

# Dielectric Waveguides

ENEE 381

Basic principles

Slab model

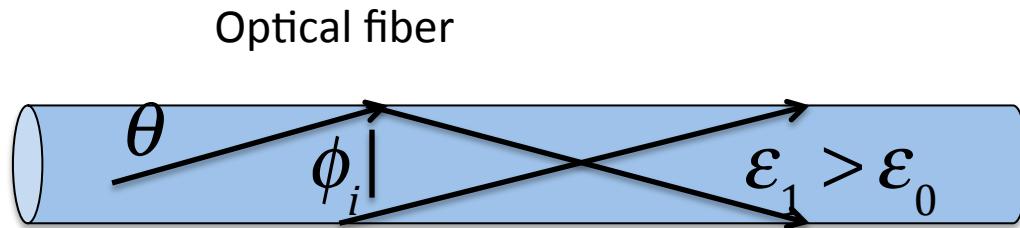
two regions

boundary conditions

dispersion relation – odd and even modes

Cylindrical geometry

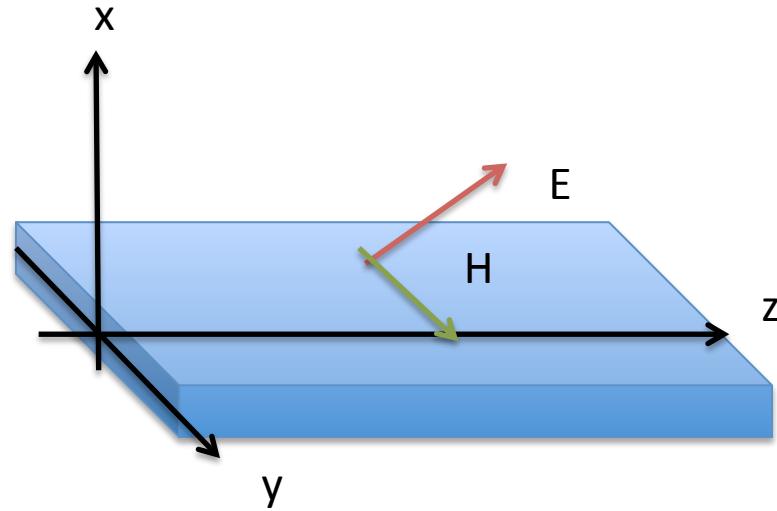
# Guiding by total internal reflection



Total reflection if

$$\sin \phi_i = \cos \theta > \sqrt{\varepsilon_1 / \varepsilon_0}$$

# Slab Model

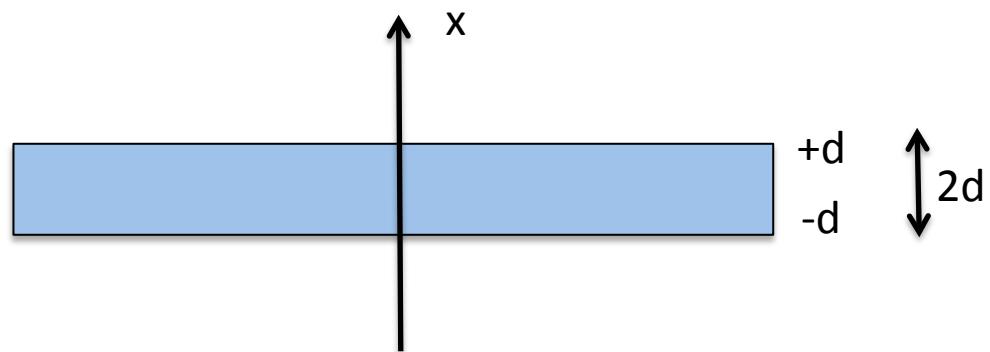


$$\epsilon = \begin{cases} \epsilon_1 & |x| < d \\ \epsilon_0 & |x| > d \end{cases}$$

Look for solutions

$$\mathbf{E} = \text{Re} \left\{ (\hat{E}_x, 0, \hat{E}_z) \exp \left[ i(k_z z - \omega t) \right] \right\}$$

$$\mathbf{H} = \text{Re} \left\{ (0, \hat{H}_y, 0) \exp \left[ i(k_z z - \omega t) \right] \right\}$$



# Maxwell's Equations

$$\nabla \times \vec{\mathbf{E}} = -\mu \frac{\partial \vec{\mathbf{H}}}{\partial t}$$

$$\frac{\partial \hat{E}_z}{\partial y} - ik_z \hat{E}_y = i\omega \mu \hat{H}_x$$

$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega \mu \hat{H}_y$$

$$\frac{\partial \hat{E}_y}{\partial x} - \frac{\partial \hat{E}_x}{\partial y} = i\omega \mu \hat{H}_z$$

$$\nabla \times \vec{\mathbf{H}} = \epsilon \frac{\partial \vec{\mathbf{E}}}{\partial t}$$

$$\frac{\partial \hat{H}_z}{\partial y} - ik_z \hat{H}_y = -i\omega \epsilon \hat{E}_x$$

$$ik_z \hat{H}_x - \frac{\partial \hat{H}_z}{\partial x} = -i\omega \epsilon \hat{E}_y$$

$$\frac{\partial \hat{H}_y}{\partial x} - \frac{\partial \hat{H}_x}{\partial y} = -i\omega \epsilon \hat{E}_z$$

Look for solutions

$$\mathbf{E} = \text{Re} \left\{ (\hat{E}_x, 0, \hat{E}_z) \exp \left[ i(k_z z - \omega t) \right] \right\}$$

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# Maxwell's Equations

$$\nabla \times \vec{\mathbf{E}} = -\mu \frac{\partial \vec{\mathbf{H}}}{\partial t}$$

$$\frac{\partial \hat{E}_z}{\partial y} - ik_z \hat{E}_y = i\omega \mu \hat{H}_x \rightarrow$$

$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega \mu \hat{H}_y$$

$$\frac{\partial \hat{E}_y}{\partial x} - \frac{\partial \hat{E}_x}{\partial y} = i\omega \mu \hat{H}_z \rightarrow$$

$$\nabla \times \vec{\mathbf{H}} = \epsilon \frac{\partial \vec{\mathbf{E}}}{\partial t}$$

$$\frac{\partial \hat{H}_z}{\partial y} - ik_z \hat{H}_y = -i\omega \epsilon \hat{E}_x$$

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$$\frac{\partial \hat{H}_y}{\partial x} - \frac{\partial \hat{H}_x}{\partial y} = -i\omega \epsilon \hat{E}_z$$

Look for solutions

$$\mathbf{E} = \text{Re} \left\{ (\hat{E}_x, 0, \hat{E}_z) \exp \left[ i(k_z z - \omega t) \right] \right\}$$

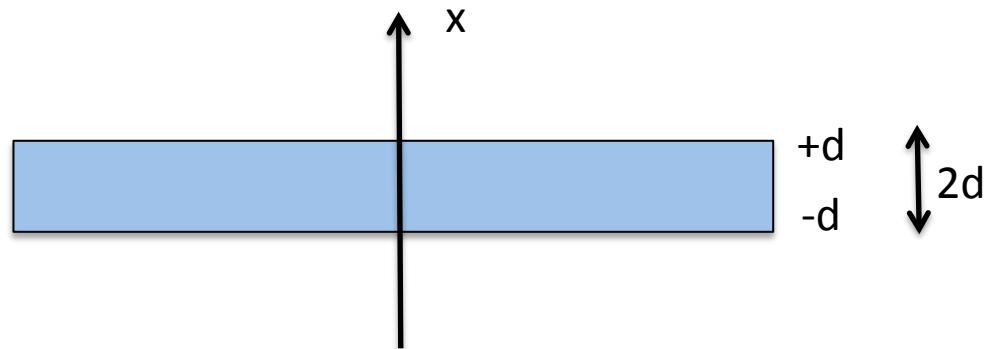
$$\mathbf{H} = \text{Re} \left\{ (0, \hat{H}_y, 0) \exp \left[ i(k_z z - \omega t) \right] \right\}$$

# Maxwell's Equations

$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega\mu \hat{H}_y \quad -ik_z \hat{H}_y = -i\omega\epsilon \hat{E}_x$$
$$\frac{\partial \hat{H}_y}{\partial x} = -i\omega\epsilon \hat{E}_z$$

$$\epsilon = \begin{cases} \epsilon_1 & |x| < d \\ \epsilon_0 & |x| > d \end{cases}$$

What are the boundary conditions at  $x = +/- d$ ?



# Maxwell's Equations

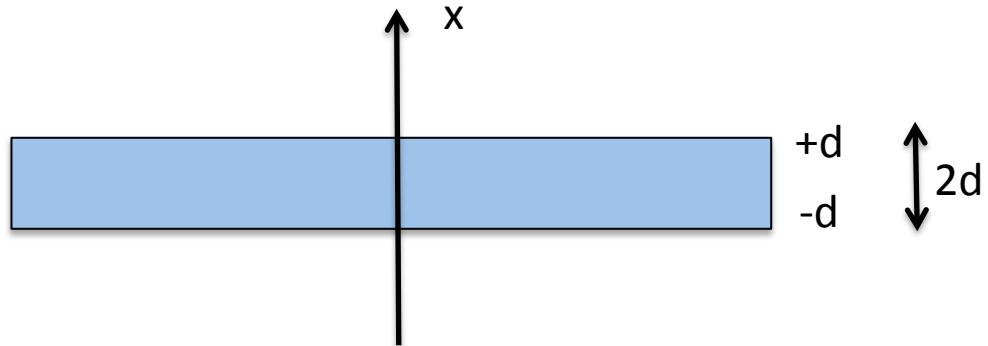
$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega\mu \hat{H}_y$$

$$\begin{aligned}-ik_z \hat{H}_y &= -i\omega\epsilon \hat{E}_x \\ \frac{\partial \hat{H}_y}{\partial x} &= -i\omega\epsilon \hat{E}_z\end{aligned}$$

$$\epsilon = \begin{cases} \epsilon_1 & |x| < d \\ \epsilon_0 & |x| > d \end{cases}$$

What are the boundary conditions at  $x = +/- d$ ?

$$\left. \begin{array}{l} \hat{E}_z \\ \hat{H}_y \end{array} \right\} \text{continuous at } x = \pm d$$



# Combine to a single equation

The diagram illustrates the derivation of a single wave equation from three initial equations. A blue box contains the first three equations:

$$ik_z \hat{E}_x - \frac{\partial \hat{E}_z}{\partial x} = i\omega \mu \hat{H}_y$$

$$-ik_z \hat{H}_y = -i\omega \epsilon \hat{E}_x$$

$$\frac{\partial \hat{H}_y}{\partial x} = -i\omega \epsilon \hat{E}_z$$

Blue arrows point from each of these three equations to a central result:

$$k_z^2 \hat{E}_x + ik_z \frac{\partial \hat{E}_z}{\partial x} = i\omega \mu (-ik_z \hat{H}_y) = \omega^2 \epsilon \mu \hat{E}_x$$

$$\hat{E}_x = \frac{ik_z}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

$$\hat{H}_y = \frac{i\omega \epsilon}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

A blue arrow points from the third equation in the box to the fourth equation below it:

$$\frac{\partial \hat{H}_y}{\partial x} = -i\omega \epsilon \hat{E}_z$$

An arrow points from this fourth equation to the final result:

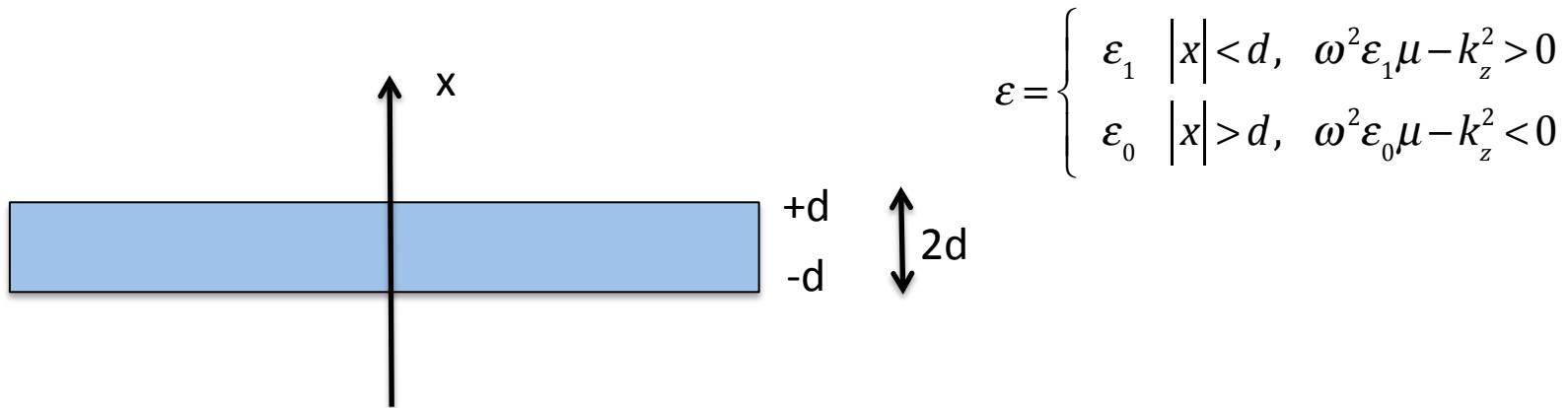
$$\frac{\partial}{\partial x} \frac{\epsilon}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x} = -\epsilon \hat{E}_z$$

If  $\epsilon$  is constant

$$\frac{\partial}{\partial x} \frac{\partial \hat{E}_z}{\partial x} + (\omega^2 \epsilon \mu - k_z^2) \hat{E}_z = 0$$

# Solve in two regions

If  $\epsilon$  is constant  $\frac{\partial}{\partial x} \frac{\partial \hat{E}_z}{\partial x} + (\omega^2 \epsilon \mu - k_z^2) \hat{E}_z = 0$

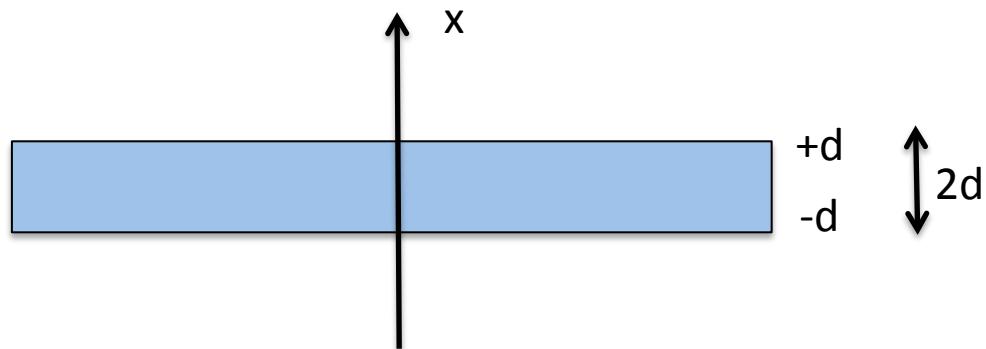


For  $x > d$ ,  $\epsilon = \epsilon_0$   $\hat{E}_z = A \exp(-\kappa x)$   $\kappa = \sqrt{k_z^2 - \omega^2 \epsilon_0 \mu}$

Evanescent (Spatially decaying) field

# Solve in two regions

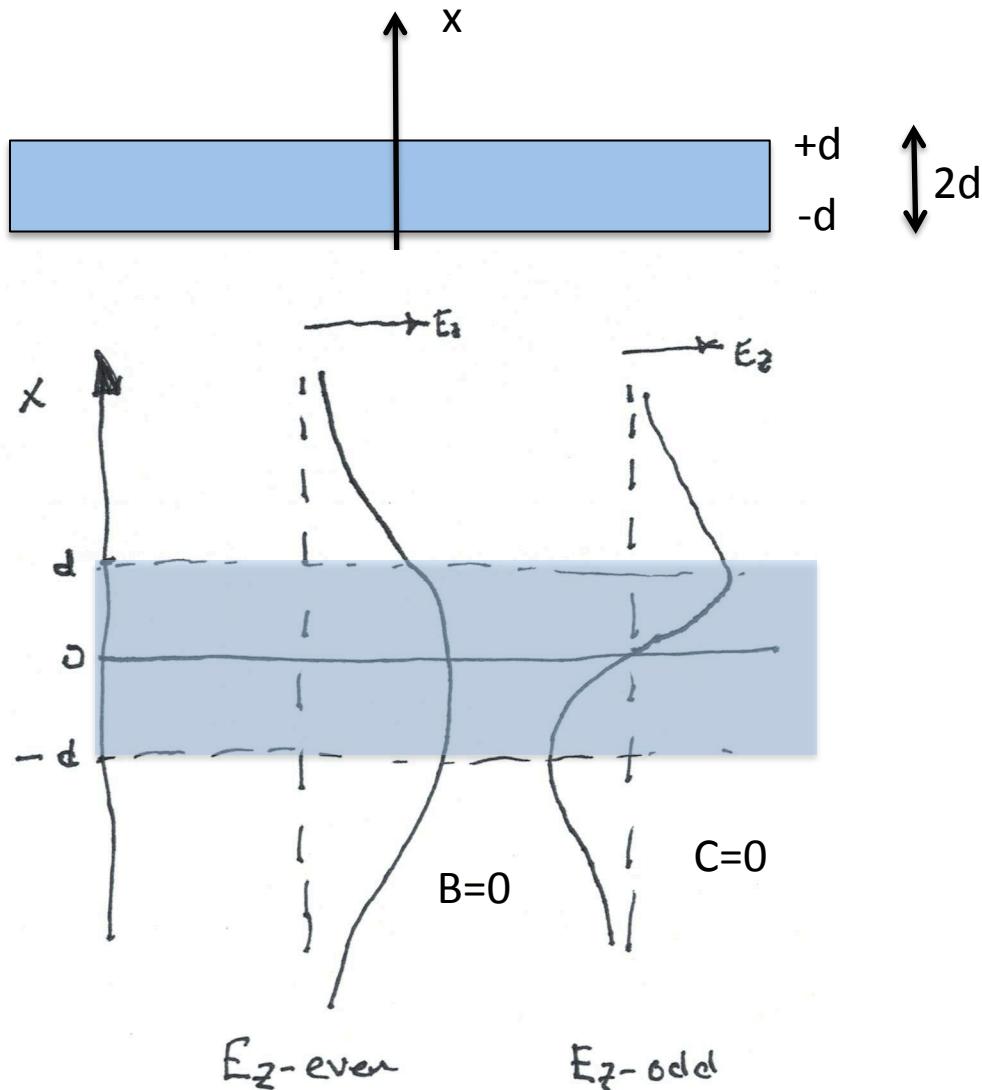
If  $\epsilon$  is constant  $\frac{\partial}{\partial x} \frac{\partial \hat{E}_z}{\partial x} + (\omega^2 \epsilon \mu - k_z^2) \hat{E}_z = 0$



For  $|x| < d$ ,  $\epsilon = \epsilon_1$   $\hat{E}_z = B \sin(k_x x) + C \cos(k_x x)$   $k_x = \sqrt{\omega^2 \epsilon_1 \mu - k_z^2}$

Spatially oscillating field

# Odd and Even solutions



For  $|x| < d$ ,  $\epsilon = \epsilon_1$

$$\hat{E}_z = B \sin(k_x x) + C \cos(k_x x)$$

$$k_x = \sqrt{\omega^2 \epsilon_1 \mu - k_z^2}$$

Either:  $B=0, C \neq 0$

or  $B \neq 0, C = 0$

$$\hat{E}_x = \frac{ik_z}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

$$\hat{H}_y = \frac{i\omega \epsilon}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

# $E_z$ -even ( $E_x$ - odd) solution

$$\text{For } |x| < d, \quad \hat{E}_z = C \cos(k_x x) \quad k_x = \sqrt{\omega^2 \epsilon_1 \mu - k_z^2} \quad \hat{E}_x = \frac{ik_z}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

$$\text{For } x > d, \quad \hat{E}_z = A \exp(-\kappa x) \quad \kappa = \sqrt{k_z^2 - \omega^2 \epsilon_0 \mu} \quad \hat{H}_y = \frac{i\omega \epsilon}{\omega^2 \epsilon \mu - k_z^2} \frac{\partial \hat{E}_z}{\partial x}$$

$$\text{At } x=d \quad \hat{E}_z = C \cos(k_x d) = A \exp(-\kappa d)$$

$$\hat{H}_y = \frac{i\omega \epsilon_1}{k_x^2} \frac{\partial}{\partial x} C \cos(k_x x) \Big|_{x=d} = -\frac{i\omega \epsilon_1}{k_x} C \sin(k_x d)$$

$$\hat{H}_y = \frac{i\omega \epsilon_0}{-\kappa^2} \frac{\partial}{\partial x} A \exp(-\kappa x) \Big|_{x=d} = \frac{i\omega \epsilon_0}{\kappa} A \exp(-\kappa d)$$

# Dispersion Relation

$$C \cos(k_x d) = A \exp(-\kappa d)$$

$$-\frac{\epsilon_1}{k_x} C \sin(k_x d) = \frac{\epsilon_0}{\kappa} A \exp(-\kappa d)$$

Even E<sub>z</sub>

$$-\frac{\epsilon_1}{\epsilon_0} \kappa d = k_x d \cot(k_x d)$$

Odd E<sub>z</sub>

$$\frac{\epsilon_1}{\epsilon_0} \kappa d = k_x d \tan(k_x d)$$

$$\kappa d = d \sqrt{k_z^2 - \omega^2 \epsilon_0 \mu} = \sqrt{\omega^2 (\epsilon_1 - \epsilon_0) \mu d^2 - k_x^2 d^2}$$

$$\kappa d = \sqrt{\omega^2 (\epsilon_1 - \epsilon_0) \mu d^2 - k_x^2 d^2}$$

# Plot both sides versus $k_x d$

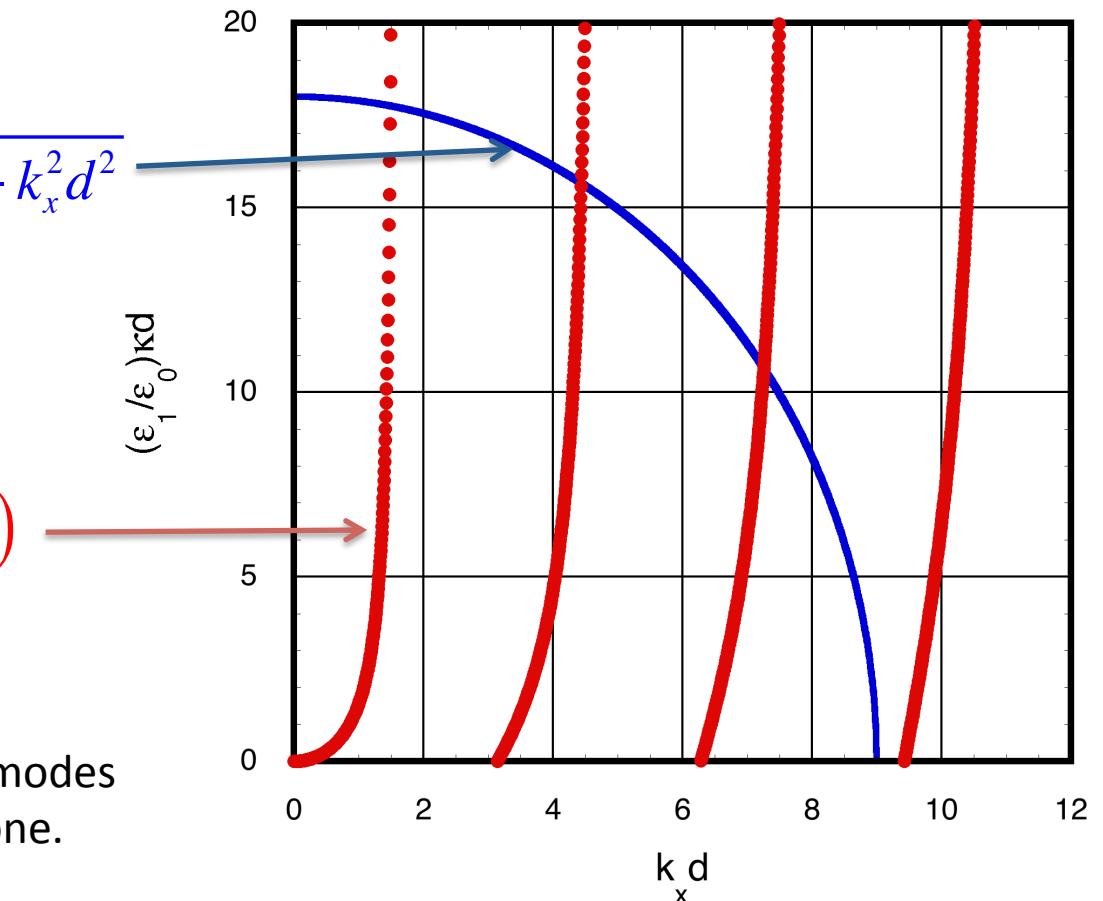
Odd  $E_z$

$$\frac{\epsilon_1}{\epsilon_0} \kappa d = k_x d \tan(k_x d)$$

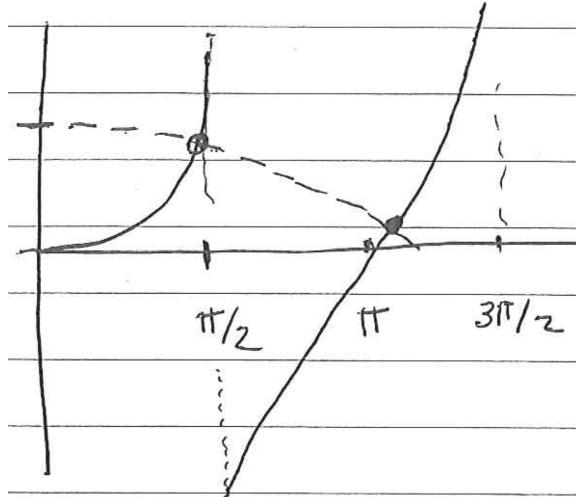
$$\frac{\epsilon_1}{\epsilon_0} \kappa d = \frac{\epsilon_1}{\epsilon_0} \sqrt{\omega^2 (\epsilon_1 - \epsilon_0) \mu d^2 - k_x^2 d^2}$$

$$\frac{\epsilon_1}{\epsilon_0} \kappa d = k_x d \tan(k_x d)$$

In this case three solutions – modes  
There will always be at least one.

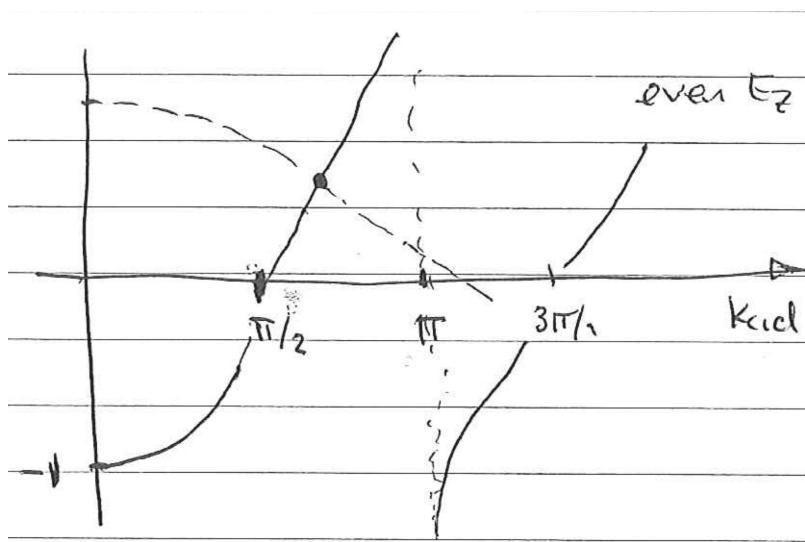


# Comments



Odd Ez –Even Ex

Always at least one solution  
For small diameter d fields outside dielectric  
Requires cladding

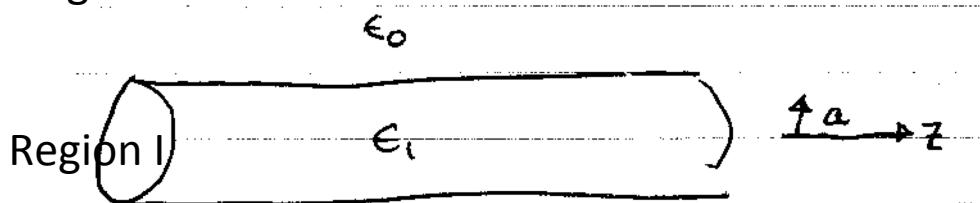


Even Ez –Odd Ex

Possibly no solutions

# Cylindrical Dielectrics

Region II



$$\text{Field} \sim \{\hat{E}, \hat{H}(r, \theta)\} \exp[ik_z z - i\omega t]$$

Express transverse components in terms of axial components

$$\hat{E}_\perp = \frac{1}{k_z^2} [ik_z \nabla_\perp \hat{E}_z - i\omega \mu \epsilon_2 \times \nabla_\perp \hat{H}_z]$$

$$\hat{H}_\perp = \frac{1}{k_z^2} [ik_z \nabla_\perp \hat{H}_z + i\omega \mu \epsilon_2 \times \nabla_\perp \hat{E}_z] \quad k_z^2 = \omega^2 \epsilon \mu - k_x^2 \quad - \text{different in each region}$$

# Equation for Axial Components

$$\nabla_1^2 \hat{E}_z + k_1^2 \hat{E}_z = 0 \quad \nabla_2^2 \hat{H}_z + k_2^2 \hat{H}_z = 0$$

$$k_1^2 = \omega^2 \epsilon \mu - k_z^2 \quad - \text{different in each region}$$

Region I:  $k_{11}^2 = \omega^2 \epsilon_1 \mu - k_z^2 > 0$

Region II  $k_{12}^2 = \omega^2 \epsilon_2 \mu - k_z^2 = -\gamma_2^2 < 0$

$$\hat{E}_z = A_1 \exp(i \ell \theta) J_\ell(k_{11} r)$$

$\ell = \text{integer}$

$$\hat{E}_z(r, \theta) = A_2 e^{i \ell \theta} K_\ell(\gamma_2 r)$$

$$\hat{H}_z = B_1 \exp(i \ell \theta) J_\ell(k_{11} r)$$

$$\hat{H}_z(r, \theta) = B_2 e^{i \ell \theta} K_\ell(\gamma_2 r)$$

Ordinary Bessel functions  
Regular at origin

Modified Bessel functions  
Goes to zero at infinity

# Unknowns and Boundary Conditions

$$\hat{E}_z = A_1 \exp(i\ell\theta) J_\ell(k_1 r)$$

$$\hat{E}_z(r, \theta) = A_2 e^{i\ell\theta} K_\ell(\gamma_2 r)$$

$$\hat{H}_z = B_1 \exp(i\ell\theta) J_\ell(k_1 r)$$

$$\hat{H}_z(r, \theta) = B_2 e^{i\ell\theta} K_\ell(\gamma_2 r)$$

Unknowns  $A_1, A_2, B_1, B_2$  & w for given  $k_z$

Boundary Conditions: at  $r=a$

Tangential components of E and H continuous

$E_z, E_\theta, H_z, H_\theta$  - continuous at  $r=a$

$H_z, E_z$  - continuous

$$A_1 J_\ell(k_1 a) = A_2 K_\ell(\gamma_2 a)$$

$$B_1 J_\ell(k_1 a) = B_2 K_\ell(\gamma_2 a)$$

# Continuity of Transverse Components

Transverse components in terms  
of axial components

$$\hat{E}_z = \frac{1}{k_z^2} [ i k_z \nabla_{\perp} \hat{E}_z - i w \mu \epsilon_z \times \nabla_{\perp} \hat{H}_z ]$$

$$\hat{H}_z = \frac{1}{k_z^2} [ i k_z \nabla_{\perp} \hat{H}_z + i w \epsilon \epsilon_z \times \nabla_{\perp} \hat{E}_z ]$$

Region I

$$\hat{E}_0(a) = \frac{1}{k_1^2} \left\{ i k_z \left( \frac{i\ell}{a} \right) A_1 J_\ell(k_1 a) - i w \mu B_1 k_1 J_\ell'(k_1 a) \right\}$$

$$H_0(r=a) = \frac{1}{k_1^2} \left\{ i k_z \left( \frac{i\ell}{a} \right) B_1 J_\ell(k_1 a) + i w \epsilon k_1 A_1 J_\ell'(k_1 a) \right\}$$

Region II

$$\hat{E}_0(a) = -\frac{1}{\gamma_2^2} \left\{ i k_z \left( \frac{i\ell}{a} \right) A_2 K_\ell(\gamma_2 a) - i w \mu \gamma_2 B_2 K_\ell'(\gamma_2 a) \right\}$$

$$k_2^2 = -\gamma_2^2$$

$$H_0(a) = -\frac{1}{\gamma_2^2} \left\{ i k_z \left( \frac{i\ell}{a} \right) B_2 K_\ell(\gamma_2 a) + i w \epsilon \gamma_2 A_2 K_\ell'(\gamma_2 a) \right\}$$

# Solution

$$A_1 J_e(k_1 a) = A_2 K_e(\gamma_2 a)$$

$$B_1 J_e(k_1 a) = B_2 K_e(\gamma_2 a)$$



$$\hat{E}_e(a) = \frac{1}{k_{21}^2} \left\{ ik_2 \left( i \frac{\ell}{a} \right) A_1 J_e(k_1 a) - iw\mu B_1 k_1 J_e'(k_1 a) \right\}$$

$$H_D(r=a) = \frac{1}{k_{21}^2} \left\{ ik_2 \left( i \frac{\ell}{a} \right) B_1 J_e(k_1 a) + iw\epsilon k_1 A_1 J_e'(k_1 a) \right\}$$

$$\hat{E}_0(a) = -\frac{1}{\gamma_2^2} \left\{ ik_2 \left( i \frac{\ell}{a} \right) A_2 K_e(\gamma_2 a) - iw\mu \gamma_2 B_2 K_e'(\gamma_2 a) \right\}$$

$$H_0(a) = -\frac{1}{\gamma_2^2} \left\{ ik_2 \left( i \frac{\ell}{a} \right) B_2 K_e(\gamma_2 a) + iw\epsilon \gamma_2 A_2 K_e'(\gamma_2 a) \right\}$$

Replace  $A_2 = A_1 J_e / K_e$      $B_2 = B_1 J_e / K_e$

$$\begin{bmatrix} 2 \times 2 \\ M \end{bmatrix} \begin{pmatrix} A_1 \\ B_1 \end{pmatrix} = 0$$

solution  $\det(M) = 0$  determines  $w(k)$