

Cavities

ENEE 381

Role of Cavities

Cavities are resonant structures: Support EM modes at specific frequencies.

Used in:

Filters

Oscillators

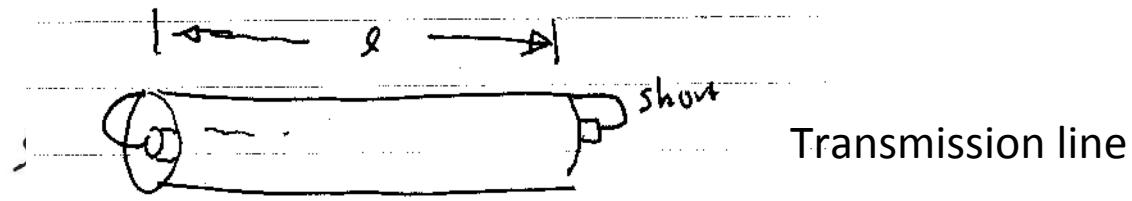
Amplifiers

Measurement of material properties

Resonance

natural frequency of oscillation of fields

Example:



Terminated in short circuits

$$k l = n\pi$$

$$\cancel{\omega = k \cdot} \quad k = \frac{\omega}{v}$$

$$\boxed{\omega_n = \frac{n\pi}{l} v}$$

$$n = 1, 2, \dots$$

$$\frac{d^2}{dx^2} V(x) + \frac{\omega^2}{v^2} V(x) = 0$$

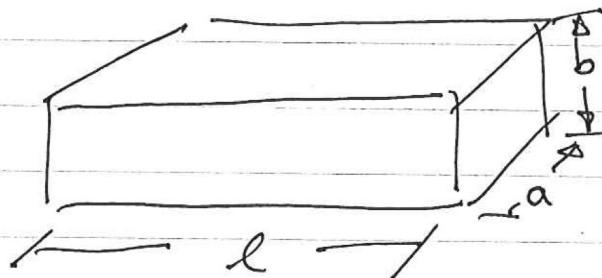
$$V(x) = V_+ e^{i\omega x/v} + V_- e^{-i\omega x/v}$$

$$V(x=0) = 0$$

$$V(x=l) = 0$$

Enclosed Rectangular Prism

~~enclosed~~ enclosed box

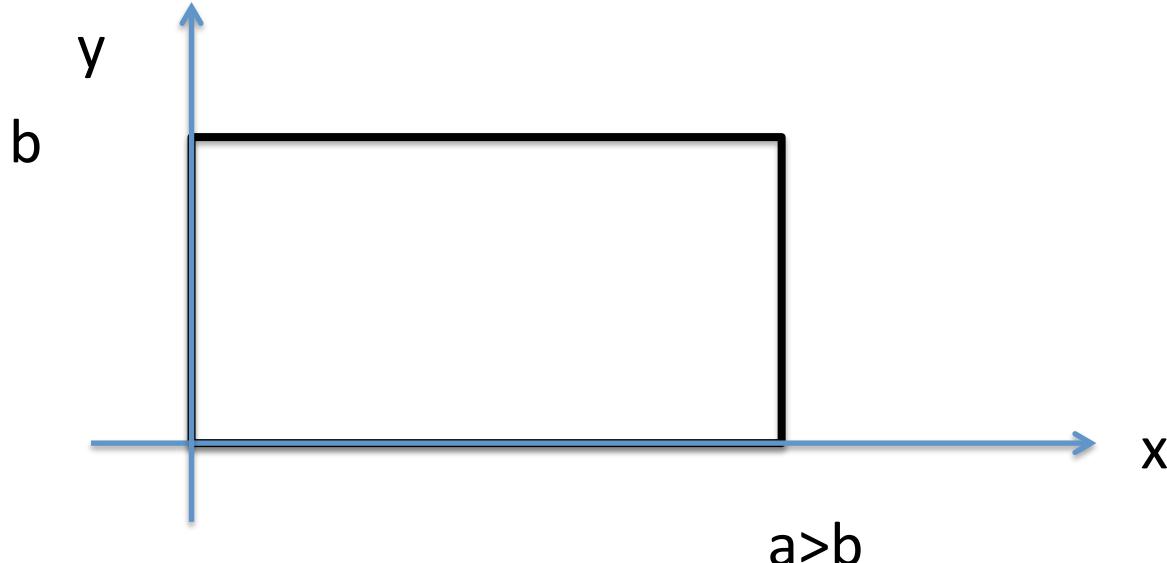


TE_{nm} or TM_{nm}

$$\omega_{nmp} = V \sqrt{\left(\frac{n\pi}{a}\right)^2 + \left(\frac{m\pi}{b}\right)^2 + \left(\frac{p\pi}{c}\right)^2}$$

n, m, p
are
integers

Modes of a Rectangular WG



$$\text{TM}_{nm}: \hat{E}_z = E_0 \sin(k_x x) \sin(k_y y) \quad k_x = \frac{n\pi}{a}, \quad k_y = \frac{m\pi}{b}; \quad n,m=1,2,3,\dots$$

$$\text{TE}_{nm}: \hat{H}_z = H_0 \cos(k_x x) \cos(k_y y) \quad k_x = \frac{n\pi}{a}, \quad k_y = \frac{m\pi}{b}; \quad n,m=0^*,1,2,3,\dots$$

* one or the other, but not both

Cut-Off frequencies

$$\omega_{c,n,m} = v \sqrt{\left(\frac{n\pi}{a}\right)^2 + \left(\frac{m\pi}{b}\right)^2}$$

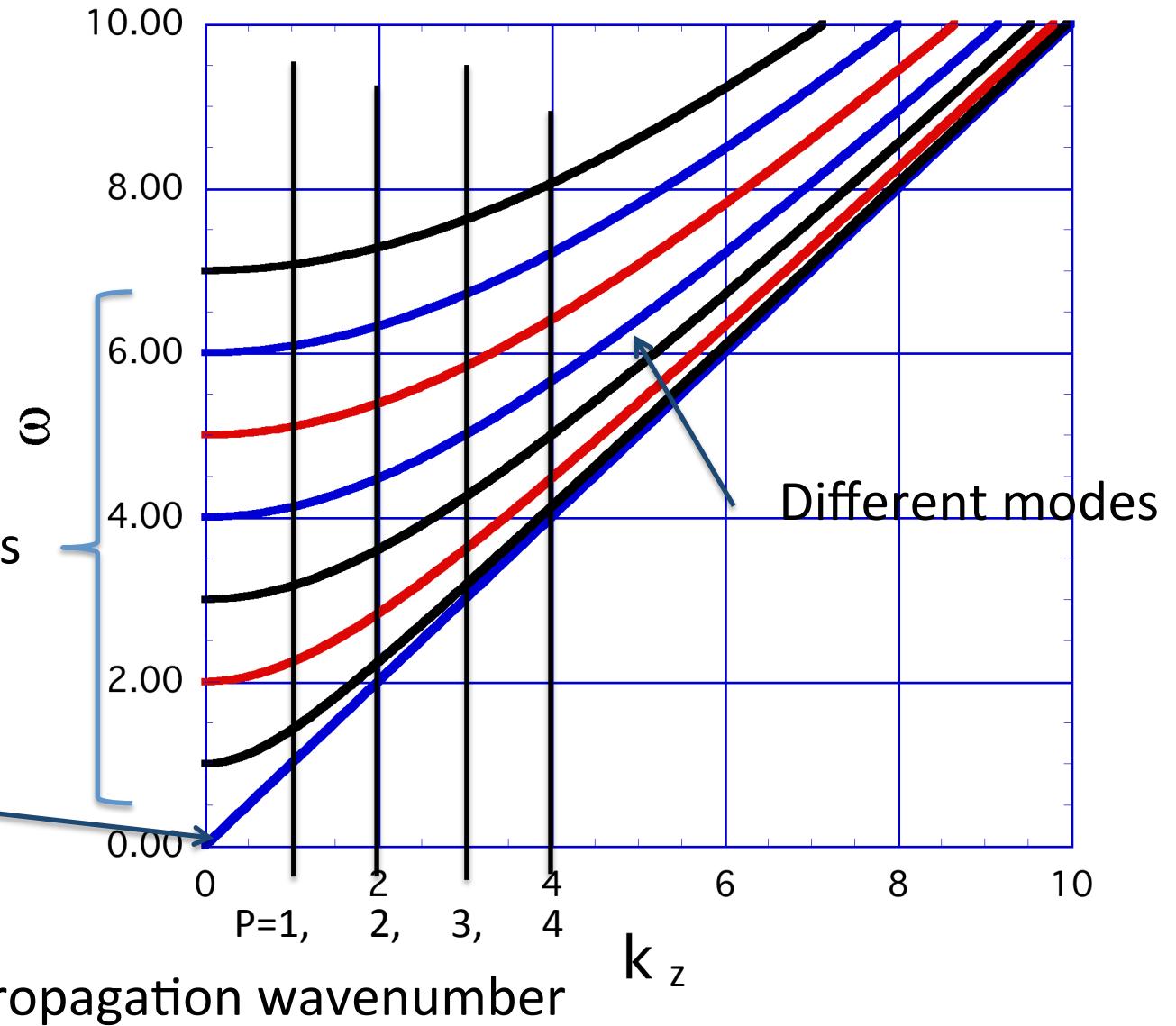
WG Dispersion Relations

$$\frac{\omega^2}{v^2} = k_{\perp}^2 + k_z^2$$

$$k_z = p \frac{\pi}{L}$$

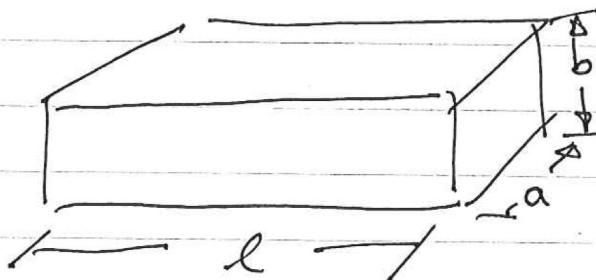
Cut off frequencies

Transmission
Lines Only



Enclosed Rectangular Prism

~~enclosed~~ enclosed box



TE_{nm} or TM_{nm}

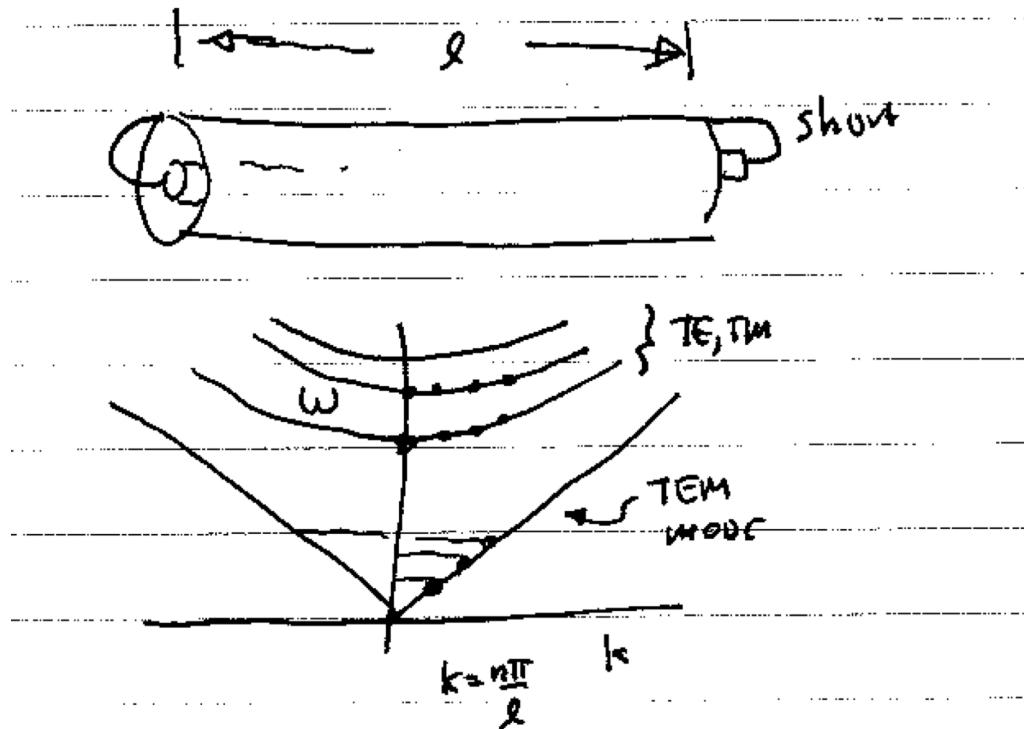
$$\omega_{nmp} = \sqrt{\left(\frac{n\pi}{a}\right)^2 + \left(\frac{m\pi}{b}\right)^2 + \left(\frac{p\pi}{l}\right)^2}$$
$$\omega^2 = k_z^2 v^2 + \omega_{c,nm}^2$$

n, m, p
are
integers

$$\omega_{c,n,m} = v \sqrt{\left(\frac{n\pi}{a}\right)^2 + \left(\frac{m\pi}{b}\right)^2}$$

$$k_z = p \frac{\pi}{L}$$

Transmission Line – TEM mode

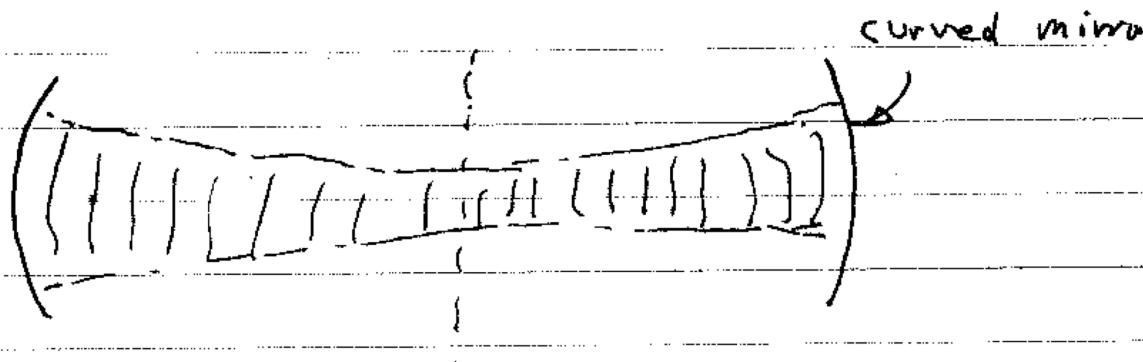


Operate at frequencies well below
TE and TM cut-off

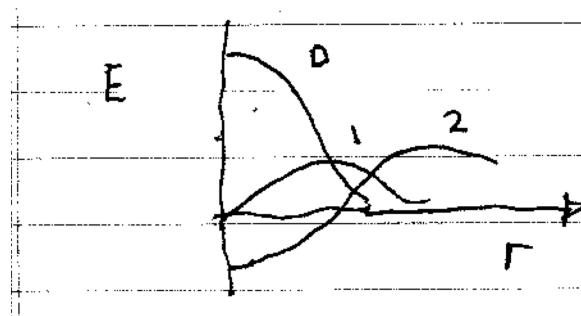
Fabry-Perot Cavity

Fabry-Perot Cavity

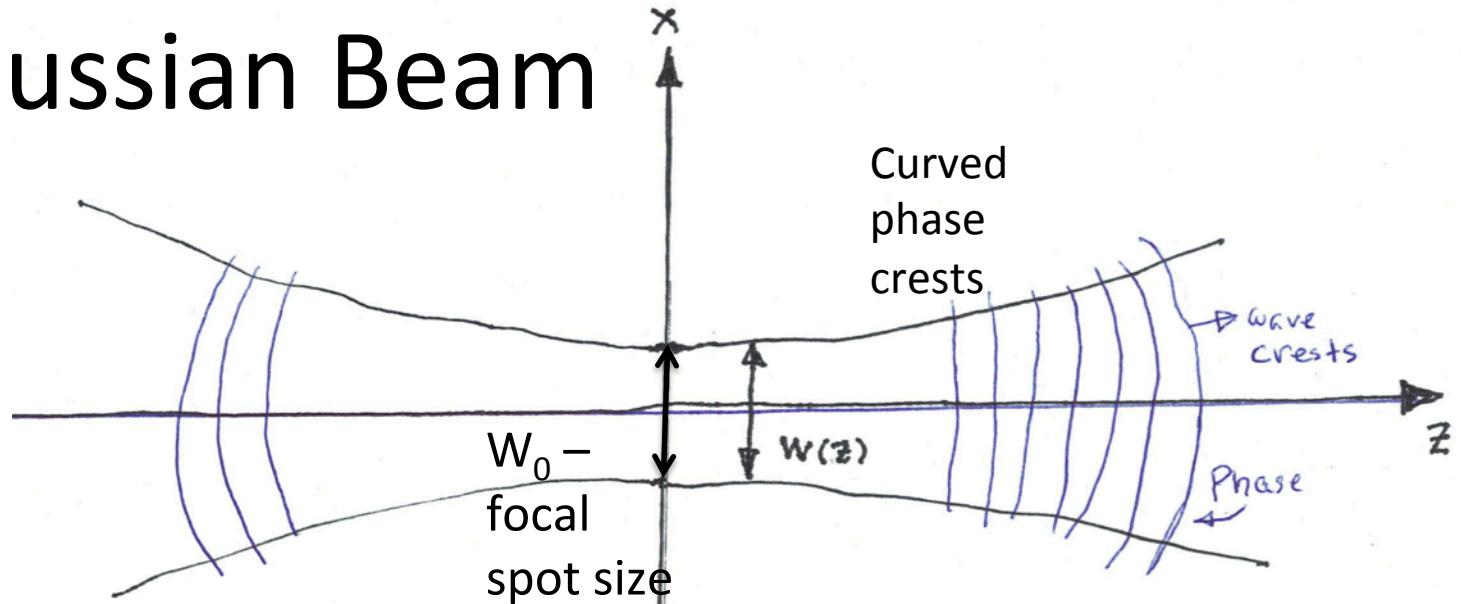
Curved
mirror



Gaussian - Laguerre
modes



Gaussian Beam



$$E_{x,y}(x,y,z) = \frac{E_0}{1 + iz/Z_R} \exp \left[-\frac{(x^2 + y^2)}{W_0^2(1 + iz/Z_R)} + ikz \right]$$

$$W(z) = W_0 \sqrt{1 + z^2 / Z_R^2} \quad Z_R = \frac{1}{2} k W_0^2 \quad \text{Rayleigh Length}$$

$$E_{x,y}(0,0,z) = \frac{E_0}{1 + iz/Z_R} \exp[+ikz] = \frac{E_0}{\sqrt{1 + (z/Z_R)^2}} \exp[+ikz + i\phi(z)]$$

$$\text{Guoy Phase} \quad \tan \phi = -z / Z_R$$

Gaussian Beam

$$E_{x,y}(x,y,z) = \frac{E_0}{1 + iz/Z_R} \exp\left[-\frac{(x^2 + y^2)}{W_0^2(1 + iz/Z_R)} + ikz\right]$$
$$= \frac{E_0}{\sqrt{1 + z^2/Z_R^2}} \exp\left[-\frac{(x^2 + y^2)}{W_0^2(1 + z^2/Z_R^2)}\right] \exp\left\{i\left[kz + \frac{z(x^2 + y^2)}{Z_R W_0^2(1 + z^2/Z_R^2)} + \phi_G\right]\right\}$$

The diagram illustrates the decomposition of the Gaussian beam amplitude and phase. It consists of two horizontal blue brackets. The left bracket groups the first two terms of the equation, labeled "Amplitude". The right bracket groups the remaining terms, labeled "Phase".

Pick parameters such that phase is constant on surface of mirror.

The pick k such that the phase changes by $p\pi$ in going from one mirror to the next

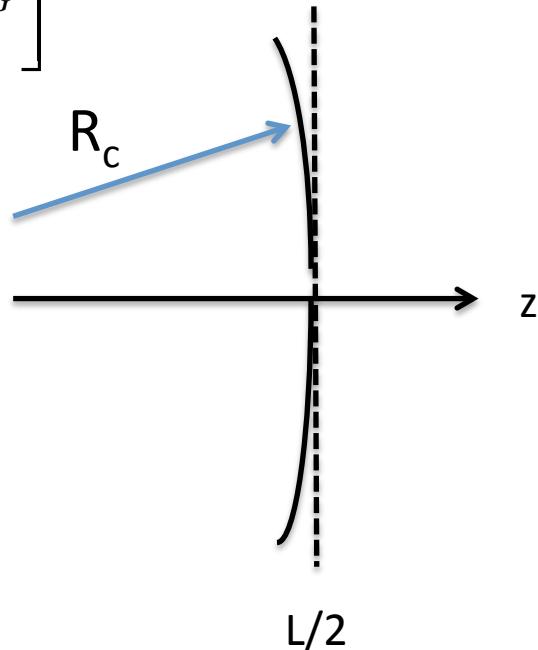
Wavebeam phase $\left[k_z + \frac{z(x^2 + y^2)}{Z_R W_0^2 (1 + z^2 / Z_R^2)} + \phi_G \right]$

Surface of mirror $z \approx \frac{L}{2} - \frac{(x^2 + y^2)}{2R_c}$

Will match if

$$\frac{L}{2R_c} = \frac{(L/2)^2}{Z_R^2 + (L/2)^2} < 1$$

Phase change mirror to mirror $2 \left(k_p \frac{L}{2} + \phi_G \right) = p\pi$



For a given L and R_c Z_R is determined above.

Hence W_0 focal spot determined.

$$Z_R = \frac{1}{2} k W_0^2$$

Design a Fabrey-Perot resonator

Requirements:

Wavelength 1 micron = 10^{-6} m.

Focal spot size = 100 microns = -10^{-4} m.

Spot size on mirrors = 300 microns = 3×10^{-4} m

Find L and R_c

Super bonus: How big must the mirrors be to keep “spill over” below 10%

Quality Factor

Quality factor \Rightarrow measures losses Q

big Q low loss low Q high loss

Time domain

$$\text{Defn} \quad \left(\frac{\omega U}{P_0} \right) = Q$$

$V(t)$

$E(t)$

$$e^{-\frac{\omega}{2Q}t}$$

decay of envelope

period of oscillation

$$\frac{1}{Q} \equiv \frac{\text{Power Dissipated}}{\omega \text{ Energy Stored}}$$

Field decay rate

$$E, H \sim \exp\left(-\frac{\omega}{2Q}t\right)$$

Frequency Domain

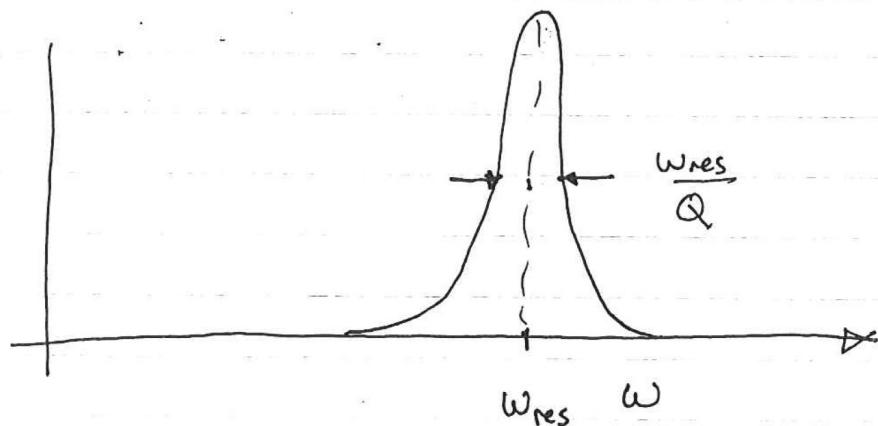
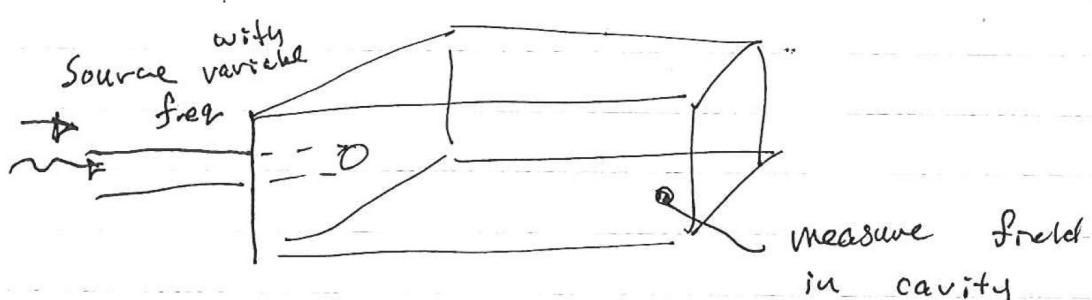
Steady state field response

$$E, H \sim \frac{\text{Source}}{\omega - \omega_{res} + i \frac{\omega_{res}}{2Q}}$$

$$|E|^2, |H|^2 \sim \frac{1}{\left(\omega - \omega_{res}\right)^2 + \left(\frac{\omega_{res}}{2Q}\right)^2}$$

$$\text{Half maximum } \omega - \omega_{res} = \pm \frac{\omega_{res}}{2Q}$$

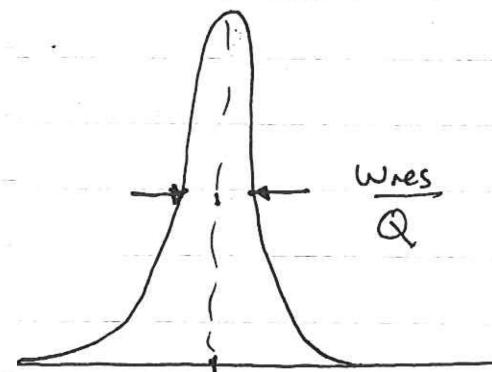
$$\text{Full Width at Half Maximum (FWHM)} = \frac{\omega_{res}}{Q}$$



Response Function

$$|E|^2 \propto$$

$$\frac{1}{(\omega - \omega_{\text{res}})^2 + \left(\frac{\omega_{\text{res}}}{2Q}\right)^2}$$



when $\omega = \omega_{\text{res}} \pm \frac{\omega_{\text{res}}}{2Q}$

$|E|^2$ is $\frac{1}{2}$ of peak value

Full width at half max of Power
(FWHM) $\rightarrow \frac{\omega_{\text{res}}}{Q}$

Multiple Contributions to Loss

if. losses are small ($\alpha \gg 1$)

losses are additive

 different loss mechanism

$$\frac{1}{Q} = \frac{P_d}{\omega T} = \frac{P_{d1}}{\omega T} + \frac{P_{d2}}{\omega T} + \dots$$

$$= \frac{1}{Q_1} + \frac{1}{Q_2} + \dots$$

Reciprocals of Q add.

Dielectric and Conductor Loss

α due to lossy dielectric

$$\epsilon = \epsilon' - j\epsilon''$$

$$Q = \frac{\epsilon'}{\epsilon''}$$

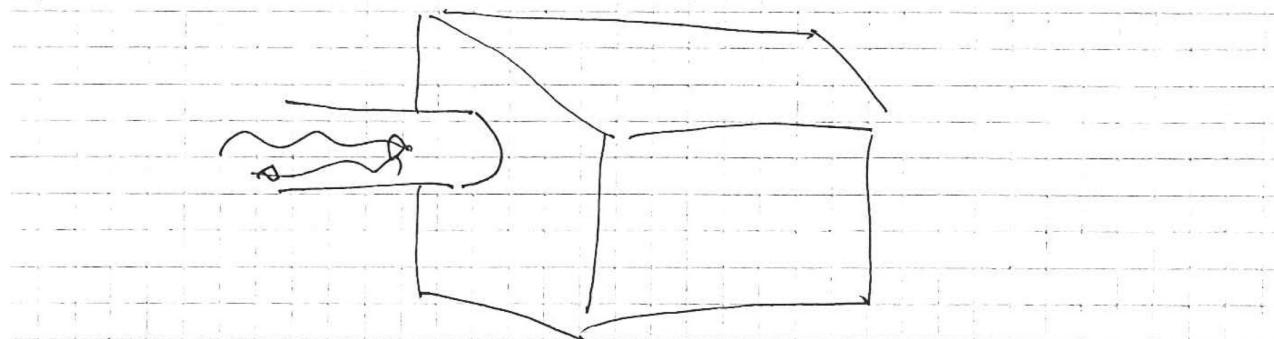
ϵ'
 ϵ must completely
fill cavity

Losses due to conductors

$$\alpha = \frac{\omega}{R_s} \left(\frac{\pi}{R_s} \right) \frac{\int d^3x |\hat{H}|^2}{\int da |\hat{H}_t|^2} \begin{matrix} \leftarrow \text{energy stored} \\ \leftarrow \text{losses} \end{matrix}$$

Coupling to Cavities

A closed box is useless

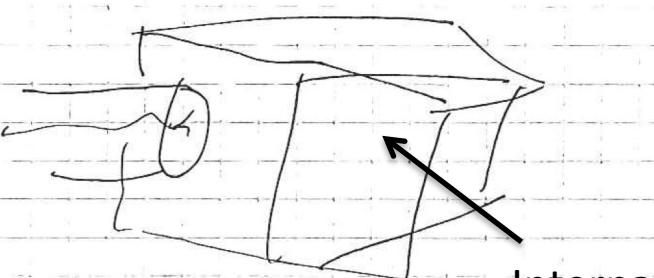


need to be able to get power in it

Adding a hole (or coupling port does two things) power can come in and power can go out.

Coupling Also Characterized by Q

Modifies Q


$$\frac{1}{Q_T} = \frac{1}{Q_{\text{internal}}} + \frac{1}{Q_{\text{coupling}}}$$

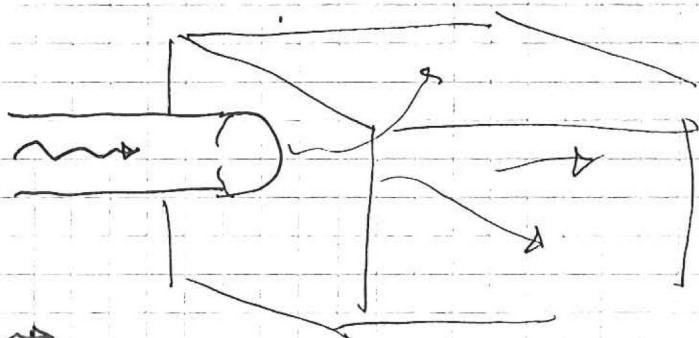
Q_{internal}

Critically Coupled

$$Q_{\text{coupling}} = Q_{\text{internal}}$$

V_{incident}

$V_{\text{reflected}}$



At resonance
all power
from source
is absorbed
in cavity!

Voltage reflection coefficient

$$\rho = - \frac{i(\omega/\omega_{\text{res}} - 1) - \left(\frac{1}{2Q_i} - \frac{1}{2Q_e} \right)}{i(\omega/\omega_{\text{res}} - 1) - \left(\frac{1}{2Q_i} + \frac{1}{2Q_e} \right)}$$

~~$\rho = \frac{V_{\text{reflected}}}{V_{\text{incident}}}$~~

Universal Response

$$\rho = - \frac{i(\omega_{res} - 1) - \left(\frac{1}{2Q_i} - \frac{1}{2Q_e} \right)}{i(\omega_{res} - 1) - \left(\frac{1}{2Q_i} + \frac{1}{2Q_e} \right)}$$

$$\rho = \frac{V_{reflected}}{V_{incident}}$$

Knowing Q_i , Q_e and ω_{res} determines

far from resonance $\rho = -1$

Reflectivity at resonance

0 - if $Q_i = Q_e$

$$|P|_{res}^2 = \left(\frac{Q_i^{-1} - Q_e^{-1}}{Q_i^{-1} + Q_e^{-1}} \right)^2$$

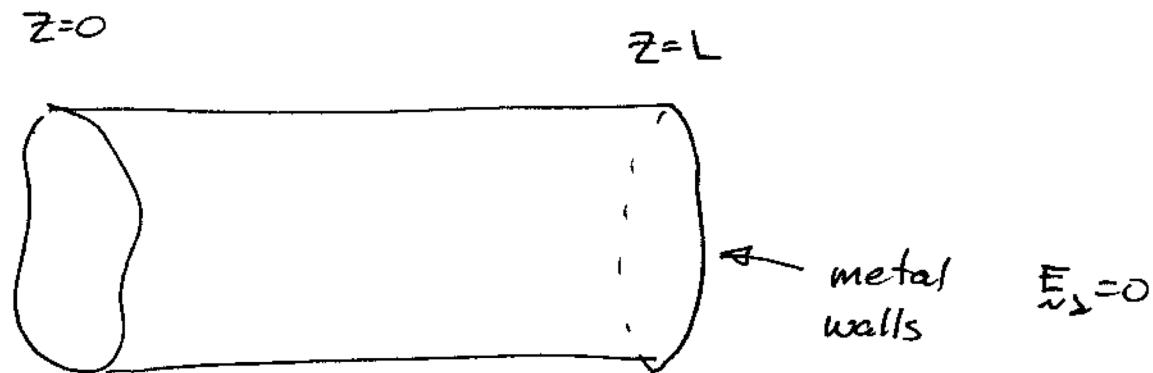
damping rate for fields $\rho \rightarrow \infty$

$$\omega_{res} = 1 - i \frac{1}{2\alpha_r}$$

$$\alpha_r = \frac{1}{Q_i} + \frac{1}{Q_e}$$

Waveguide Cavities

CAVITIES constructed from cylindrical
waveguides will have normal modes



Waveguide fields

$$\mathbf{E} = \operatorname{Re} \left\{ \hat{\mathbf{E}}(x, y) \exp \left[i(k_z z - \omega t) \right] \right\}$$

$$\mathbf{H} = \operatorname{Re} \left\{ \hat{\mathbf{H}}(x, y) \exp \left[i(k_z z - \omega t) \right] \right\} \quad \frac{\partial^2 \hat{H}_z}{\partial x^2} + \frac{\partial^2 \hat{H}_z}{\partial y^2} = - \left((\omega / v)^2 - k_z^2 \right) \hat{H}_z$$

$$\frac{\partial^2 \hat{E}_z}{\partial x^2} + \frac{\partial^2 \hat{E}_z}{\partial y^2} = - \left((\omega / v)^2 - k_z^2 \right) \hat{E}_z$$

$$\hat{\mathbf{E}}_{\perp} = \frac{i}{(\omega / v)^2 - k_z^2} \left[k_z \nabla_{\perp} \hat{E}_z - \omega \mu \hat{\mathbf{z}} \times \nabla_{\perp} \hat{H}_z \right] \quad \hat{E}_z \Big|_{wall} = 0, \quad \mathbf{n} \cdot \nabla_{\perp} \hat{H}_z \Big|_{wall} = 0$$

$$\hat{\mathbf{H}}_{\perp} = \frac{i}{(\omega / v)^2 - k_z^2} \left[k_z \nabla_{\perp} \hat{H}_z + \omega \epsilon \hat{\mathbf{z}} \times \nabla_{\perp} \hat{E}_z \right]$$

Forward and Backward Waves

TM modes

$$\hat{E}_z = \hat{E}_{z,nm}(x, y) \left(A_+ \exp(i k_z z) + A_- \exp(-i k_z z) \right)$$

$$\hat{\mathbf{E}}_\perp = \frac{i k_z \nabla_\perp \hat{E}_{z,nm}}{(\omega/v)^2 - k_z^2} \left(A_+ \exp(i k_z z) - A_- \exp(-i k_z z) \right) \quad k_z L = p\pi, \quad p = 0, 1, 2, \dots$$

BC: $\hat{\mathbf{E}}_\perp(z=0, L) = 0$

$$A_+ = A_- \quad A_+ = A_- e^{-2ik_z L}$$

TE modes

$$\hat{H}_z = H_{z,nm}(x, y) \left(A_+ \exp(i k_z z) + A_- \exp(-i k_z z) \right)$$

$$\hat{\mathbf{E}}_\perp = \frac{-i\omega\mu\hat{\mathbf{z}} \times \nabla_\perp \hat{H}_{z,nm}}{(\omega/v)^2 - k_z^2} \left(A_+ \exp(i k_z z) + A_- \exp(-i k_z z) \right) \quad k_z L = p\pi, \quad p = 1, 2, \dots$$

$$A_+ = -A_- \quad A_+ = -A_- e^{-2ik_z L}$$

Resonant Frequencies

$$k_{\parallel} = \frac{\pi p}{L}$$

$p \neq 0$ not allowed
all fields zero

$$\frac{\omega^2}{c^2} \epsilon \mu = k_c^2 + \left(\frac{\pi p}{L} \right)^2$$

determined by cross section

$$\omega_{\text{Res}} = \frac{c}{\epsilon \mu} \sqrt{k_c^2 + \left(\frac{\pi p}{L} \right)^2}$$

Rectangular cross section

TE_{nmp}
cavite mode

$$\omega_{\text{Res}} = \omega_{nmp} = \frac{c}{\epsilon \mu} \sqrt{\left(\frac{n\pi}{a} \right)^2 + \left(\frac{m\pi}{b} \right)^2 + \left(\frac{p\pi}{L} \right)^2}$$

TM Modes

$$E_z = \frac{1}{2} \left\{ \hat{E}_{\parallel}(x_2) \left[A_+ e^{i(k_{\parallel}z - \omega t)} + A_- e^{-i(k_{\parallel}z - \omega t)} \right] + c.c. \right\}$$

$$\hat{E}_{\perp} = \frac{1}{2} \left\{ i \frac{k_{\parallel} \nabla_{\perp}}{k_c^2} \hat{E}_{\parallel} \left[A_+ e^{i(k_{\parallel}z - \omega t)} - A_- e^{-i(k_{\parallel}z - \omega t)} \right] \right\}$$

at $z=0$ $A_+ - A_- = 0$ $A_+ = A_-$

at $z=L$ $A_+ e^{ik_{\parallel}L} = A_- e^{-ik_{\parallel}L}$

again

$$2k_{\parallel}L = 2\pi p \quad p=0 \quad \text{O.K.}$$

Cavity Losses

$$Q = \frac{\omega U}{P_d} \leftarrow \begin{array}{l} \text{energy stored} \\ \text{Power dissipated} \end{array}$$

Quality Factor

Poynting Theorem

$$\frac{\partial}{\partial t} \int d^3x \frac{1}{4} (\epsilon |\hat{E}|^2 + \mu |\hat{H}|^2) + \int_s dA n \cdot \frac{1}{2} \operatorname{Re} \{ \hat{E} \times \hat{H} \} = 0$$

$\underbrace{\phantom{d^3x \frac{1}{4} (\epsilon |\hat{E}|^2 + \mu |\hat{H}|^2) + \int_s dA n \cdot \frac{1}{2} \operatorname{Re} \{ \hat{E} \times \hat{H} \}}_U}$

Power
to walls



For cavity modes

$$\int d^3x \frac{\epsilon |\hat{E}|^2}{4} = \int d^3x \frac{\mu |\hat{H}|^2}{4}$$

average energy
stored in $E \times H$
equal

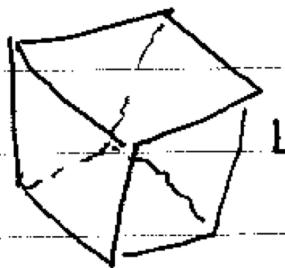
$$P_d = \int dA \frac{1}{2} R_s |\hat{H}_s|^2$$

$$Q = \frac{\omega}{c} \sqrt{\epsilon} \frac{1}{R_s} \frac{\int d^3x |\hat{H}|^2}{\int dA |\hat{H}_s|^2}$$

Weyl's Formula

How many modes in a cavity of volume V
have $\omega_{res} < \omega$?

Consider a cubic cavity of side L $V=L^3$



Resonant Frequencies

$$\underline{\omega_n} = \frac{\pi c}{L} \sqrt{n_x^2 + n_y^2 + n_z^2}$$

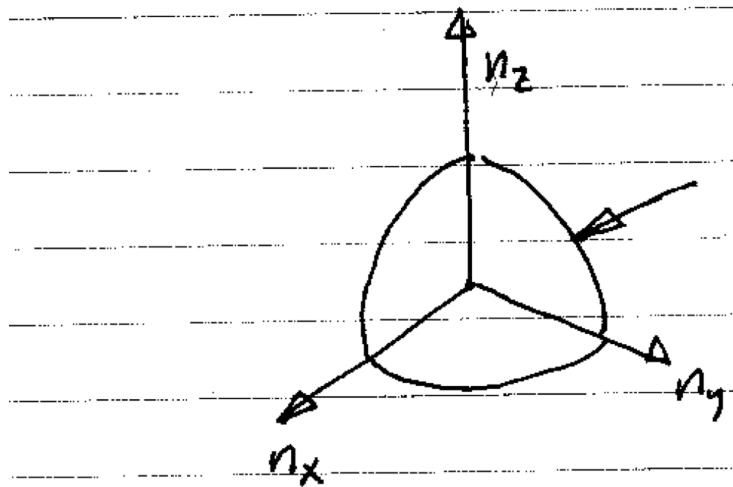
$$\underline{n} = (n_x, n_y, n_z)$$

Estimate of Number of Modes

How many combinations of integers (n_x, n_y, n_z)

have

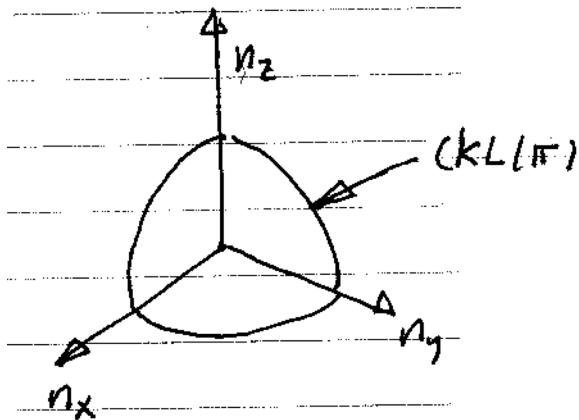
$$n_x^2 + n_y^2 + n_z^2 \leq \left(\frac{wL}{\pi c}\right)^2 = \left(\frac{kL}{\pi}\right)^2 \quad w_n = \frac{\pi c}{L} \sqrt{n_x^2 + n_y^2 + n_z^2}$$



Spherical surface
or radius (kL/π)

Each combination
occupies a cube of
unit volume

Volume Inside Spherical Surface



$$N = \frac{1}{8} \frac{4}{3} \pi \left(\frac{kL}{\pi} \right)^3$$

fraction of sphere

$$N(k) = \frac{(kL)^3}{6\pi^2} = \frac{k^3 V}{6\pi^2}$$

But wait!

For each set of integers

there are 2 polarizations

For EM modes

$$N(k) = \frac{1}{3} \frac{k^3 V}{\pi^2}$$

Example

$$\text{Volume} = 1 \text{ m}^3$$

$$f = 1 \text{ GHz}$$

$$k = \frac{2\pi f}{c} = \frac{2\pi \times 10^9}{3 \times 10^8} = 21 \text{ m}^{-1}$$

$$N(k) \approx 310$$

What is the typical spacing

$$\delta k = \text{spacing in } k = \omega/c$$

$$N(k+\delta k) \approx N(k) + 1$$

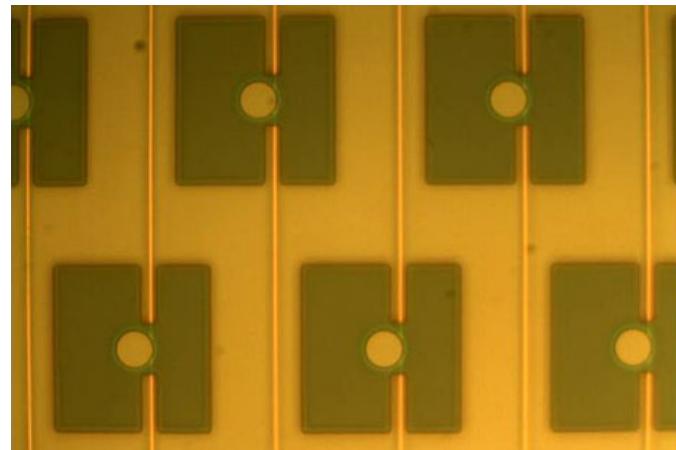
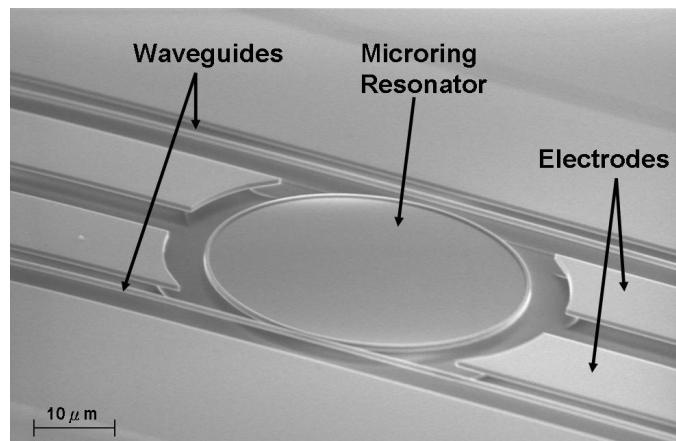
$$N(k+\delta k) \approx N(k) + \frac{\delta k}{k} \frac{dN}{dk}$$

$$\frac{dN}{dk} = \frac{k^2 V}{\pi^2}$$

$$\text{fractional Spacing} \quad \frac{\delta k}{k} = \frac{1}{k \frac{dN}{dk}} = \frac{\pi^2}{k^3 V} \approx 1.07 \times 10^{-3}$$

Integrated photonics

(Courtesy Edo Waks)

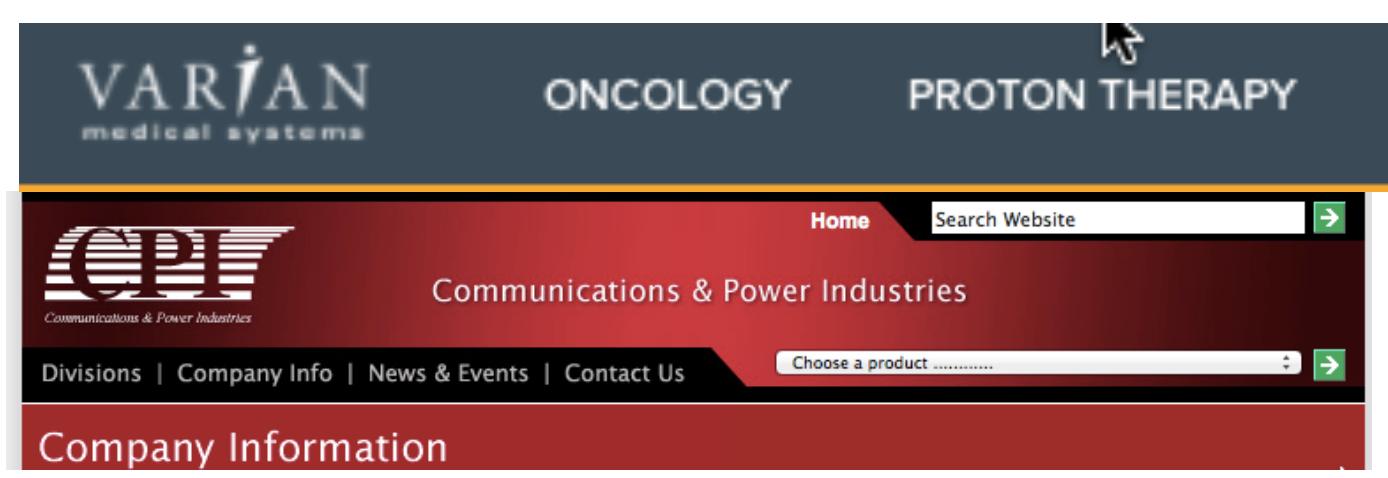
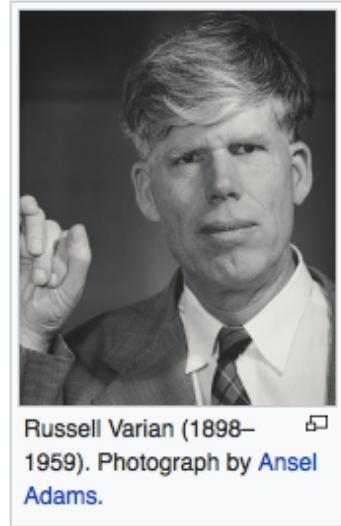
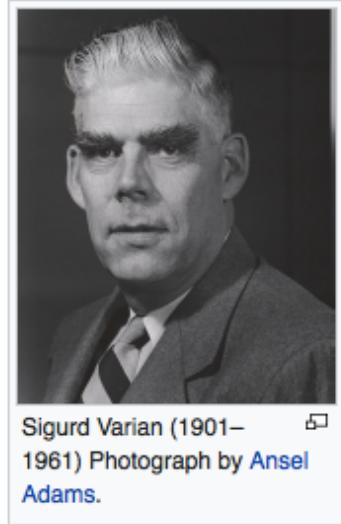


Klystron – Beam Driven HPM source

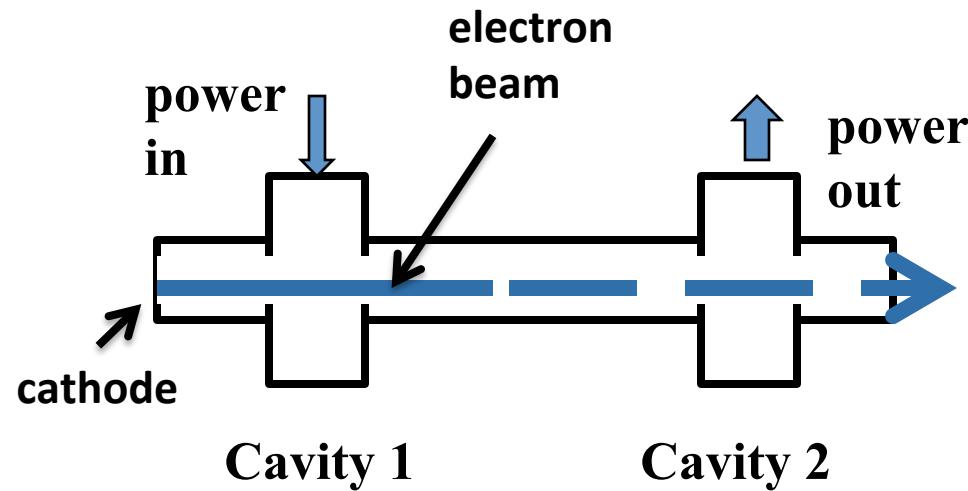
Klystron: invented in 1937 by the Varian brothers. One of the first Palo Alto High Tech. firms.

High Power Source of Microwaves

Radar, Particle Accelerators, (LHC 16 x 300 kW), etc

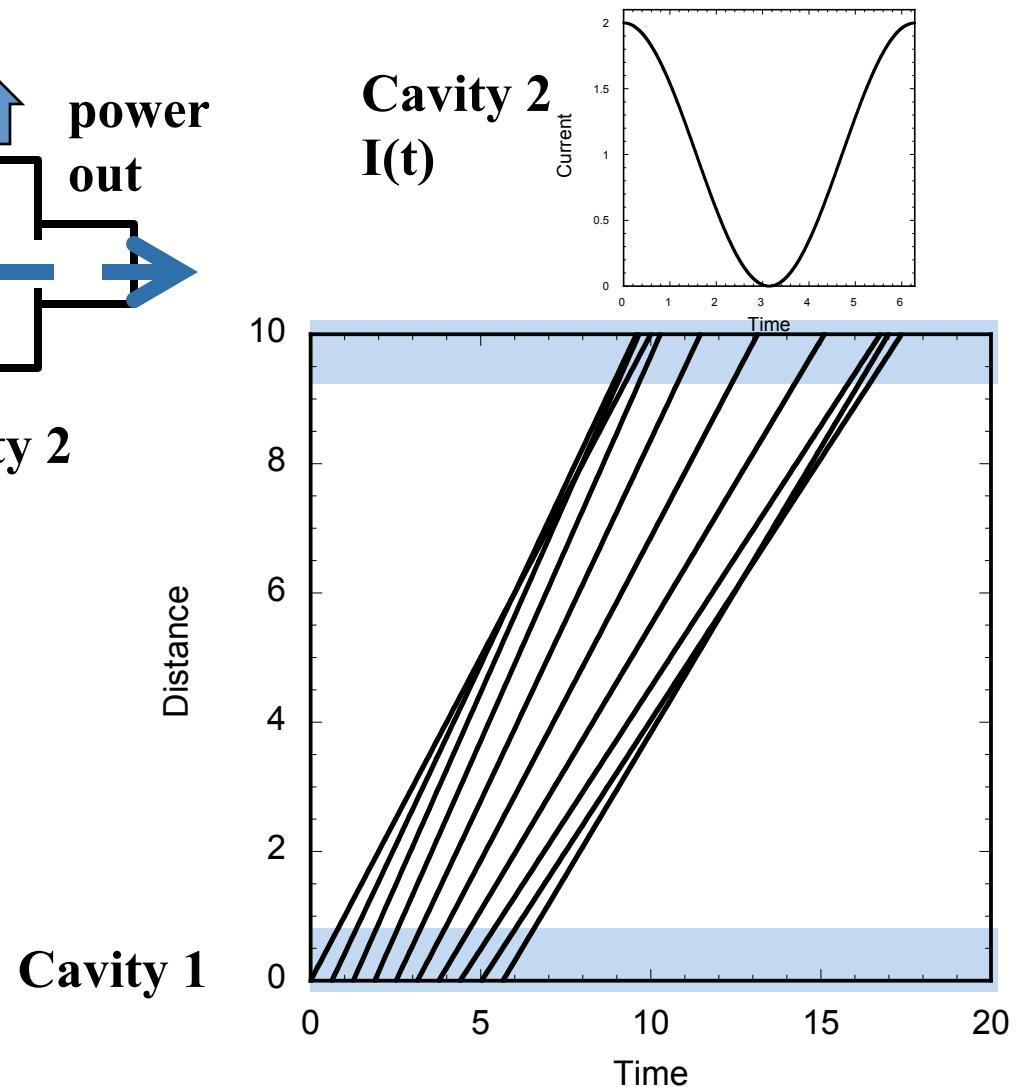
A screenshot of the Varian medical systems website. The top navigation bar includes links for 'ONCOLOGY' and 'PROTON THERAPY'. Below this is a red banner featuring the CPI logo ('Communications & Power Industries') and links for 'Divisions | Company Info | News & Events | Contact Us'. A search bar and a dropdown menu for 'Choose a product' are also visible.

Velocity Modulation Ballistic Bunching



Field in cavity 1 gives small time dependent velocity modulation

Fast electrons catch up to slow electrons giving large current modulation.



Examples

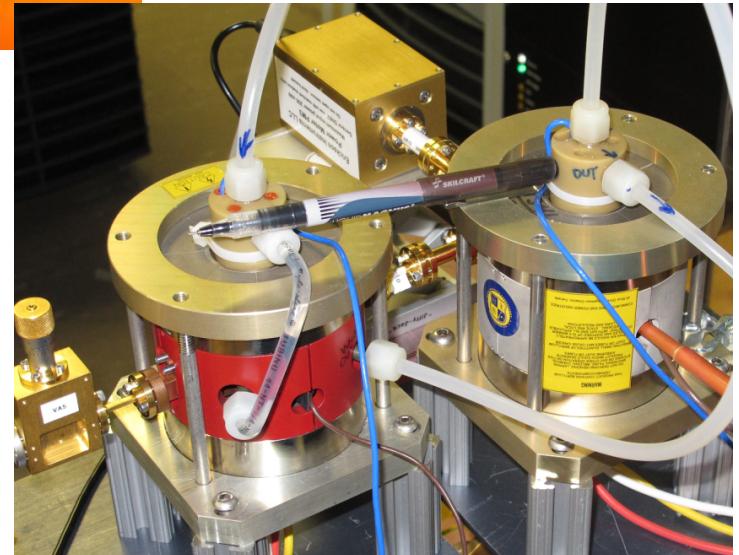


Monica
Blank

170 GHz CPI Gyrotron
IEEE IVEC
<http://ieeexplore.ieee.org>



**L3 Ka Band
Power Module**
[http://
www.linkmicrot
ek.com](http://www.linkmicrotek.com)



Experimental high power set-up showing the CPI 218.4 GHz EIK driving the compact NRL Serpentine Waveguide (SWG) TWT.