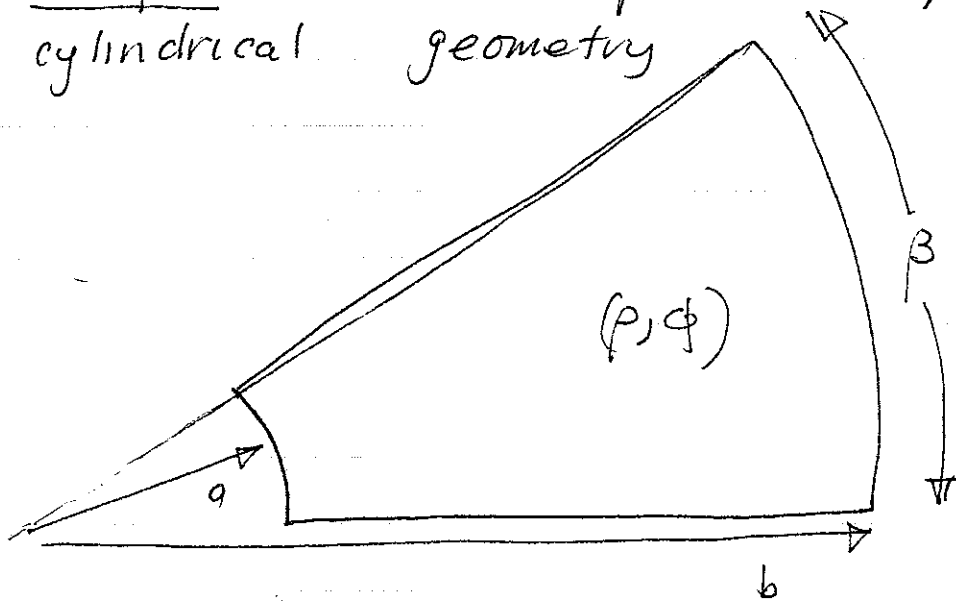


Separation in Cylindrical Coord

67

Examples of Laplace's Equation in cylindrical geometry



Separation $\Phi(\rho, \phi) = R(\rho)\psi(\phi)$

$$\nabla^2 \Phi = 0$$

Laplace's Equation

$$\nabla^2 \Phi = \frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial \Phi}{\partial \rho} + \frac{1}{\rho^2} \frac{\partial^2 \Phi}{\partial \phi^2}$$

$$= \frac{R\psi}{\rho^2} \left\{ \underbrace{\frac{\rho}{R} \frac{\partial}{\partial \rho} \rho \frac{\partial R}{\partial \rho}}_{\text{separation constant } \nu^2} + \underbrace{\frac{1}{\psi} \frac{\partial^2 \psi}{\partial \phi^2}}_{-\nu^2} \right\}$$

separation constant ν^2 $-\nu^2$

ν^2 is arbitrary

for a particular $\nu^2 > 0$

~~particular~~ solutions can be written

$$\psi = C \exp(i\nu\varphi) \quad \text{or} \quad A \sin \nu\varphi + B \cos \nu\varphi$$

$$R = \cancel{A_1 \rho^\nu} \quad a \rho^\nu + b \rho^{-\nu}$$

~~general so~~

$$\text{IF } \nu^2 = 0$$

$$\frac{\partial^2 \psi}{\partial \varphi^2} = 0$$

$$\psi = A_0 + B_0 \varphi$$

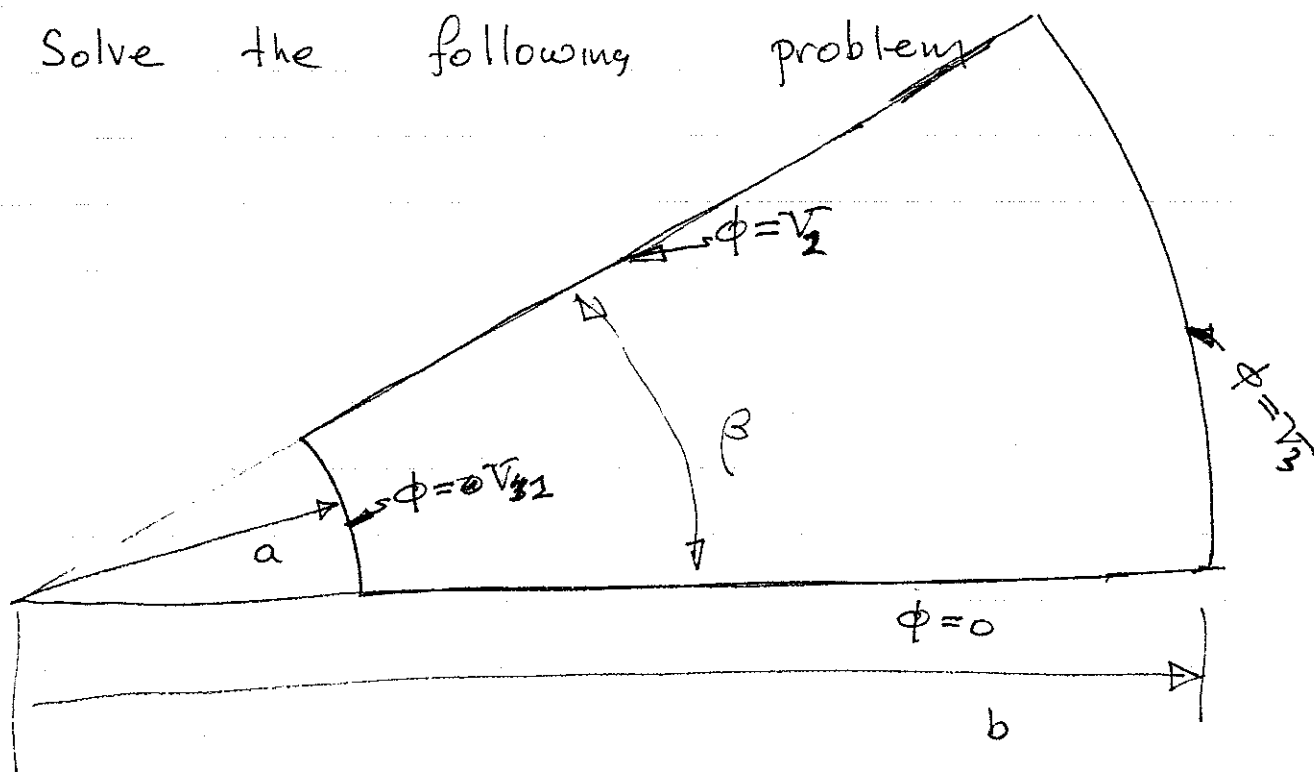
arbitrary constants

$$\rho \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} R = 0 \quad R = a_0 + b_0 \ln \rho$$

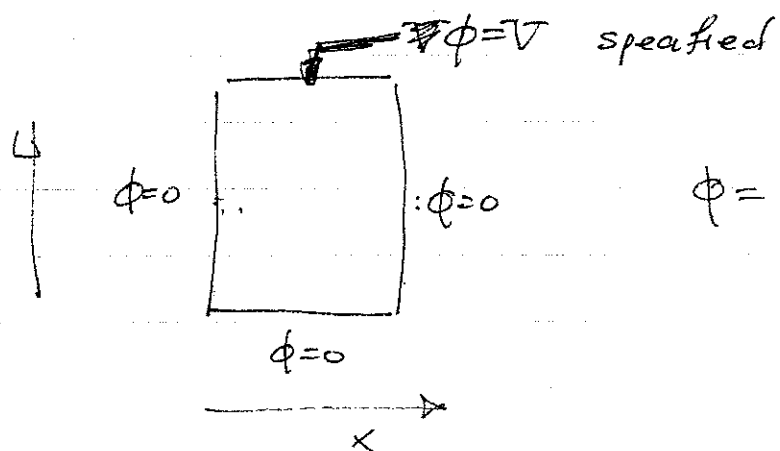
if region of solution corresponds to full range of φ such that

$$\psi(\varphi + 2\pi) = \psi(\varphi) \quad (\text{periodic}) \quad \nu = \text{integer}$$

Solve the following problem



Similar to our problem



$$\phi = \sum_n a_n \sin\left(\frac{n\pi x}{a}\right) \sinh \frac{n\pi y}{a}$$

~~$$\phi = B_0 r$$~~

start with a ~~separa~~ solution corresponding to zero separation constant,

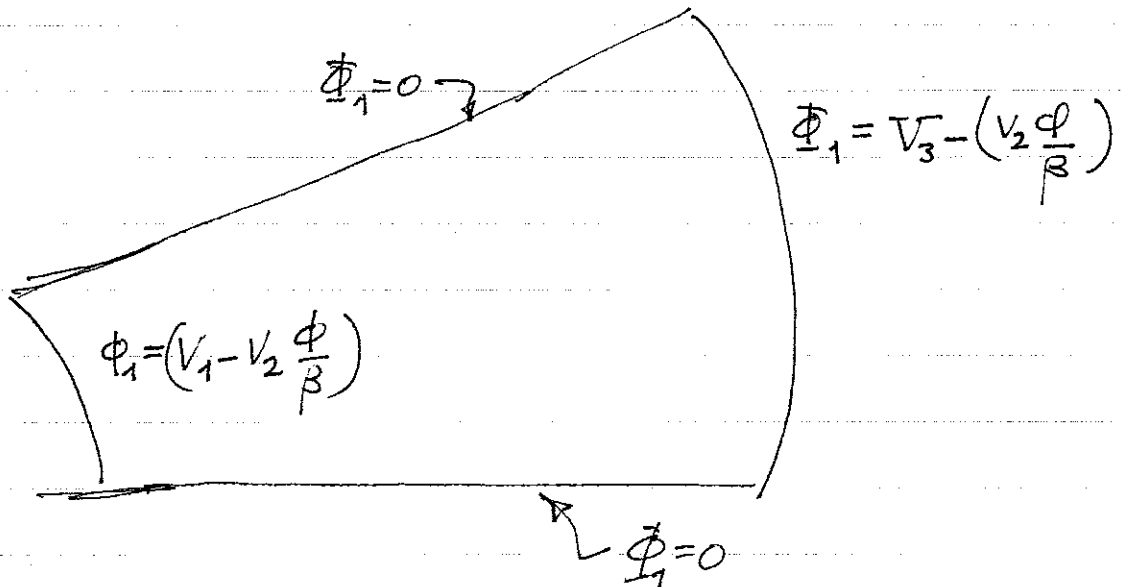
$$v_{13} \quad \psi = B_0 \phi \quad R = a_0$$

why? this gives $\bar{\Phi} = \text{constant}$ on upper boundary

$$\bar{\Phi} = V_2 \frac{\phi}{\beta} + \bar{\Phi}_1$$

~~where~~
$$\nabla^2 \bar{\Phi} = \nabla^2 V_2 \frac{\phi}{\beta} + \nabla^2 \bar{\Phi}_1 = 0$$

$$\nabla^2 \bar{\Phi}_1 = 0$$



$$\Phi_1 = \sum_{n=1}^{\infty} [a_n \rho^{\nu_n} + b_n \rho^{-\nu_n}] \sin \nu_n \phi$$

$$\nu_n = \frac{n\pi}{\beta} \quad n^{\text{th}} \text{ separation constant}$$

at $\rho = a$

$$\Phi_1 = \left(V_1 - V_2 \frac{\rho}{\beta} \right) = \sum_{n=1}^{\infty} [a_n a^{\nu_n} + b_n a^{-\nu_n}] \sin \nu_n \phi$$

multiply by $\sin \nu_m \phi$ and integrate 0 to β

$$[a_m a^{\nu_m} + b_m a^{-\nu_m}] \beta/2 = \int_0^{\beta} d\phi \sin \nu_m \phi \left[V_1 - V_2 \frac{\rho}{\beta} \right]$$

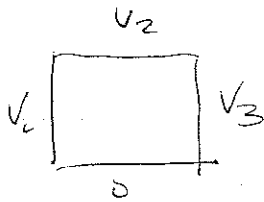
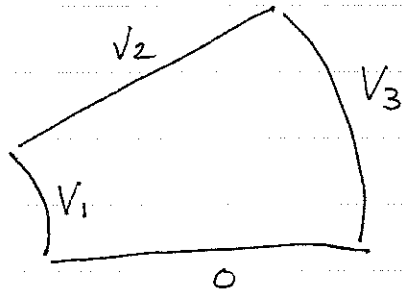
$$\equiv \beta/2 c_m$$

at $\rho = b$

$$[a_m b^{\nu_m} + b_m b^{-\nu_m}] \beta/2 = \int_0^{\beta} d\phi \sin \nu_m \phi \left[V_1 - V_2 \frac{\rho}{\beta} \right]$$

We are going to do the following problems:

1)



2)

$$\nabla^2 G = -4\pi \delta(\underline{x} - \underline{x}') \quad \text{in 2D}$$

WRITE

$$\phi_1 = \sum_{n=1}^{\infty} \left[c_n \left[\left(\frac{\rho}{a}\right)^{\nu_n} - \left(\frac{a}{\rho}\right)^{\nu_n} \right] + d_n \left[\left(\frac{b}{\rho}\right)^{\nu_n} - \left(\frac{\rho}{b}\right)^{\nu_n} \right] \right] \sin \nu_n \phi$$

$\stackrel{=0}{=} \text{on } a$
 $\stackrel{=0}{=} \text{on } b$

at $\rho = a$ $\phi_1 = V_1 - V_2 \frac{\rho}{\beta}$

$$= \sum_{n=1}^{\infty} d_n \left[\left(\frac{b}{a}\right)^{\nu_n} - \left(\frac{a}{b}\right)^{\nu_n} \right] \sin \nu_n \phi$$

multiply by $\sin \nu_m \phi$ and integrate 0 to β

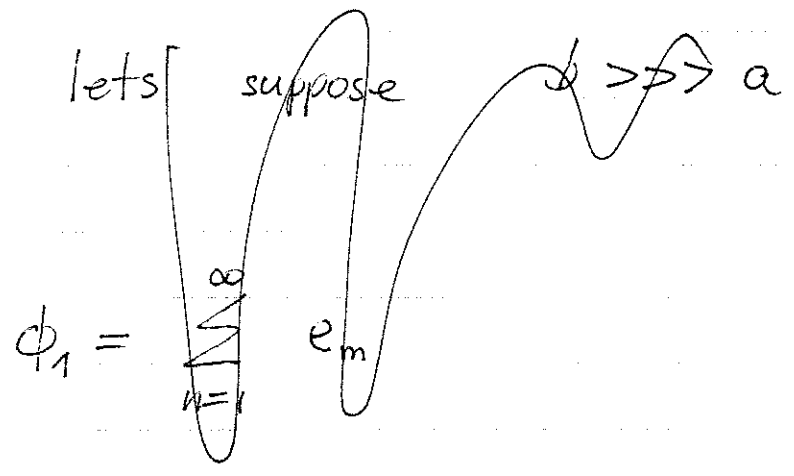
$$d_m \left[\left(\frac{b}{a}\right)^{\nu_m} - \left(\frac{a}{b}\right)^{\nu_m} \right] \frac{\beta}{2} = \int_0^{\beta} d\beta \sin \nu_m \phi \left[V_1 - V_2 \frac{\rho}{\beta} \right]$$

$$\equiv \frac{\beta}{2} d_m \quad \text{can be done}$$

Like wise

$$C_m \left[\left(\frac{b}{a}\right)^{\nu_m} - \left(\frac{a}{b}\right)^{\nu_m} \right] \frac{\beta}{2} = \int_0^\beta d\beta \sin \nu_m \varphi \left[\frac{V_3 - V_2 \frac{\varphi}{\beta}}{\beta} \right]$$

$$\equiv \frac{\beta}{2} e_m$$



(74)

Suppose $b \gg a, \rho$

and ~~$b \gg a$~~

$C_m \rightarrow 0$ as $\left(\frac{\rho}{b}\right)^{\nu_m} \rightarrow 0$

~~amplitude~~

$$\phi_1 \approx \sum_{n=1}^{\infty} \left(\frac{a}{\rho}\right)^{\nu_n} d_n \sin \nu_n \varphi$$

when $a = \rho$ gives $V_1 - V_2 \frac{\phi}{\beta}$ as it should

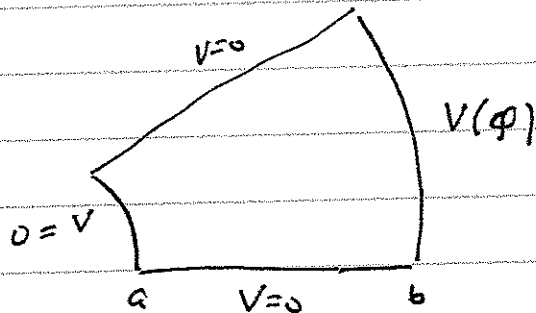
as $\rho/a \rightarrow \infty$ (but $\rho/b \ll 1$)

$$\phi_1 \rightarrow 0$$

as $b/a \rightarrow \infty$

$$d_m \sim \left(\frac{a}{b}\right)^{V_m}$$

$$c_m \sim \left(\frac{a}{b}\right)^{V_m}$$



skip

$$\phi_{\frac{1}{2}} = \sum_{n=1}^{\infty} \left[c_n \left(\frac{p}{a}\right)^{V_n} + d_n \left(\frac{a}{p}\right)^{V_n} \right] \sin V_n \phi$$

when $a = p = a$ $\phi_n = 0$ $d_n = -c_n$

as $p = b$

$$\phi_{\frac{1}{2}} = \sum_n c_n \left[\left(\frac{b}{a}\right)^{V_n} - \left(\frac{a}{b}\right)^{V_n} \right] \sin V_n \phi$$

$$= \sum_n V_n \sin V_n \phi$$

$$C_n = \frac{V_n}{\left[\left(\frac{b}{a}\right)^{\nu_n} - \left(\frac{a}{b}\right)^{\nu_n} \right]}$$

$$\phi = \sum_{n=1}^{\infty} \frac{\left[(\rho/a)^{\nu_n} - (a/\rho)^{\nu_n} \right] V_n \sin n\phi}{\left[\left(\frac{b}{a}\right)^{\nu_n} - \left(\frac{a}{b}\right)^{\nu_n} \right]}$$

$$\nu_n = \frac{n\pi}{\beta} \quad \text{as } (b/a) \rightarrow \infty \quad \frac{\left[(\rho/a)^{\nu_n} - (a/\rho)^{\nu_n} \right] V_n}{(b/a)^{\nu_n}}$$

leading term is
one with smallest ν_n ($n=1$)

$$\phi \approx \frac{(\rho/a) - (a/\rho)}{(b/a)} V_1 \sin \left(\frac{\pi \phi}{\beta} \right)$$

Green's Function in cylindrical coordinates (2D dimensional case)

$$\nabla^2 G(\mathbf{r}, \mathbf{r}') = -4\pi \delta(\mathbf{r} - \mathbf{r}')$$

In two dimensions ρ, ϕ

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} G + \frac{1}{\rho^2} \frac{\partial^2}{\partial \phi^2} G = -4\pi \frac{\delta(\phi - \phi') \delta(\rho - \rho')}{\rho}$$

two dimensional delta fun

$$\int d\phi' d\rho' \frac{\delta(\phi - \phi') \delta(\rho - \rho')}{\rho'} = 1$$

assume periodic in ϕ

write $G = \sum_{\nu=-\infty}^{\infty} \hat{G}_{\nu}(\rho, \rho', \phi') e^{i\nu\phi}$

$e^{-i\nu\phi} \int \frac{d\phi}{2\pi}$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} \hat{G}_{\nu} - \frac{\nu^2}{\rho^2} \hat{G}_{\nu} = -4\pi \int_0^{2\pi} \frac{d\phi}{2\pi} \frac{\delta(\phi - \phi') \delta(\rho - \rho')}{\rho} e^{-i\nu\phi}$$

multiply by $e^{i\nu\phi}$

can do ϕ integral

$$= -4\pi \frac{1}{2\pi} e^{-i\nu\phi'} \frac{\delta(\rho-\rho')}{\rho}$$

let $\hat{G}_\nu = e^{-i\nu\phi'} G_\nu(\rho, \rho')$

$$G = \sum_{\bar{\nu}=-\infty}^{\infty} G_{\bar{\nu}} e^{i\bar{\nu}(\phi-\phi')}$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} G_\nu - \frac{\nu^2}{\rho^2} G_\nu = -2 \frac{\delta(\rho-\rho')}{\rho}$$

How to solve such an equation

Option #1 Represent $G_\nu(\rho)$ in
a complete set of functions
on the interval $\rho=0$ $\rho=\infty$

Range is infinite separation
constant will be
continuous

$$G_\nu(\rho) = \int dk A_\nu(k) G_\nu(\rho, k)$$

where

Bessel transform
(Like Fourier transform)

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} G_\nu(\rho, k) - \frac{\nu^2}{\rho^2} G_\nu(\rho, k) + k^2 G_\nu(\rho, k) = 0$$

Bessel's equation

$$G_\nu(\rho, k) = J_\nu(k\rho)$$

Option #2 since we have an
ordinary equation solve it on

either side of δ function &
match

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for $\rho < \rho'$

$$G_V = a_V \rho^{|\nu|} + b_V \rho^{-|\nu|}$$

as $\rho \rightarrow 0$ G_V finite requires

we set $b_V = 0$

$$G_V = a_V \rho^{|\nu|} \quad \rho < \rho'$$

for $\rho > \rho'$

$$G_V = c_V \rho^{|\nu|} + d_V \rho^{-|\nu|}$$

as $\rho \rightarrow \infty$ G_V finite gives $c_V = 0$

$$G_V = d_V \rho^{-|\nu|}$$

conditions at $\rho = \rho'$

G is continuous

but $\frac{\partial G}{\partial \rho}$ is discontinuous

why

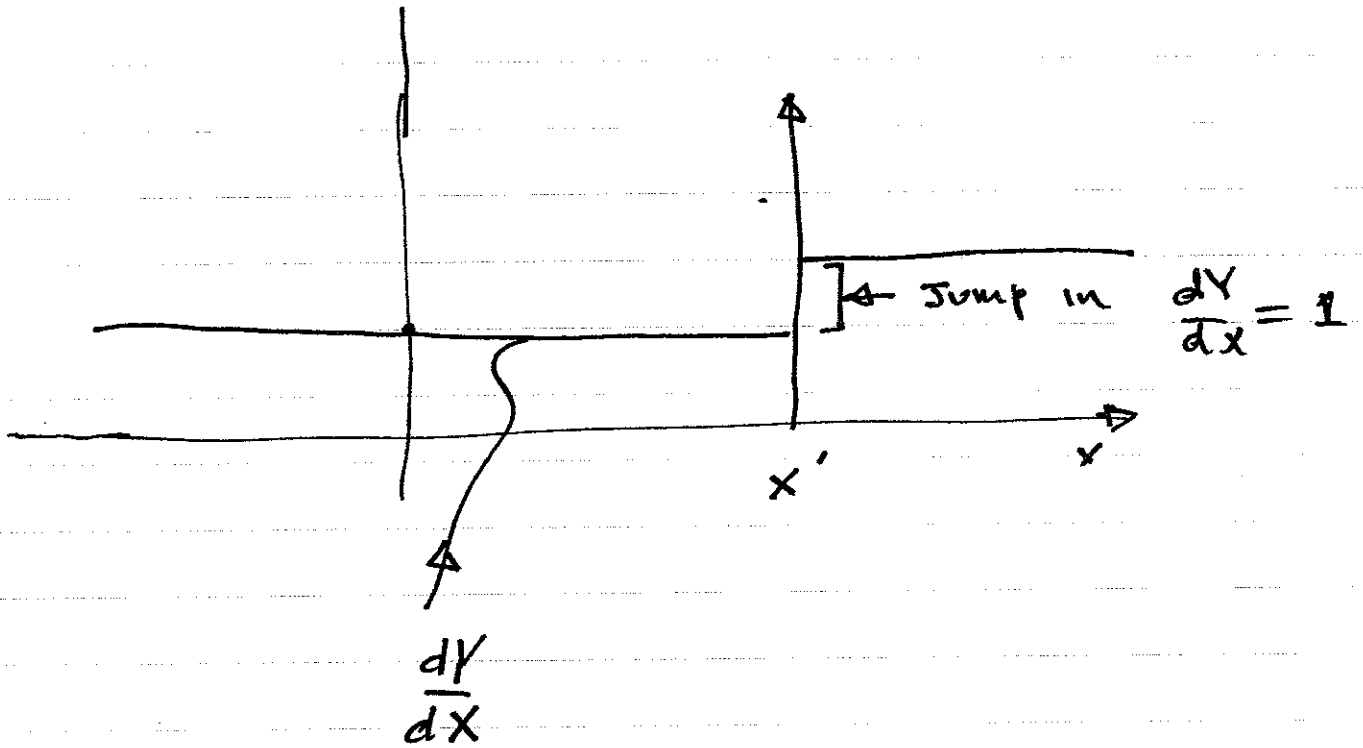
consider the equation

$$\frac{d^2}{dx^2} Y = \delta(x-x')$$

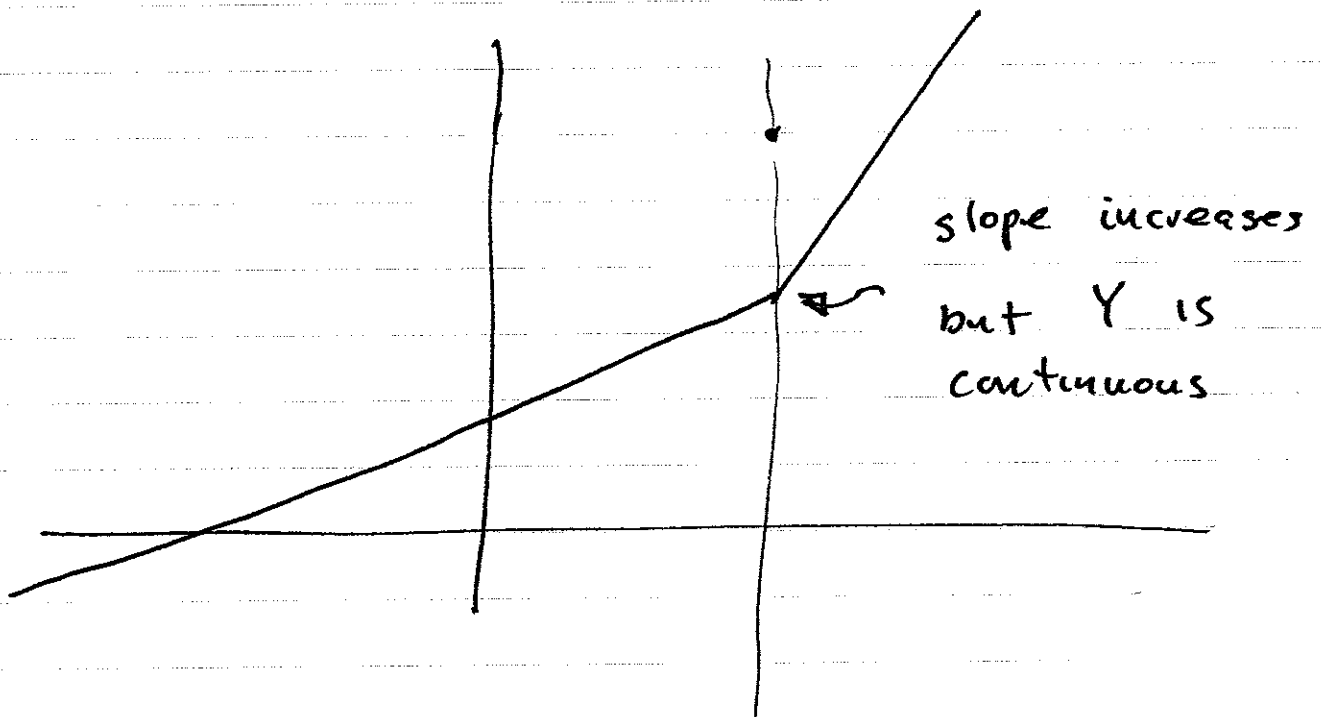
$$\frac{dY}{dx} = \int_{-\infty}^x dx' \delta(x-x')$$

$$= \frac{dY}{dx} \Big|_{-\infty} \quad x < x'$$

$$= \frac{dY}{dx} \Big|_{-\infty} + 1 \quad x > x'$$



integrate once more to get Y



continuity of G_{ν}

$$a_{\nu} \rho^{|\nu|} = d_{\nu} \rho^{-|\nu|}$$

discontinuity of $\partial G / \partial \rho$

for ρ near ρ'

$$\int_{\rho'-\epsilon}^{\rho'+\epsilon} d\rho \left[\frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} G_{\nu} - \frac{\nu^2}{\rho} G_{\nu} \right] = -2 \int_{\rho'-\epsilon}^{\rho'+\epsilon} d\rho \delta(\rho - \rho')$$

$\epsilon \rightarrow 0$

$$\rho \frac{\partial}{\partial \rho} G_{\nu} \Big|_{\rho'-\epsilon}^{\rho'+\epsilon} \sim \epsilon = -2$$

↑
Jump

$$-|\nu| d_{\nu} \rho^{-|\nu|} - |\nu| a_{\nu} \rho^{|\nu|} = -2$$

\uparrow
 $= d_{\nu} \rho^{-|\nu|}$

from continuity of G

$$d_\nu = \frac{1}{|\nu|} \rho'^{|\nu|}$$

$$a_\nu = \frac{1}{|\nu|} \rho'^{-|\nu|}$$

$$G_\nu(\rho, \rho') = \frac{1}{|\nu|} \begin{cases} \left(\frac{\rho}{\rho'}\right)^{|\nu|} & \rho < \rho' \\ \left(\frac{\rho'}{\rho}\right)^{|\nu|} & \rho > \rho' \end{cases}$$

$$G(\rho, \rho', \phi, \phi') = \sum_{\nu=-\infty}^{\infty} \frac{1}{|\nu|} e^{i\nu(\phi-\phi')} \begin{cases} \left(\frac{\rho}{\rho'}\right)^{|\nu|} & \rho < \rho' \\ \left(\frac{\rho'}{\rho}\right)^{|\nu|} & \rho > \rho' \end{cases}$$

~~Symmetric~~

~~what if $\nu=0$ take limit~~

must do $v < 0$ dem

$$G = a \quad p < p'$$

$$G = b + c \ln p \quad p > p' \quad [\text{diverges as } p \rightarrow \infty]$$

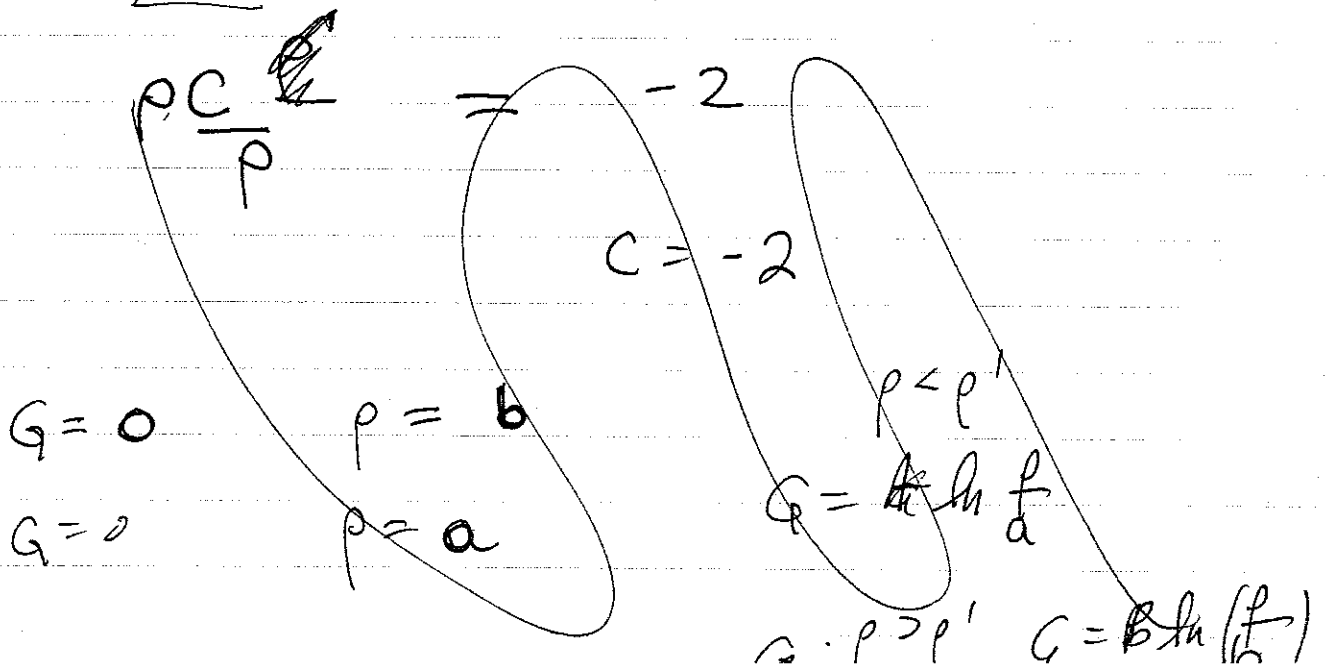
Match G jump $\frac{\partial G}{\partial p}$

potential due to line charge diverges

$$G = a \quad p < p'$$

$$G = a + c \ln(p/p') \quad p > p'$$

\angle constant undetermined



$\nu=0$ case

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial G}{\partial \rho} = -2 \frac{\delta(\rho - \rho')}{\rho}$$

integrate $\int_0^{\rho} \rho' d\rho' \left\{ \frac{1}{\rho'} \frac{\partial}{\partial \rho'} \rho' \frac{\partial G}{\partial \rho'} \right\} = \begin{cases} -2 & \rho > \rho' \\ 0 & \rho < \rho' \end{cases}$

$$\rho \frac{\partial G}{\partial \rho} = \cancel{0} \frac{\partial G}{\partial \rho} \Big|_0 = \begin{cases} -2 & \rho > \rho' \\ 0 & \rho < \rho' \end{cases}$$

$$\frac{\partial G}{\partial \rho} = \begin{cases} -2/\rho & \rho > \rho' \\ 0 & \rho < \rho' \end{cases}$$

$$G(\rho) = G(0) \quad \rho < \rho'$$

$$G(\rho) = G(0) - 2 \ln \frac{\rho}{\rho'} \quad \rho > \rho'$$

note $G \neq 0$ as $\rho \rightarrow \infty$