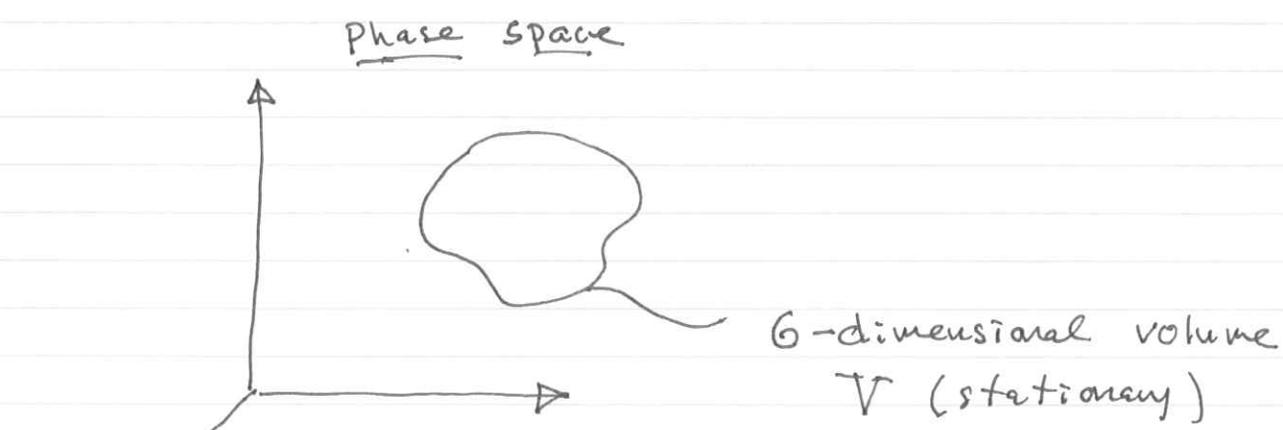
TRUE FOR ANY V

$$\therefore \frac{\partial n}{\partial t} + \nabla \cdot (\vec{u} n) = 0$$

Now lets derive the Vlasov Equation

Obtain the derivation by analogy

Conservation of # of particles



$$N_V = \int d^3x \int_V d^3v f(x, v, t)$$

$N_V = \#$  of particles in  $V$

$\frac{dN}{dt} = -$  Rate at which particles leave volume by passing through surface

$$= - \int dS_m \cdot \underline{U} f$$

↳ 6D velocity in phase space

$dS_m =$  5D surface area direction out of  $V$

Also by divergence theorem

$$\int_S d\vec{s} \cdot \vec{U} f = \int_V d^3x d^3v \nabla_6 \cdot (\vec{U} f)$$

$\left\{ \begin{array}{l} \text{six dimensional} \\ \text{divergence} \end{array} \right.$

$$\nabla_6 \cdot (\vec{U} f) = \frac{\partial}{\partial x} (v_x f) + \frac{\partial}{\partial v_x} \left( \frac{F_x}{m} f \right) + \frac{\partial}{\partial y} + \dots$$

Continuity equation in phase space

~~$$\frac{\partial f}{\partial t} \int_V d^3x d^3v f = \int_V d^3x d^3v \frac{\partial f}{\partial t}$$~~

$$\int d^3x d^3v \left[ \frac{\partial f}{\partial t} + \nabla_6 \cdot (\vec{U} f) \right] = 0$$

$$\frac{\partial f}{\partial t} + \nabla_6 \cdot (\vec{U} f) = 0$$

OR

$$\frac{\partial f}{\partial t} + \cancel{\nabla_6 \cdot (\vec{U} f)} \frac{\partial}{\partial x} \cdot (\vec{v} f) + \frac{\partial}{\partial v_x} \cdot \left( \frac{F_x}{m} f \right) = 0$$

NOT THERE YET

$$\frac{\partial}{\partial x} \cdot (\underline{v} f) = \underline{v} \cdot \frac{\partial f}{\partial x} + f \frac{\partial}{\partial x} \cdot \underline{v}$$

What is  $\frac{\partial}{\partial x} \cdot \underline{v}$       Ans 0

$\underline{v}$  and  $\underline{x}$  are both independent variables

— fluid velocity  
 $(\frac{\partial}{\partial x} \cdot \underline{u}(x) \neq 0)$

$$\frac{\partial}{\partial \underline{v}} \cdot \left( \frac{\underline{v} \cdot \underline{F}}{m} f \right)$$

where

$$\underline{F} = \underline{q} \left( \underline{E}(\underline{x}, t) + \frac{\underline{v} \times \underline{B}(\underline{x}, t)}{c} \right)$$

$$= \frac{\underline{v} \cdot \underline{F}}{m} \cdot \frac{\partial f}{\partial \underline{v}} + \frac{f}{m} \frac{\partial}{\partial \underline{v}} \cdot \underline{F}$$

$$\frac{\partial}{\partial \underline{v}} \cdot \underline{F} = \frac{\partial}{\partial \underline{v}} \cdot \underline{q} \frac{\underline{v} \times \underline{B}}{c}$$

$$= \frac{\partial}{\partial v_x} \frac{q}{c} (v_y B_z - v_z B_y) + \frac{\partial}{\partial v_y} \frac{q}{c} (v_z B_x - v_x B_z)$$

$$+ \frac{\partial}{\partial v_z} \frac{q}{c} (v_x B_y - v_y B_x) = 0$$

Interaction with "smooth" density and currents" effects of particle discreteness absent

Examples of missing effects

- 1) collisions
- 2) spontaneous radiation  
Bremsstrahlung  
cyclotron radiation



- 1) stimulated radiation

Flow is incompressible in phase space!

$$\nabla_6 \cdot \underline{U} = \frac{\partial}{\partial \underline{x}} \cdot \underline{v} + \frac{\partial}{\partial \underline{v}} \cdot \frac{q}{m} \underline{F} = 0$$

Important consequences

Q: How do particles influence each other

A: Through E&M Forces due to average charge and current distributions

$$\underline{F} = q \left( \underline{E} + \underline{v} \times \underline{B} \right)$$

$$\nabla \cdot \underline{E} = 4\pi \rho = 4\pi \sum_q q \int n_q(\underline{x}, t)$$

$$n_q(\underline{x}, t) = \int d^3v f_q(\underline{x}, \underline{v}, t)$$

$$\nabla \times \underline{B} = \frac{4\pi}{c} \underline{j} + \frac{1}{c} \frac{\partial \underline{E}}{\partial t}$$

$$\underline{j} = \sum_q q n_q \underline{u}_q = \sum_q q \int d^3v \underline{v} f_q$$