

Dielectric Waveguides

ENEE 681

Basic principles

Slab model

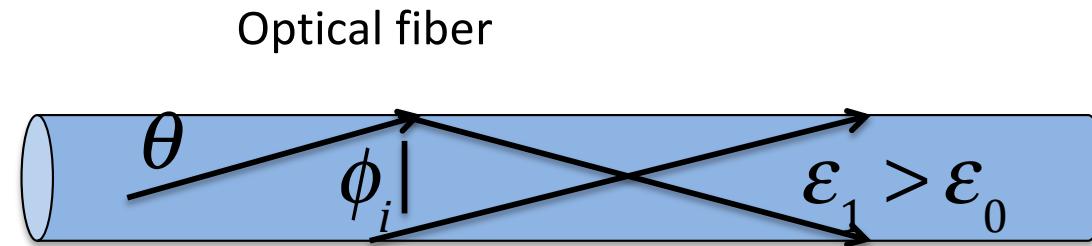
two regions

boundary conditions

dispersion relation – odd and even modes

Cylindrical geometry

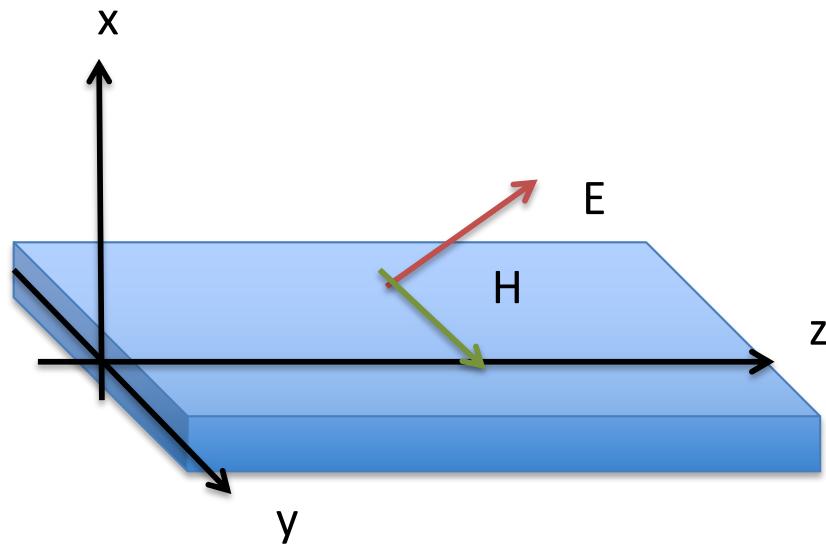
Guiding by total internal reflection



Total reflection if

$$\sin \phi_i = \cos \theta > \sqrt{\varepsilon_0 / \varepsilon_1}$$

Slab Model

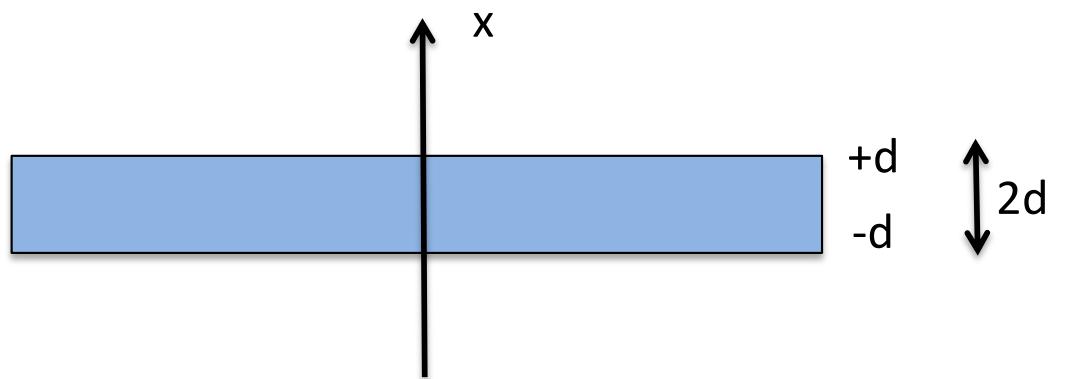


$$\epsilon = \begin{cases} \epsilon_1 & |x| < d \\ \epsilon_0 & |x| > d \end{cases}$$

Look for solutions

$$\mathbf{E} = \text{Re} \left\{ (\hat{E}_x, 0, \hat{E}_z) \exp \left[i(k_z z - \omega t) \right] \right\}$$

$$\mathbf{H} = \text{Re} \left\{ (0, \hat{H}_y, 0) \exp \left[i(k_z z - \omega t) \right] \right\}$$



Maxwell's Equations

$$\nabla \times \vec{\mathbf{E}} = -\mu \frac{\partial \vec{\mathbf{H}}}{\partial t}$$

$$\frac{\partial \hat{E}_z}{\partial y} - ik_z \hat{E}_y = i\omega \mu \hat{H}_x$$

$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega \mu \hat{H}_y$$

$$\frac{\partial \hat{E}_y}{\partial x} - \frac{\partial \hat{E}_x}{\partial y} = i\omega \mu \hat{H}_z$$

$$\nabla \times \vec{\mathbf{H}} = \epsilon \frac{\partial \vec{\mathbf{E}}}{\partial t}$$

$$\frac{\partial \hat{H}_z}{\partial y} - ik_z \hat{H}_y = -i\omega \epsilon \hat{E}_x$$

$$ik_z \hat{H}_x - \frac{\partial \hat{H}_z}{\partial x} = -i\omega \epsilon \hat{E}_y$$

$$\frac{\partial \hat{H}_y}{\partial x} - \frac{\partial \hat{H}_x}{\partial y} = -i\omega \epsilon \hat{E}_z$$

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Maxwell's Equations

$$\nabla \times \vec{\mathbf{E}} = -\mu \frac{\partial \vec{\mathbf{H}}}{\partial t}$$

$$\frac{\partial \hat{E}_z}{\partial y} - ik_z \hat{E}_y = i\omega \mu \hat{H}_x \rightarrow$$

$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega \mu \hat{H}_y$$

$$\frac{\partial \hat{E}_y}{\partial x} - \frac{\partial \hat{E}_x}{\partial y} = i\omega \mu \hat{H}_z \rightarrow$$

$$\nabla \times \vec{\mathbf{H}} = \epsilon \frac{\partial \vec{\mathbf{E}}}{\partial t}$$

$$\frac{\partial \hat{H}_z}{\partial y} - ik_z \hat{H}_y = -i\omega \epsilon \hat{E}_x$$

$$ik_z \hat{H}_x - \frac{\partial \hat{H}_z}{\partial x} = -i\omega \epsilon \hat{E}_y \rightarrow$$

$$\frac{\partial \hat{H}_y}{\partial x} - \frac{\partial \hat{H}_x}{\partial y} = -i\omega \epsilon \hat{E}_z$$

Look for solutions

$$\mathbf{E} = \text{Re} \left\{ (\hat{E}_x, 0, \hat{E}_z) \exp \left[i(k_z z - \omega t) \right] \right\}$$

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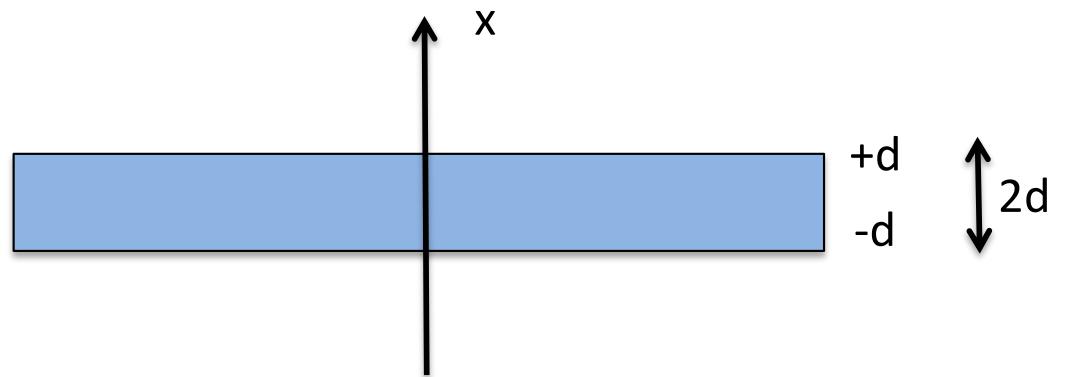
Maxwell's Equations

$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega \mu \hat{H}_y$$

$$\begin{aligned}-ik_z \hat{H}_y &= -i\omega \epsilon \hat{E}_x \\ \frac{\partial \hat{H}_y}{\partial x} &= -i\omega \epsilon \hat{E}_z\end{aligned}$$

$$\epsilon = \begin{cases} \epsilon_1 & |x| < d \\ \epsilon_0 & |x| > d \end{cases}$$

What are the boundary conditions at $x = +/- d$?



Maxwell's Equations

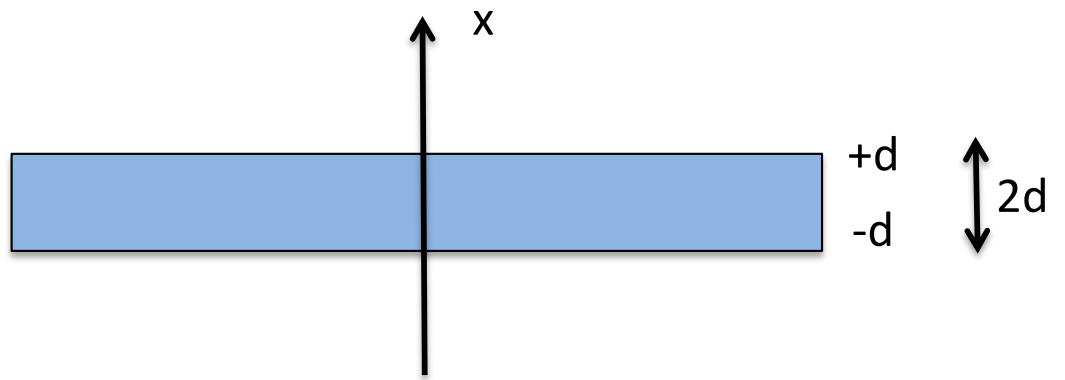
$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega \mu \hat{H}_y$$

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$$\varepsilon = \begin{cases} \varepsilon_1 & |x| < d \\ \varepsilon_0 & |x| > d \end{cases}$$

What are the boundary conditions at $x = +/- d$?

$$\left. \begin{array}{l} \hat{E}_z \\ \hat{H}_y \end{array} \right\} \text{continuous at } x = \pm d$$



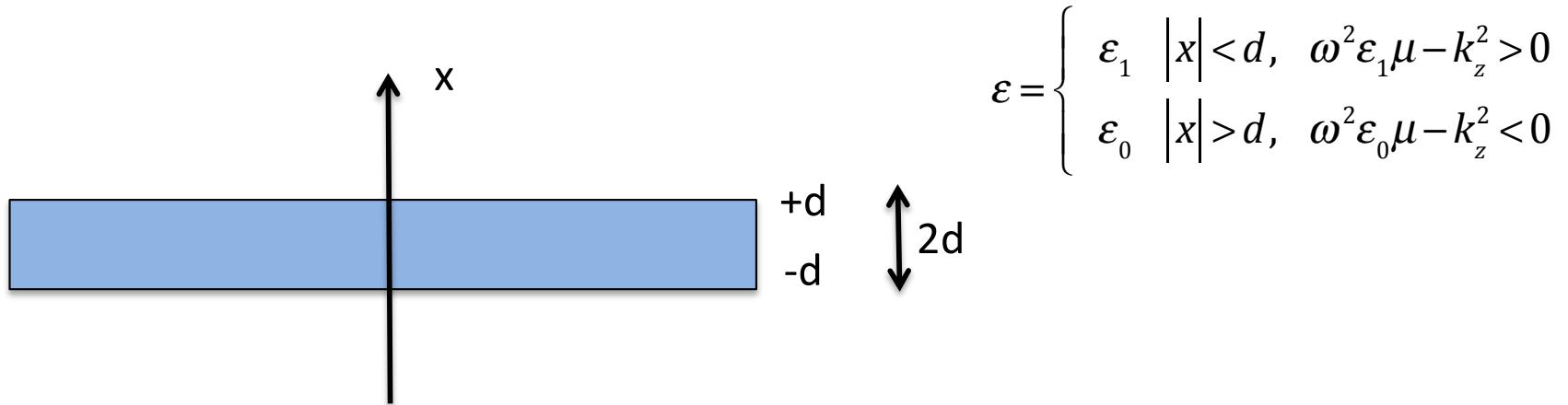
Combine to a single equation

$$\begin{aligned}
 ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} &= i\omega\mu \hat{H}_y \\
 -ik_z \hat{H}_y &= -i\omega\epsilon \hat{E}_x \\
 \frac{\partial \hat{H}_y}{\partial x} &= -i\omega\epsilon \hat{E}_z
 \end{aligned}
 \quad \rightarrow \quad
 \begin{aligned}
 k_z^2 \hat{E}_x + ik_z \frac{\partial \hat{E}_z}{\partial x} &= i\omega\mu (-ik_z \hat{H}_y) = \omega^2 \epsilon \mu \hat{E}_x \\
 \hat{E}_x &= \frac{ik_z}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x} \\
 \hat{H}_y &= \frac{i\omega\epsilon}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x} \\
 \frac{\partial \hat{H}_y}{\partial x} &= -i\omega\epsilon \hat{E}_z
 \end{aligned}
 \quad \rightarrow \quad
 \frac{\partial}{\partial x} \frac{\epsilon}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x} = -\epsilon \hat{E}_z$$

If ϵ is constant $\frac{\partial}{\partial x} \frac{\partial \hat{E}_z}{\partial x} + (\omega^2 \epsilon \mu - k_z^2) \hat{E}_z = 0$

Solve in two regions

If ϵ is constant $\frac{\partial}{\partial x} \frac{\partial \hat{E}_z}{\partial x} + (\omega^2 \epsilon \mu - k_z^2) \hat{E}_z = 0$

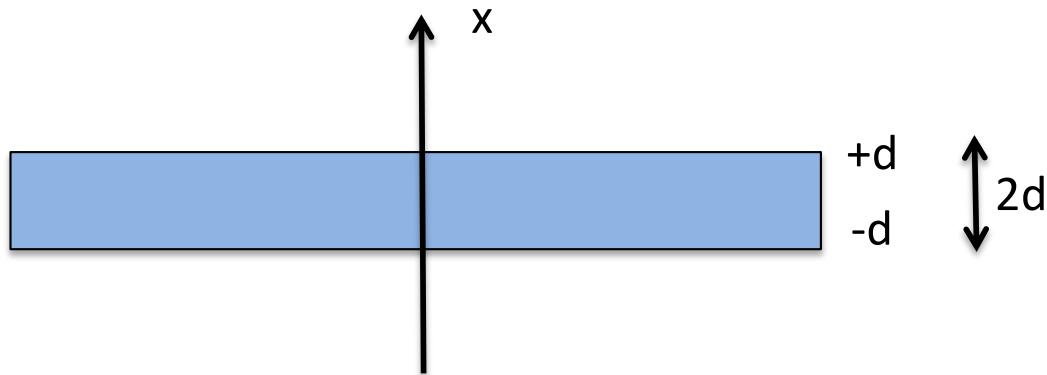


For $x > d$, $\epsilon = \epsilon_0$ $\hat{E}_z = A \exp(-\kappa x)$ $\kappa = \sqrt{k_z^2 - \omega^2 \epsilon_0 \mu}$

Evanescence (Spatially decaying) field

Solve in two regions

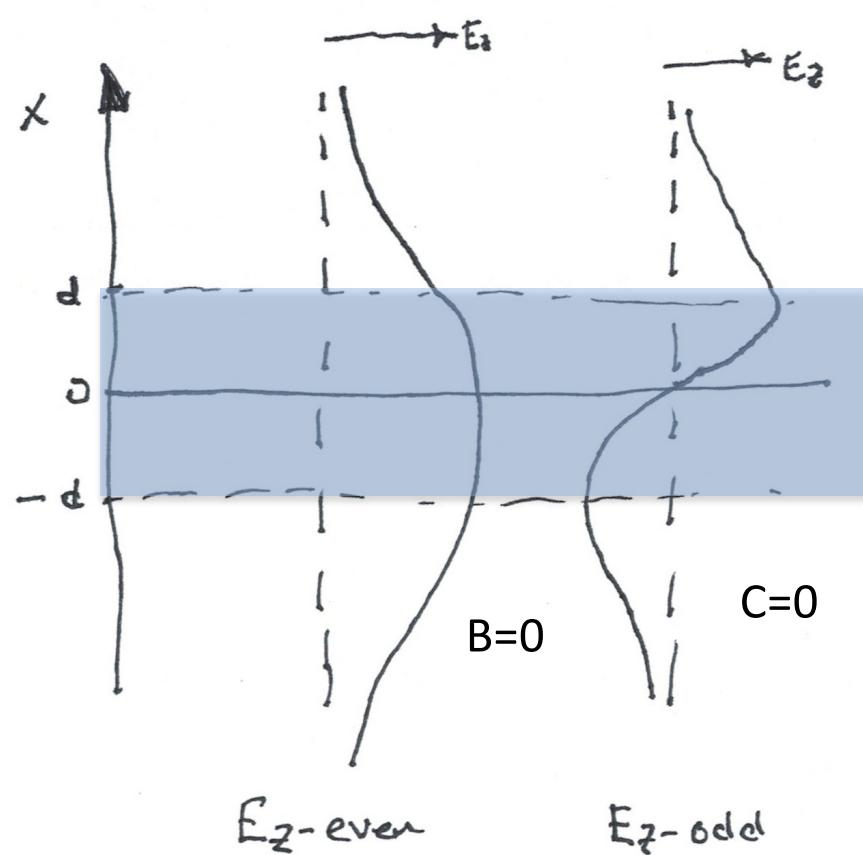
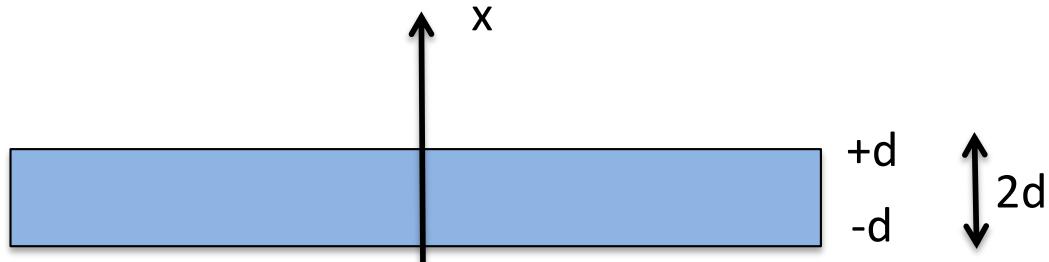
$$\text{If } \epsilon \text{ is constant} \quad \frac{\partial}{\partial x} \frac{\partial \hat{E}_z}{\partial x} + (\omega^2 \epsilon \mu - k_z^2) \hat{E}_z = 0$$



$$\text{For } |x| < d, \quad \epsilon = \epsilon_1 \quad \hat{E}_z = B \sin(k_x x) + C \cos(k_x x) \quad k_x = \sqrt{\omega^2 \epsilon_1 \mu - k_z^2}$$

Spatially oscillating field

Odd and Even solutions



For $|x| < d$, $\epsilon = \epsilon_1$

$$\hat{E}_z = B \sin(k_x x) + C \cos(k_x x)$$

$$k_x = \sqrt{\omega^2 \epsilon_1 \mu - k_z^2}$$

Either: $B=0, C \neq 0$
 or $B \neq 0, C=0$

$$\hat{E}_x = \frac{ik_z}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

$$\hat{H}_y = \frac{i\omega \epsilon}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

E_z –even (E_x – odd) solution

For $|x| < d$, $\hat{E}_z = C \cos(k_x x)$ $k_x = \sqrt{\omega^2 \epsilon_1 \mu - k_z^2}$

$$\hat{E}_x = \frac{ik_z}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

For $x > d$, $\hat{E}_z = A \exp(-\kappa x)$ $\kappa = \sqrt{k_z^2 - \omega^2 \epsilon_0 \mu}$

$$\hat{H}_y = \frac{i\omega \epsilon}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

At $x=d$ $\hat{E}_z = C \cos(k_x d) = A \exp(-\kappa d)$

$$\hat{H}_y = \frac{i\omega \epsilon_1}{k_x^2} \frac{\partial}{\partial x} C \cos(k_x x) \Big|_{x=d} = -\frac{i\omega \epsilon_1}{k_x} C \sin(k_x d)$$

$$\hat{H}_y = \frac{i\omega \epsilon_0}{-\kappa^2} \frac{\partial}{\partial x} A \exp(-\kappa x) \Big|_{x=d} = \frac{i\omega \epsilon_0}{\kappa} A \exp(-\kappa d)$$

Dispersion Relation

$$C \cos(k_x d) = A \exp(-\kappa d)$$

$$-\frac{\epsilon_1}{k_x} C \sin(k_x d) = \frac{\epsilon_0}{\kappa} A \exp(-\kappa d)$$

Even E_z

$$-\frac{\epsilon_1}{\epsilon_0} \kappa d = k_x d \cot(k_x d)$$

Odd E_z

$$\frac{\epsilon_1}{\epsilon_0} \kappa d = k_x d \tan(k_x d)$$

$$\kappa d = d \sqrt{k_z^2 - \omega^2 \epsilon_0 \mu} = \sqrt{\omega^2 (\epsilon_1 - \epsilon_0) \mu d^2 - k_x^2 d^2}$$

$$\kappa d = \sqrt{\omega^2 (\epsilon_1 - \epsilon_0) \mu d^2 - k_x^2 d^2}$$

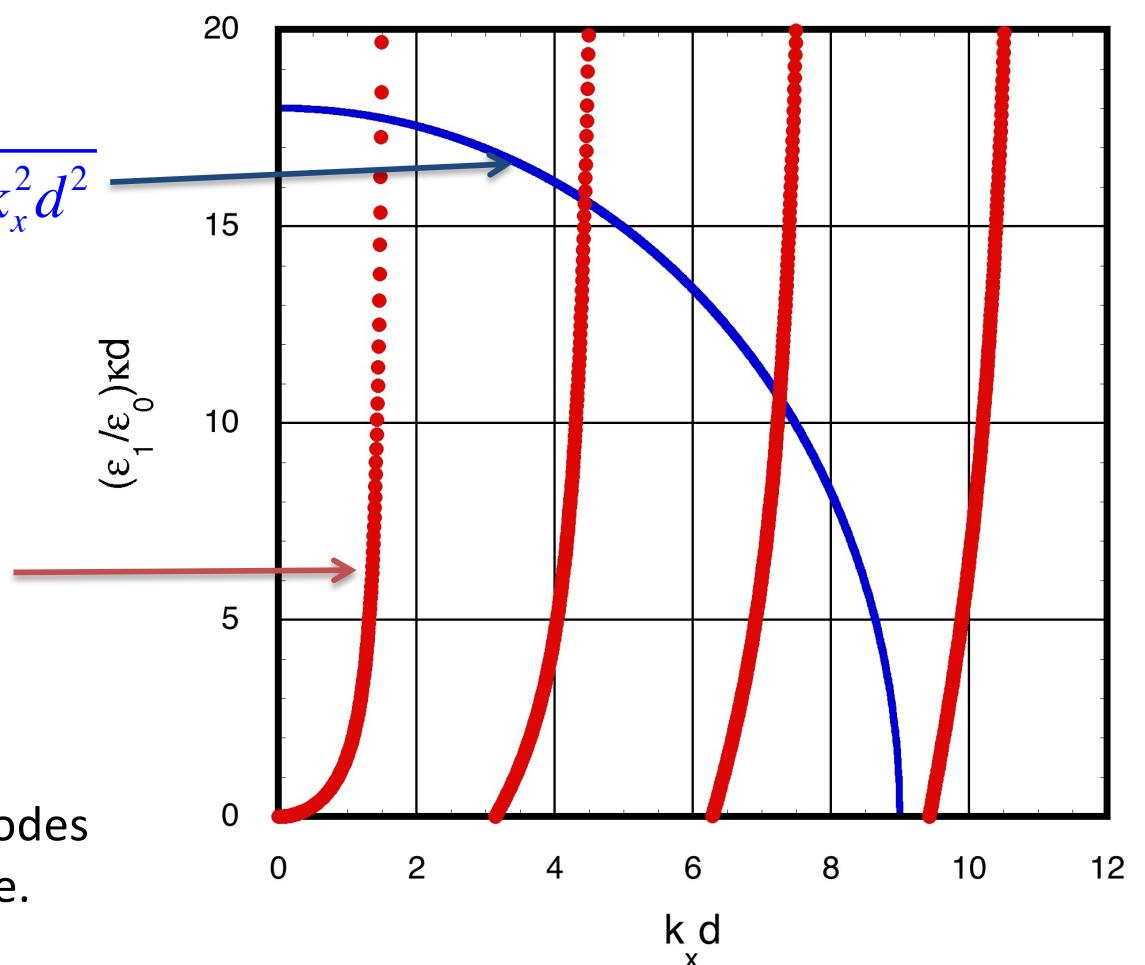
Plot both sides versus $k_x d$

Odd E_z

$$\frac{\epsilon_1}{\epsilon_0} \kappa d = \frac{\epsilon_1}{\epsilon_0} \sqrt{\omega^2 (\epsilon_1 - \epsilon_0) \mu d^2 - k_x^2 d^2}$$

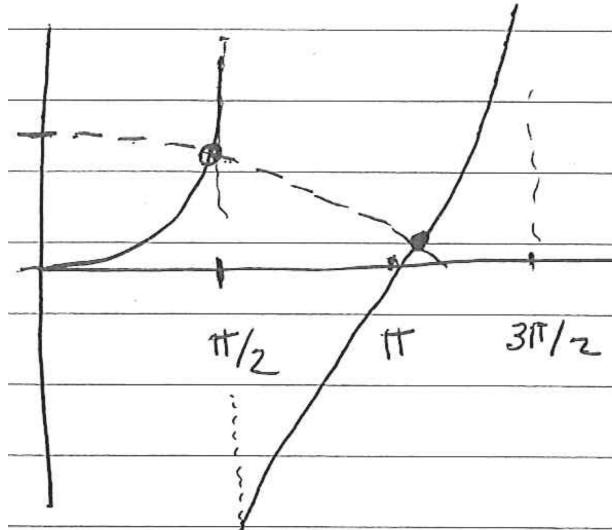
$$\frac{\epsilon_1}{\epsilon_0} \kappa d = k_x d \tan(k_x d)$$

$$\frac{\epsilon_1}{\epsilon_0} \kappa d = k_x d \tan(k_x d)$$



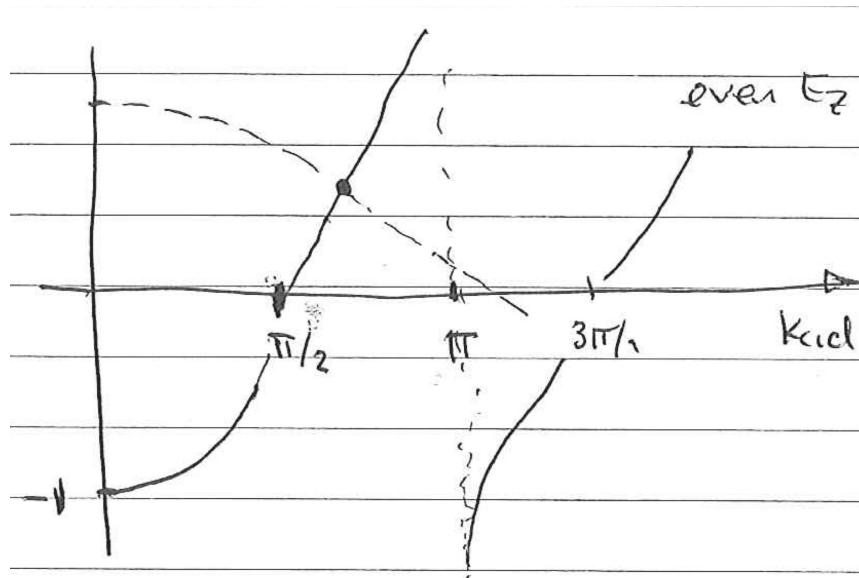
In this case three solutions – modes
There will always be at least one.

Comments



Odd E_z –Even E_x

Always at least one solution
For small diameter d fields outside dielectric
Requires cladding

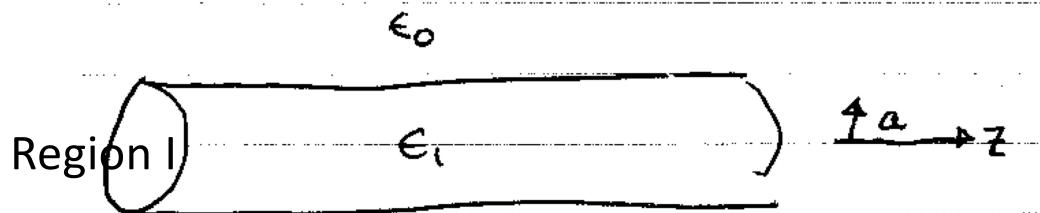


Even E_z –Odd E_x

Possibly no solutions

Cylindrical Dielectrics

Region II



$$\text{Field} \sim \{\hat{E}, \hat{H}(r_\theta)\} \exp[ik_z z - i\omega t]$$

Express transverse components in terms of axial components

$$\hat{E}_1 = \frac{1}{k_1^2} [ik_z \nabla_1 \hat{E}_z - i\omega \mu \epsilon_2 \times \nabla_1 \hat{H}_z]$$

$$\hat{H}_1 = \frac{1}{k_1^2} [ik_z \nabla_1 \hat{H}_z + i\omega \epsilon \epsilon_2 \times \nabla_1 \hat{E}_z]$$

$$k_1^2 = \omega^2 \epsilon \mu - k_z^2 \quad - \text{different in each region}$$

Equation for Axial Components

$$\nabla^2 \hat{E}_z + k_1^2 \hat{E}_z = 0 \quad \nabla^2 \hat{H}_z + k_2^2 \hat{H}_z = 0$$

$$k_1^2 = \omega^2 \epsilon \mu - k_z^2 \quad \text{different in each region}$$

Region I: $k_1^2 = \omega^2 \epsilon \mu - k_z^2 > 0$

Region II: $k_{12}^2 = \omega^2 \epsilon_2 \mu - k_z^2 = -\gamma_2^2 < 0$

$$\hat{E}_z = A_1 \exp(i\ell\theta) J_\ell(k_{1r}r)$$

$\ell = \text{integer}$

$$\hat{E}_z(r_1\theta) = A_2 e^{i\ell\theta} K_\ell(\gamma_2 r)$$

$$\hat{H}_z = B_1 \exp(i\ell\theta) J_\ell(k_{2r}r)$$

$$\hat{H}_z(r_2\theta) = B_2 e^{i\ell\theta} K_\ell(\gamma_2 r)$$

Ordinary Bessel functions
Regular at origin

Modified Bessel functions
Goes to zero at infinity

Unknowns and Boundary Conditions

$$\hat{E}_z = A_1 \exp(i\ell\theta) J_\ell(k_1 r)$$

$$\hat{E}_z(r, \theta) = A_2 e^{i\ell\theta} K_\ell(\gamma_2 r)$$

$$\hat{H}_z = B_1 \exp(i\ell\theta) J_\ell(k_1 r)$$

$$\hat{H}_z(r, \theta) = B_2 e^{i\ell\theta} K_\ell(\gamma_2 r)$$

Unknowns A_1, A_2, B_1, B_2 & ω for given k_z

Boundary Conditions : at $r=a$

Tangential components of E and H continuous

$E_z, E_\theta, H_z, H_\theta$ - continuous at $r=a$

H_z, E_z - continuous

$$A_1 J_\ell(k_1 a) = A_2 K_\ell(\gamma_2 a)$$

$$B_1 J_\ell(k_1 a) = B_2 K_\ell(\gamma_2 a)$$

Continuity of Transverse Components

Transverse components in terms
of axial components

$$\hat{E}_z = \frac{1}{k_z} \left[i k_z \nabla_{\perp} \hat{E}_z - i \omega \mu_0 \epsilon_0 \times \nabla_{\perp} \hat{H}_z \right]$$

$$\hat{H}_z = \frac{1}{k_z} \left[i k_z \nabla_{\perp} \hat{H}_z + i \omega \epsilon_0 \epsilon_0 \times \nabla_{\perp} \hat{E}_z \right]$$

Region I

$$\hat{E}_{\perp}(r) = \frac{1}{k_{11}^2} \left\{ i k_z \left(\frac{i\ell}{a} \right) A_1 J_{\ell}(k_{11}a) - i \omega \mu_0 B_1 k_{11} J_{\ell}'(k_{11}a) \right\}$$

$$H_{\perp}(r=a) = \frac{1}{k_{11}^2} \left\{ i k_z \left(\frac{i\ell}{a} \right) B_1 J_{\ell}(k_{11}a) + i \omega \epsilon_0 k_{11} A_1 J_{\ell}'(k_{11}a) \right\}$$

Region II

$$\hat{E}_0(a) = -\frac{1}{\gamma_2^2} \left\{ i k_z \left(\frac{i\ell}{a} \right) A_2 K_{\ell}(x_2 a) - i \omega \mu_0 \gamma_2 B_2 K_{\ell}'(x_2 a) \right\}$$

$$k_z^2 = -\gamma_2^2$$

$$H_0(a) = -\frac{1}{\gamma_2^2} \left\{ i k_z \left(\frac{i\ell}{a} \right) B_2 K_{\ell}(x_2 a) + i \omega \epsilon_0 \gamma_2 A_2 K_{\ell}'(x_2 a) \right\}$$

Solution

$$A_1 J_e(k_1 a) = A_2 K_e(\gamma_2 a)$$

$$B_1 J_e(k_1 a) = B_2 K_e(\gamma_2 a)$$



$$\hat{E}_0(a) = \frac{1}{k_{21}^2} \left\{ i k_2 \left(\frac{i\ell}{a} \right) A_1 J_e(k_1 a) - i w \mu B_1 K_e J_e'(k_1 a) \right\}$$

$$H_0(v=a) = \frac{1}{k_{21}^2} \left\{ i k_2 \left(\frac{i\ell}{a} \right) B_1 J_e(k_1 a) + i w \epsilon k_2 A_1 J_e'(k_1 a) \right\}$$

$$\hat{E}_0(a) = -\frac{1}{\gamma_2^2} \left\{ i k_2 \left(\frac{i\ell}{a} \right) A_2 K_e(\gamma_2 a) - i w \mu \gamma_2 B_2 K_e J_e'(k_1 a) \right\}$$

$$H_0(a) = -\frac{1}{\gamma_2^2} \left\{ i k_2 \left(\frac{i\ell}{a} \right) B_2 K_e(\gamma_2 a) + i w \epsilon \gamma_2 A_2 K_e J_e'(k_1 a) \right\}$$

Replace $A_2 = A_1 J_e / K_e$ $B_2 = B_1 J_e / K_e$

$$\begin{bmatrix} 2 \times 2 \\ M \end{bmatrix} \begin{pmatrix} A_1 \\ B_1 \end{pmatrix} = 0$$

solution $\det(M) = 0$ determines $w(k)$