

Electrostatic potential Energy

consider $i=1, \dots, j-1$ charges q_i at points \vec{x}_i

now bring in the j^{th} charge to point \vec{x}_j this takes

an amount of work $W_j = q_j \phi(\vec{x}_j)$

$$W_j = q_j \sum_{i=1}^{j-1} \frac{q_i q_j}{|\vec{x}_j - \vec{x}_i|} \frac{1}{4\pi\epsilon_0}$$

THIS was the work required to bring the j^{th} charge

The work required to assemble all charges is n is

$$W = \sum_{j=1}^n W_j = \sum_{j=1}^n \sum_{i=1}^{j-1} \frac{q_i q_j}{|\vec{x}_j - \vec{x}_i|} \frac{1}{4\pi\epsilon_0}$$

$$= \frac{1}{2} \sum_{j=1}^n \sum_{\substack{i=1 \\ i \neq j}}^n \frac{q_i q_j}{|\vec{x}_j - \vec{x}_i|} \frac{1}{4\pi\epsilon_0}$$

double counting

for continuous charge density

$$W = \frac{1}{2} \int d^3x \int d^3x' \frac{\rho(x)\rho(x')}{4\pi\epsilon_0 |\underline{x} - \underline{x}'|}$$

(note no longer
worry about
 $\underline{x} = \underline{x}'$)

also can be written

$$W = \frac{1}{2} \int d^3x \frac{\rho(x)\phi(x)}{2}$$

BUT $\rho(x) = -\epsilon_0 \nabla^2 \phi(x)$

so

$$W = -\frac{1}{8\pi} \int d^3x \phi(x) \nabla^2 \phi(x) \frac{\epsilon_0}{2}$$

$$= \int d^3x \frac{|\nabla\phi|^2}{8\pi} \frac{\epsilon_0}{2}$$

assuming

$\phi \rightarrow 0$ as $x \rightarrow \infty$

$$W = \int d^3x \frac{\epsilon_0 |\underline{E}|^2}{8\pi}$$

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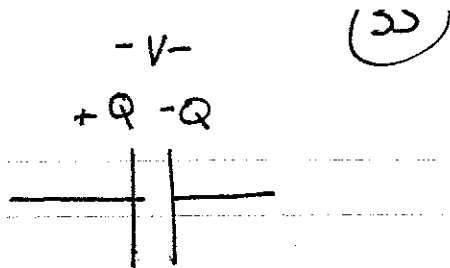
we can identify this as ^{potential} energy stored in the form of ~~electric~~ the electric field. One immediately runs into trouble here if one reverts to point charges. Why?

~~$$\int d^3x \frac{\epsilon_0 |\underline{E}|^2}{8\pi} \rightarrow \infty$$~~

The electric field energy associated with a point charge is ∞

self energy is included

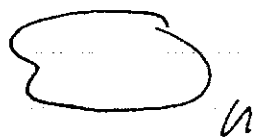
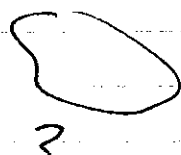
can compare difference of two states



$$Q = VC$$

~~Force on a~~

Capacitance: Consider a
System of n conductors



on each conductor there is a
surface charge density $\sigma_i(x)$ $i=1, n$
and a total charge $Q_i = \int da \sigma_i(x)$

Assuming the conductors are
 $\phi(\infty) = 0$

perfect conductors we have

$$\phi(\underline{x}) = V_i \text{ a constant on each conductor}$$

THUS ENERGY STORED

$$W = \frac{1}{2} \int d^3x \rho(x) \phi(x) = \frac{1}{2} \sum_{i=1}^n \int da \delta_i(x) \phi(x)$$

$$= \frac{1}{2} \sum_{i=1}^n V_i Q_i$$

~~$$V_i = \phi(\underline{x}) = \sum_{j=1}^n \int da \delta_j(x)$$

\underline{x} on i th conductor~~

in general inverse is C_{ij}

V_i is proportional to each of the Q_j { vice versa

capacitance matrix

$V_i = \sum_j M_{ij} Q_j$ Reciprocals $M_{ij} = M_{ji}$

This implies

37a

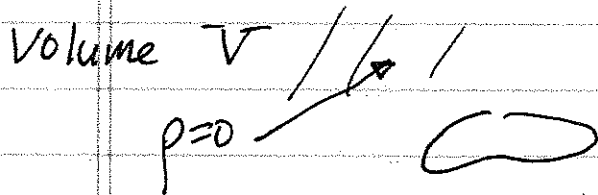
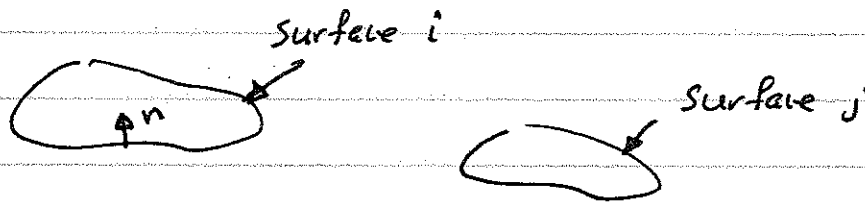
$$Q_i = \sum_j C_{ij} V_j$$

C_{ij} capacitance matrix

Follows from Green's theorem

$$\phi(x') = \int \frac{d^3x}{4\pi\epsilon_0} \rho(x) G_0(x, x') - \frac{1}{4\pi} \int_S dA \phi(x) \underline{n} \cdot \nabla G_0(x, x')$$

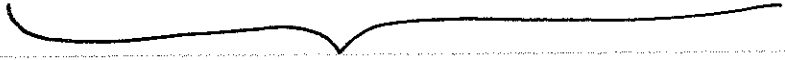
↙ outward normal



$$Q_i = \int_{S_i} \epsilon_0 \underline{E}(x') \cdot (-\underline{n}) dA' = \int_{S_i} \epsilon_0 \underline{n} \cdot \nabla' \phi(x') dA'$$

$$\underline{n} \cdot \nabla' \phi(x') = -\frac{1}{4\pi} \sum_j V_j \int_{S_j} dA \underline{n} \cdot \nabla' (\underline{n} \cdot \nabla' G_0(x', x'))$$

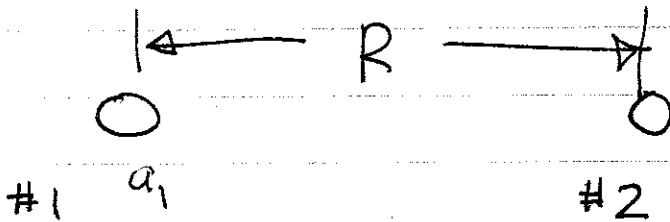
$$Q_i = \sum_j V_j \left(- \frac{1}{4\pi} \int_{S_i} dA' \int_{S_j} dA \underline{n} \cdot \nabla (\underline{n} \cdot \nabla G_0(\underline{x}, \underline{x}')) \right)$$



 $C_{ij} =$

$$C_{ij} = C_{ji} \quad \text{because} \quad G_0(\underline{x}, \underline{x}') = G_0(\underline{x}', \underline{x})$$

Example, calculate the capacitance matrix
~~due to~~ for two small spheres
 of radii^u a_1, a_2 ~~and~~ separated
 by a distance R with $R \gg a_1, a_2$



$R \gg a$ allows us to neglect the rearrangement of charge on spheres due to presence of other spheres.

Calculate potential due to point charge

$$V_{\phi_1} = \frac{Q_1}{4\pi\epsilon_0 a_1} + \frac{Q_2}{4\pi\epsilon_0 R} \quad \left(V_{\phi_2} = \frac{Q_1}{4\pi\epsilon_0 R} + \frac{Q_2}{a_2} \right)$$

$$\begin{pmatrix} V_1 \\ V_2 \end{pmatrix} = \frac{1}{4\pi\epsilon_0} \begin{bmatrix} 1/a_1 & 1/R \\ 1/R & 1/a_2 \end{bmatrix} \begin{pmatrix} Q_1 \\ Q_2 \end{pmatrix}$$

Invert

$$\begin{pmatrix} Q_1 \\ Q_2 \end{pmatrix} = \frac{4\pi\epsilon_0}{\frac{1}{a_1 a_2} - \frac{1}{R^2}} \begin{bmatrix} \frac{1}{a_2} & -\frac{1}{R} \\ -\frac{1}{R} & \frac{1}{a_1} \end{bmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix}$$

neglect

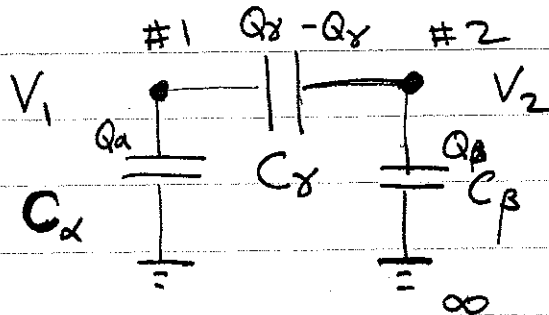
$$\begin{pmatrix} Q_1 \\ Q_2 \end{pmatrix} = 4\pi\epsilon_0 \begin{bmatrix} C_{11} & -\frac{C_{12}}{R} \\ -\frac{C_{12}}{R} & C_{22} \end{bmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix}$$

Sym

Symmetric

SKIP

~~What is the capacitance~~



Equivalent circuit

$$Q_\gamma = C_\gamma (V_1 - V_2)$$

$$Q_\alpha = C_\alpha V_1$$

$$Q_\beta = C_\beta V_2$$

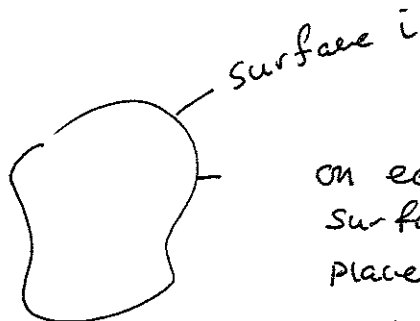
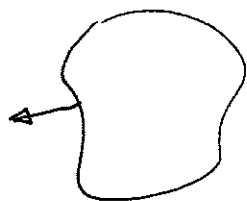
$$Q_1 = (Q_\alpha + Q_\gamma) = \overset{C_{11}}{C_\alpha V_1} + \overset{C_{12}}{C_\gamma (V_1 - V_2)}$$
$$= (C_\alpha + C_\gamma) V_1 - C_\gamma V_2$$

$$Q_2 = (Q_\beta - Q_\gamma) = \overset{C_{21}}{C_\beta V_2} - \overset{C_{21}}{C_\gamma V_1} + \overset{C_{22}}{(C_\beta + C_\gamma) V_2}$$

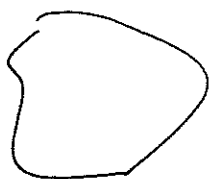
①

1.13

Thomson's



on each
surface
place Q_i
Coulombs
of charge
distributed
according to
surface charge
density σ_i



Form of σ_i is unknown

$$\int_{S_i} dS \sigma_i = Q_i$$

Minimize

$$W = \int d^3x \frac{\epsilon_0 |\nabla\phi|^2}{2}$$

where $\nabla^2\phi = 0$
except at surfaces
where

$$\sigma_i = \epsilon_0 [n \cdot \nabla(\phi_{out}) - n \cdot \nabla\phi_{in}]$$

let $\sigma_i = \sigma_{i0} + \epsilon \sigma_{it}$ ← test
 ↙ actual minimizing charge

②

$$W = \int_{\substack{V \\ \text{outside} \\ \text{all}}} d^3x \epsilon_0 \frac{|\nabla\phi|^2}{2} + \sum_i \int_{V_i} d^3x \epsilon_0 \frac{|\nabla\phi|^2}{2}$$

$$\cancel{W} = \int_V d^3x \frac{\epsilon_0}{2} \nabla\phi \cdot \nabla\phi$$

$$= \int_S ds \frac{\epsilon_0}{2} \phi \overset{\text{outward normal}}{n} \cdot \nabla\phi - \int_V d^3x \frac{\epsilon_0}{2} \phi \nabla^2\phi$$

For our problem $\nabla^2\phi = 0$

$$W = \sum_i \int_{S_i} ds \frac{\epsilon_0}{2} \phi (n \cdot \nabla\phi_{\text{out}} - n \cdot \nabla\phi_{\text{in}})$$

$$\phi = \phi_0 + \epsilon\phi_\epsilon$$

$$W(\epsilon) = \int_{\substack{V \\ \text{ALL-V}}} d^3x \frac{\epsilon_0}{2} \left[|\nabla\phi_0|^2 + 2\epsilon \nabla\phi_0 \cdot \nabla\phi_\epsilon + \epsilon^2 |\nabla\phi_\epsilon|^2 \right]$$

$$\left. \frac{\partial W}{\partial \epsilon} \right|_{\epsilon=0} = \int_{\text{ALL-V}} d^3x \epsilon_0 \nabla\phi_0 \cdot \nabla\phi_\epsilon$$

3

$$= \sum_i \int_{S_i} ds \frac{\epsilon_0}{4\pi} \phi_{0i} \underbrace{(n \cdot \nabla \phi_{t, \text{out}} - n \cdot \nabla \phi_{t, \text{in}})}$$

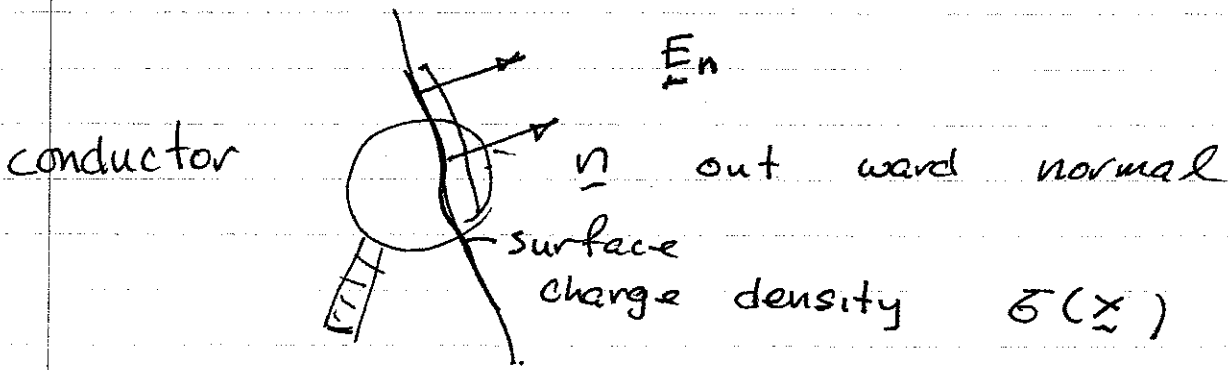
$$= \sum_i \int_{S_i} ds \phi_{0i} \bar{\epsilon}_{it} = 0$$

$$\int \bar{\epsilon}_{i0} ds = Q_i \quad \int \bar{\epsilon}_{it} ds = 0$$

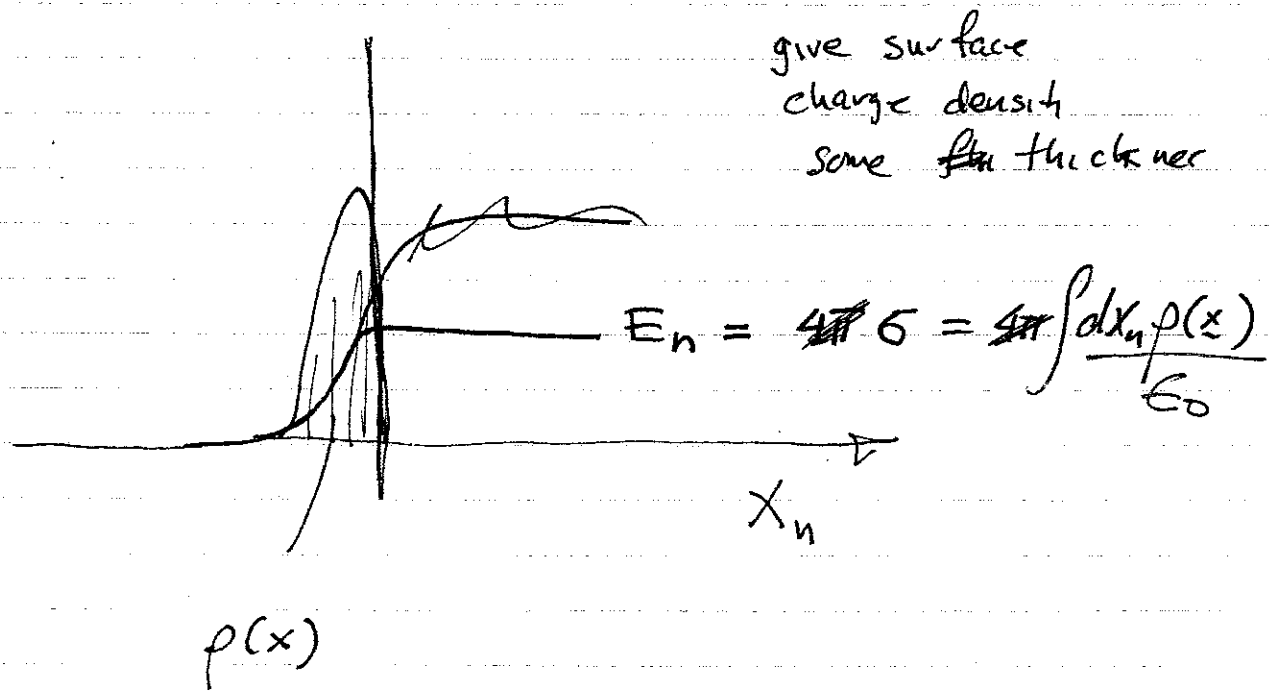
only way integral can be zero is if

$$\underline{\underline{\phi_{0i} = \text{const}}}$$

Force on the surface of a conductor



Blow up of surface



$$\frac{\partial E_n}{\partial x_n} = \frac{4\pi \rho(x_n)}{\epsilon_0} \quad \sigma = \int dx_n \rho(x')$$

small force act

$$\vec{dF}_x =$$

$$\vec{F} = \int d^3x \rho(x) \vec{E}$$

$$= \int da^n \int dx_n \rho E_n$$

normal direction

$$\frac{\partial E_n}{\partial x_n} = 4\pi\rho/\epsilon_0$$

$$= \int da^n \int \frac{\epsilon_0 E_n}{4\pi} \frac{\partial E_n(x_n)}{\partial x_n}$$

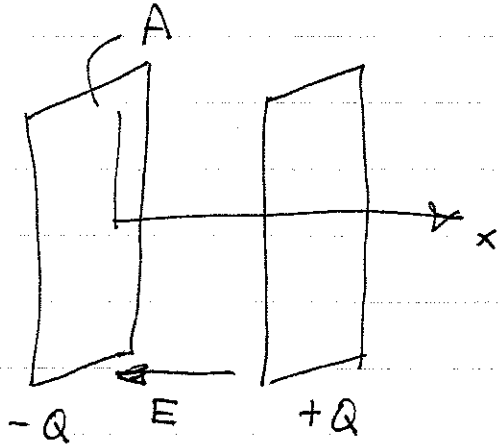
$$= \int da^n \frac{\epsilon_0}{4\pi} \frac{|E_n|_v^2}{2}$$

out.

electric pressure = $-\frac{\epsilon_0}{4\pi} \frac{|E_n|_v^2}{2}$

Skop

FORCE on a parallel plates

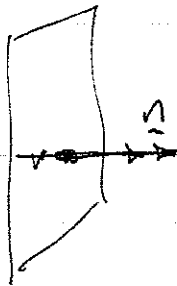


surface charge density

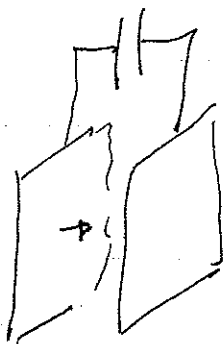
$$\vec{E}_a = -4\pi \left(\frac{Q}{A}\right) \hat{a}_x$$

Force on plate

$$\vec{F} = \int da \hat{n} \frac{|\vec{E}_n|^2}{8\pi}$$



attraction to other plate



Check out Problem 1.9 use virtual works

$$\delta W = \vec{F} \cdot \delta \vec{x} = \delta \left[\frac{1}{2} C V^2 \text{ or } \frac{Q^2}{2C} \right]$$

must include $V \delta Q$ / constant \checkmark