

# ENEE381

Plane Electromagnetic Waves

In One Dimension

Impedance, Sinusoidal Waves, Phasors

# Where We Stand

$$\nabla \cdot \vec{\mathbf{D}} = \rho_f$$

$$\nabla \cdot \vec{\mathbf{B}} = 0$$

$$\nabla \times \vec{\mathbf{E}} = -\frac{\partial \vec{\mathbf{B}}}{\partial t}$$

$$\nabla \times \vec{\mathbf{H}} = \vec{\mathbf{J}}_f + \frac{\partial \vec{\mathbf{D}}}{\partial t}$$

Assume the following:

Linear, isotropic,  
instantaneous, media

$$\vec{\mathbf{D}} = \epsilon \vec{\mathbf{E}}$$

$$\vec{\mathbf{B}} = \mu \vec{\mathbf{H}}$$

Propagation in free  
space, no free charge or  
current.

$$\rho_f = 0, \quad \vec{\mathbf{J}}_f = 0$$

# Write out component by component

$$\varepsilon \nabla \cdot \vec{\mathbf{E}} = 0$$

$$\varepsilon \left( \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \right) = 0$$

$$\nabla \times \vec{\mathbf{E}} = -\mu \frac{\partial \vec{\mathbf{H}}}{\partial t}$$

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = -\mu \frac{\partial H_x}{\partial t}$$

$$\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = -\mu \frac{\partial H_y}{\partial t}$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = -\mu \frac{\partial H_z}{\partial t}$$

$$\mu \nabla \cdot \vec{\mathbf{H}} = 0$$

$$\mu \left( \frac{\partial H_x}{\partial x} + \frac{\partial H_y}{\partial y} + \frac{\partial H_z}{\partial z} \right) = 0$$

$$\nabla \times \vec{\mathbf{H}} = \varepsilon \frac{\partial \vec{\mathbf{E}}}{\partial t}$$

$$\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} = \varepsilon \frac{\partial E_x}{\partial t}$$

$$\frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} = \varepsilon \frac{\partial E_y}{\partial t}$$

$$\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} = \varepsilon \frac{\partial E_z}{\partial t}$$

# “Look” for solutions in the form

$$\vec{\mathbf{E}} = (E_x(z,t), E_y = 0, E_z = 0) \quad \vec{\mathbf{H}} = (H_x = 0, H_y(z,t), H_z = 0)$$

Then verify that all  $2 \times 3 + 2 = 8$  equations can be satisfied

$$\epsilon \nabla \cdot \vec{\mathbf{E}} = 0$$

$$\epsilon \left( \frac{\partial E_x}{\partial x} + \frac{\partial 0}{\partial y} + \frac{\partial 0}{\partial z} \right) = 0$$


$$0$$

$$\mu \nabla \cdot \vec{\mathbf{H}} = 0$$

$$\mu \left( \frac{\partial 0}{\partial x} + \frac{\partial H_y}{\partial y} + \frac{\partial 0}{\partial z} \right) = 0$$


$$0$$

$$\vec{\mathbf{E}} = (E_x(z,t), E_y = 0, E_z = 0)$$

$$\vec{\mathbf{H}} = (H_x = 0, H_y(z,t), H_z = 0)$$

Insert fields

$$\nabla \times \vec{\mathbf{E}} = -\mu \frac{\partial \vec{\mathbf{H}}}{\partial t}$$

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = -\mu \frac{\partial H_x}{\partial t}$$

$$\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = -\mu \frac{\partial H_y}{\partial t}$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = -\mu \frac{\partial H_z}{\partial t}$$



$$\nabla \times \vec{\mathbf{E}} = -\mu \frac{\partial \vec{\mathbf{H}}}{\partial t}$$

$$\frac{\partial 0}{\partial y} - \frac{\partial 0}{\partial z} = -\mu \frac{\partial 0}{\partial t}$$

OK!

$$\frac{\partial E_x}{\partial z} - \frac{\partial 0}{\partial x} = -\mu \frac{\partial H_y}{\partial t}$$

$$\frac{\partial 0}{\partial x} - \frac{\partial E_x}{\partial y} = -\mu \frac{\partial 0}{\partial t}$$

0

$$\vec{\mathbf{E}} = (E_x(z,t), E_y = 0, E_z = 0)$$

$$\vec{\mathbf{H}} = (H_x = 0, H_y(z,t), H_z = 0)$$

Insert fields

$$\nabla \times \vec{\mathbf{H}} = \epsilon \frac{\partial \vec{\mathbf{E}}}{\partial t}$$

$$\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} = \epsilon \frac{\partial E_x}{\partial t}$$

$$\frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} = \epsilon \frac{\partial E_y}{\partial t}$$

$$\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} = \epsilon \frac{\partial E_z}{\partial t}$$



$$\nabla \times \vec{\mathbf{H}} = \epsilon \frac{\partial \vec{\mathbf{E}}}{\partial t}$$

$$\frac{\partial 0}{\partial y} - \frac{\partial H_y}{\partial z} = \epsilon \frac{\partial E_x}{\partial t}$$

$$\frac{\partial 0}{\partial z} - \frac{\partial 0}{\partial x} = \epsilon \frac{\partial 0}{\partial t}$$

$$\frac{\partial H_y}{\partial x} - \frac{\partial 0}{\partial y} = \epsilon \frac{\partial 0}{\partial t}$$

OK!

0

# When the dust settles

$$\frac{\partial E_x}{\partial z} = -\mu \frac{\partial H_y}{\partial t} \quad -\frac{\partial H_y}{\partial z} = \epsilon \frac{\partial E_x}{\partial t}$$

Combining

$$\frac{\partial}{\partial z} \frac{\partial E_x}{\partial z} = \frac{\partial^2 E_x}{\partial z^2} = -\mu \frac{\partial}{\partial z} \frac{\partial H_y}{\partial t} = -\mu \frac{\partial}{\partial t} \frac{\partial H_y}{\partial z} = \epsilon \mu \frac{\partial^2}{\partial t^2} E_x$$

The 1D wave equation

$$\frac{\partial^2 E_x}{\partial z^2} = \epsilon \mu \frac{\partial^2}{\partial t^2} E_x = \frac{1}{v^2} \frac{\partial^2}{\partial t^2} E_x \quad v^2 = \frac{1}{\epsilon \mu}$$

Once  $E_x$  is found you can go back and integrate to find  $H_y$

# General Solution of 1D W.E.

$$\frac{\partial^2 E_x}{\partial z^2} = \frac{1}{v^2} \frac{\partial^2}{\partial t^2} E_x$$

$$E_x = E_+(z - vt) + E_-(z + vt)$$

$E_+$  is a pulse moving in the  $+z$  direction

$E_-$  is a pulse moving in the  $-z$  direction

If you move with speed  $v$  such that your position is  $z = vt + z_0$

Then  $E_+$  does not change.

Then using

$$\frac{\partial E_x}{\partial z} = -\mu \frac{\partial H_y}{\partial t}$$

$$H_y = \frac{1}{\eta} [E_+(z - vt) - E_-(z + vt)] \quad v = \frac{1}{\sqrt{\epsilon\mu}} \quad \eta = \sqrt{\frac{\mu}{\epsilon}} - \text{medium's impedance}$$

The most important minus sign this semester

# Plane Wave

$$E_x = E_+(z - vt) + E_-(z + vt)$$

$$H_y = \frac{1}{\eta} [E_+(z - vt) - E_-(z + vt)]$$

$$v = \frac{1}{\sqrt{\epsilon\mu}} \quad \eta = \sqrt{\frac{\mu}{\epsilon}} \text{ - medium's impedance}$$

Poynting Flux

$$S_z = E_x H_y = \frac{1}{\eta} \left[ (E_+(z - vt))^2 - (E_-(z + vt))^2 \right]$$

Forward      Backward

Units

[E] – volts/meter

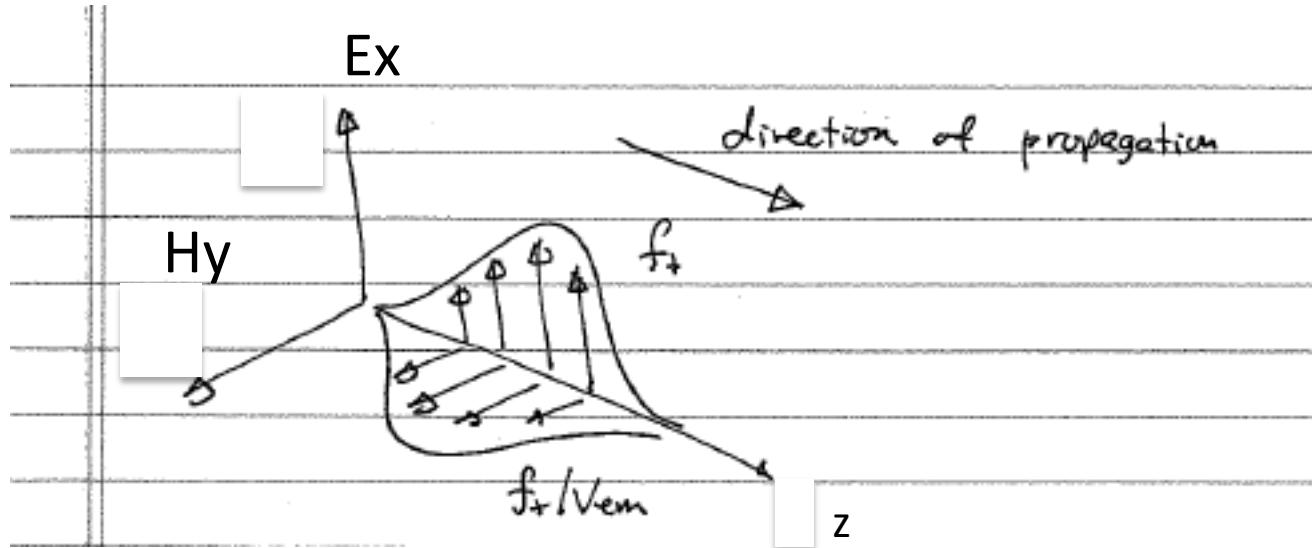
[H] – Amps/meter

$$\eta = \sqrt{\frac{\mu}{\epsilon}} \text{ - [Ohms]}$$

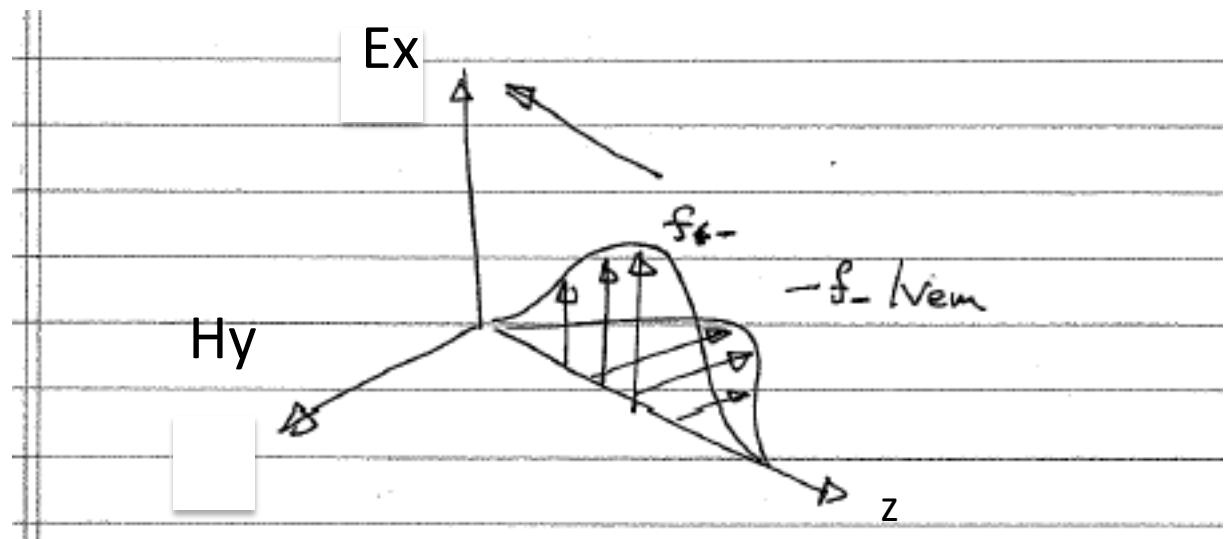
E and B fields in waves and Right Hand Rule:

f+ solution

Wave propagates in  $\mathbf{E} \times \mathbf{B}$  direction



f- solution



## Energy Density and Intensity of EM Waves

Energy density associated with electric and magnetic fields

$$u_E = \frac{\epsilon_0 |\vec{E}|^2}{2} \quad u_B = \frac{|\vec{B}|^2}{2\mu_0}$$

For a wave:  $|\vec{B}| = \frac{1}{v_{em}} |\vec{E}| = \sqrt{\epsilon_0 \mu_0} |\vec{E}|$

Thus:

$$u_E = u_B \quad \text{Units: J/m}^3$$

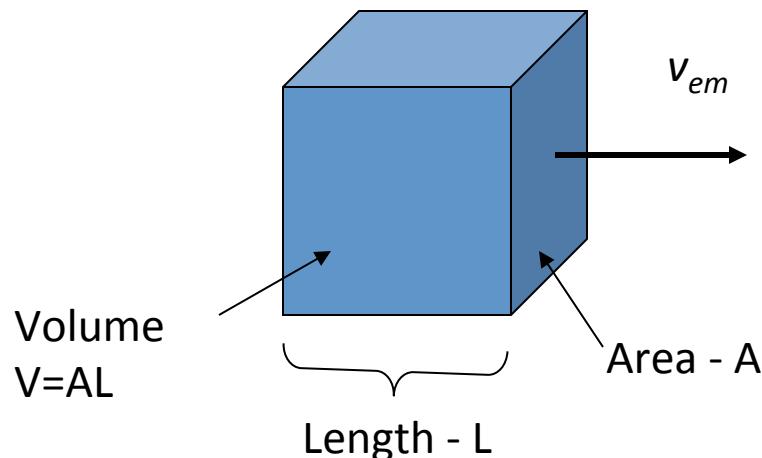
Energy density in electric and magnetic fields are equal for a wave in vacuum.

## Wave Intensity - Power/area

Energy density inside cube

$$u_E = \frac{\epsilon_0 |\vec{E}|^2}{2} = u_B = \frac{|\vec{B}|^2}{2\mu_0}$$

In time  $T=L/v_{em}$  an amount of energy



$$U = V(u_E + u_B) = AL\epsilon_0 |\vec{E}|^2$$

comes through the area A.

Intensity

$I = \text{Power}/\text{Area}$

$$I = \frac{U}{TA} = \sqrt{\frac{\epsilon_0}{\mu_0}} |\vec{E}|^2$$

## Poynting Vector

The power per unit area flowing in a given direction

$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$$

$$|\vec{S}| = I = \sqrt{\frac{\epsilon_0}{\mu_0}} |\vec{E}|^2$$

What are the units of

$I$  - W/m<sup>2</sup>,  $E$  - V/m

$$\sqrt{\frac{\mu_0}{\epsilon_0}}$$

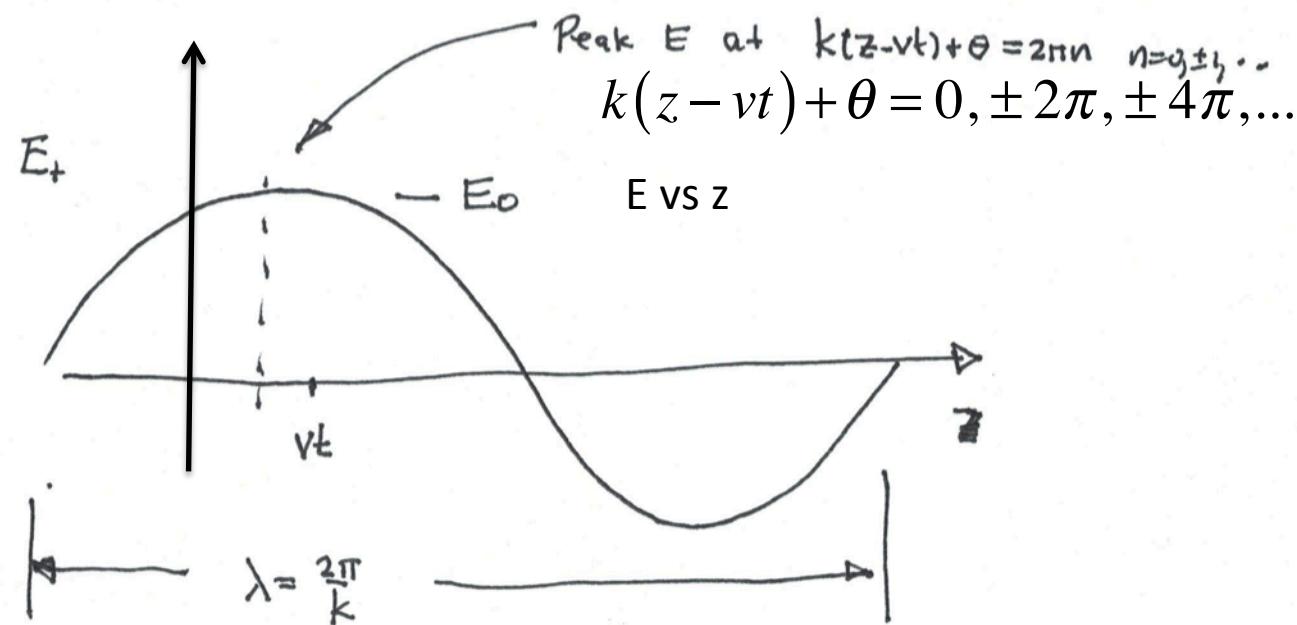
Ans: Ohms

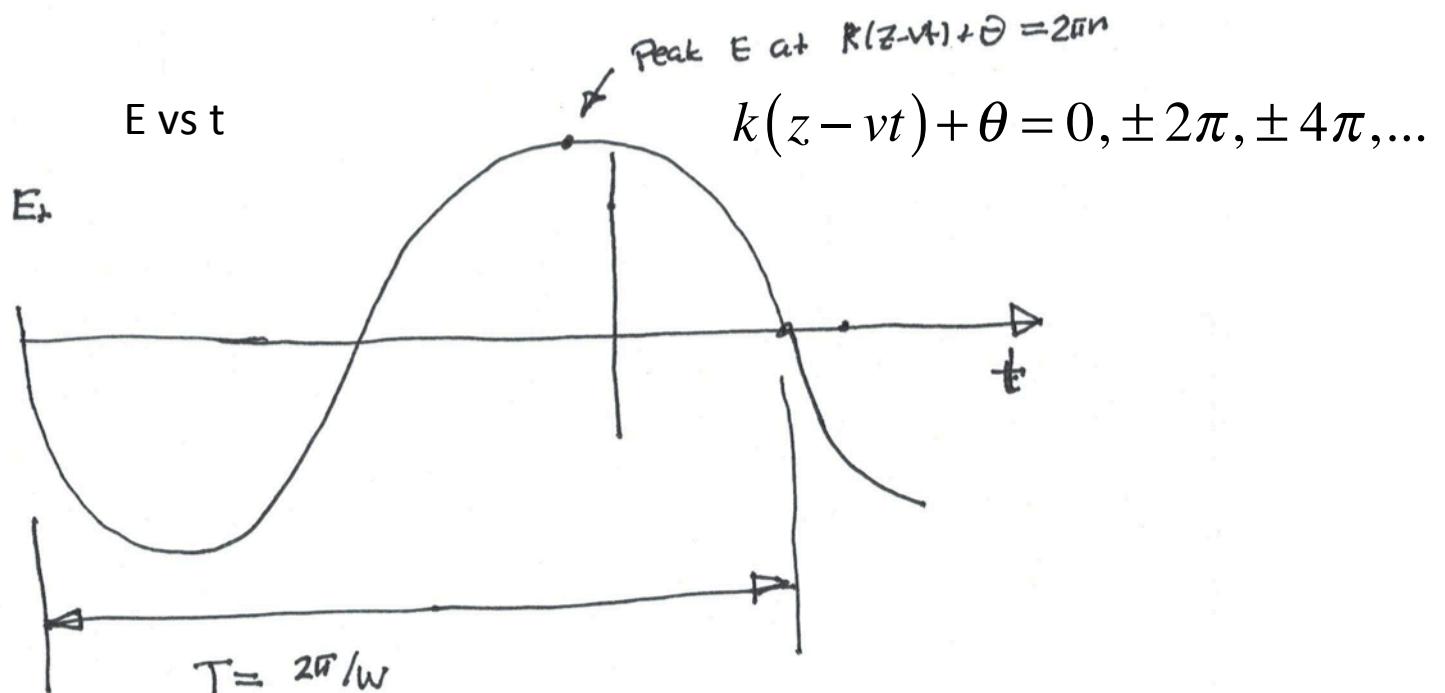
$$\sqrt{\frac{\mu_0}{\epsilon_0}} = 377\Omega$$

# Sinusoidal Waves

$$E_x = E_+(z - vt) + E_-(z + vt)$$

$$E_+(z - vt) = E_0 \cos[k(z - vt) + \theta]$$

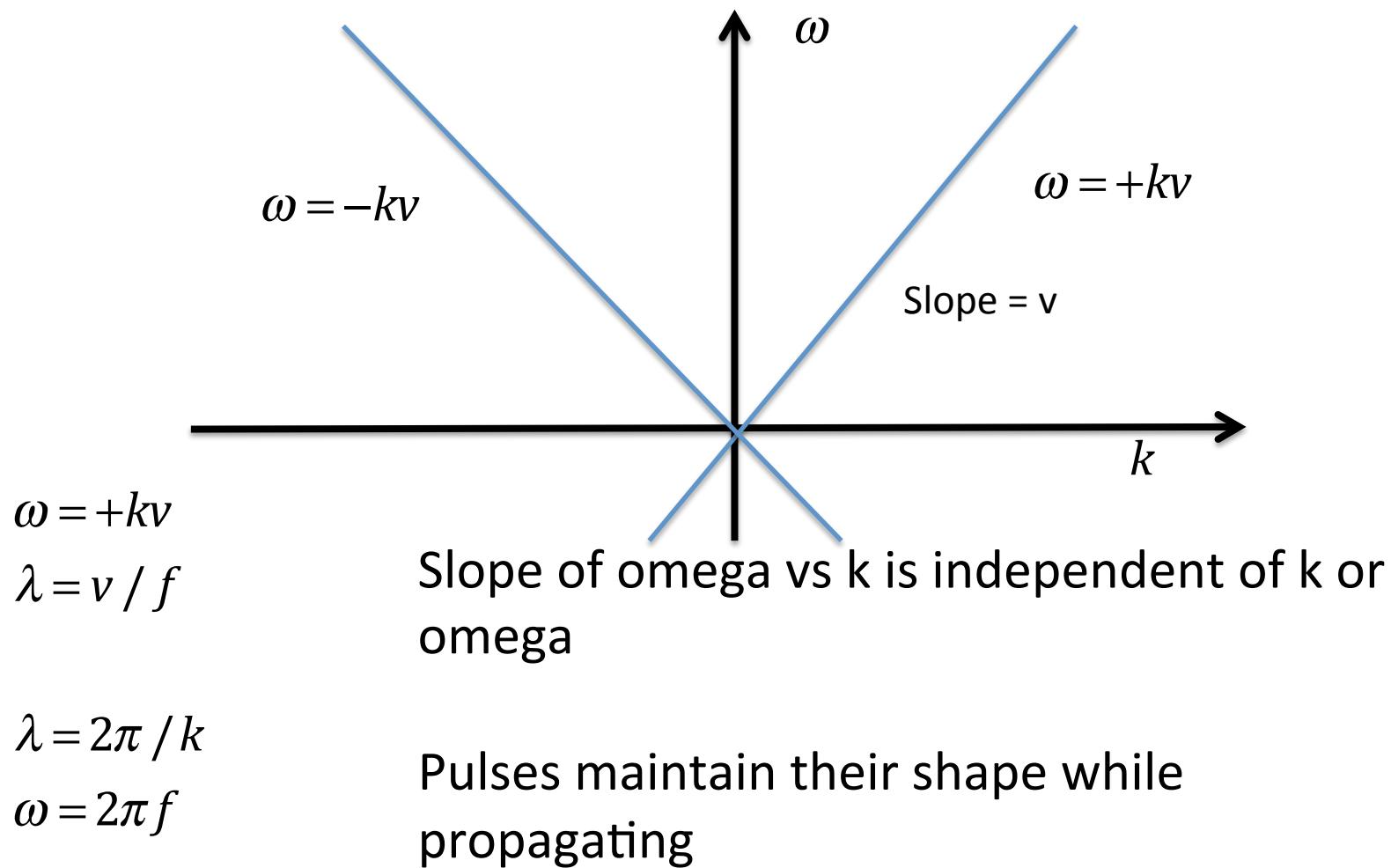




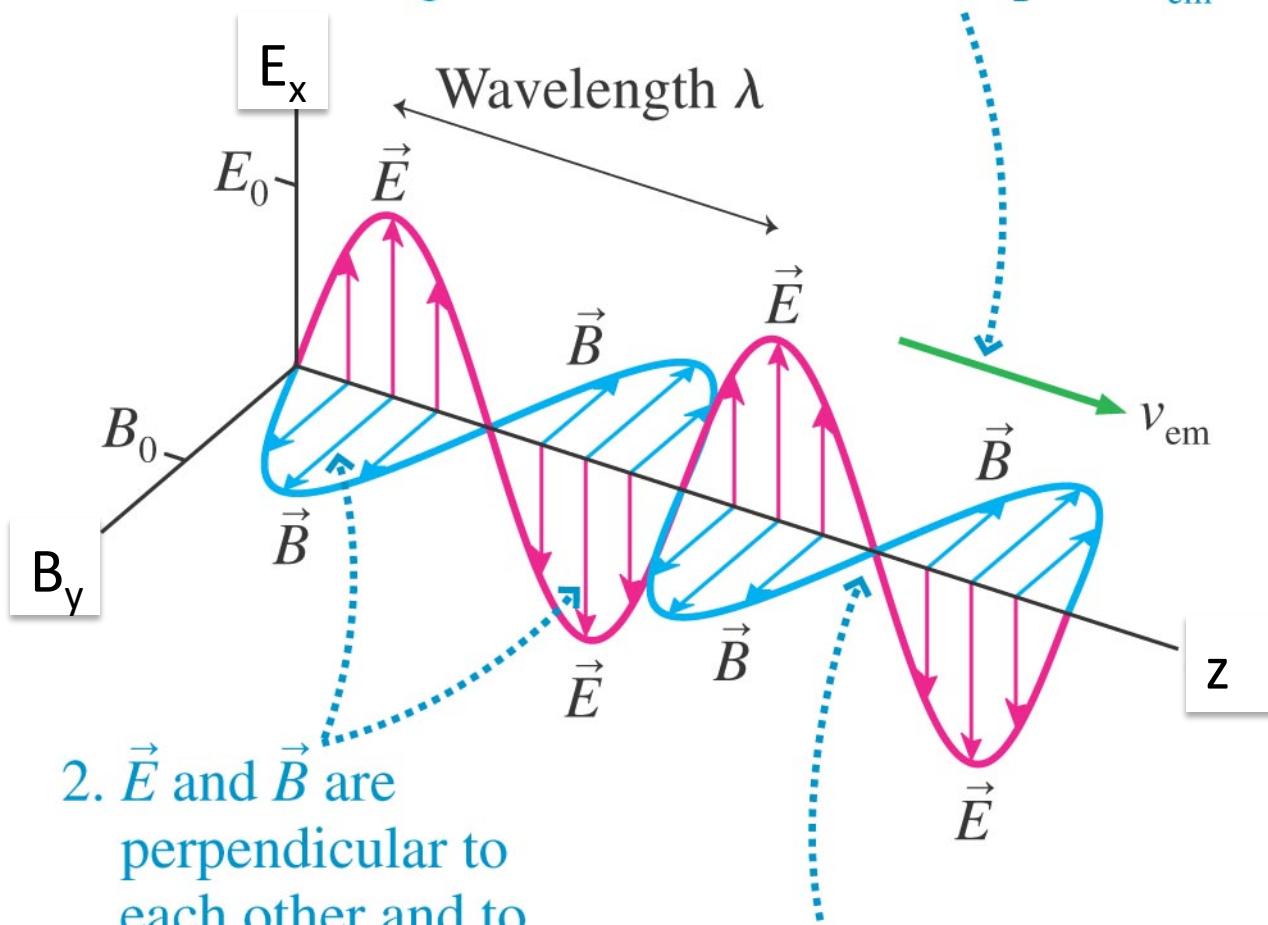
$$\omega = kv \quad \text{Dispersion relation given } k \text{ what is } \omega$$

# Dispersion Relation

Plane waves in a nondispersive medium



1. A sinusoidal wave with frequency  $f$  and wavelength  $\lambda$  travels with wave speed  $v_{\text{em}}$ .



2.  $\vec{E}$  and  $\vec{B}$  are perpendicular to each other and to the direction of travel. The fields have amplitudes  $E_0$  and  $B_0$ .

3.  $\vec{E}$  and  $\vec{B}$  are in phase. That is, they have matching crests, troughs, and zeros.

# Complex Phasors

## Sinusoidal Wave Parameters:

Properties of the medium

$$v = \frac{1}{\sqrt{\epsilon\mu}}, \quad \eta = \sqrt{\frac{\mu}{\epsilon}}$$

Determined by the source

$$\omega \text{ or } k \\ \text{usually } \omega$$

But  $\omega = kv$

Local amplitude and phase

$$E_0 \\ H_0 = E_0 / \eta \\ \theta$$

Three Equivalent representations

$$E_x = E_0 \cos[kz - \omega t + \theta] \text{ or } k$$

Physicists  
Engineers

$$E_x = \operatorname{Re} \left[ E_0 \exp[i(kz - \omega t + \theta)] \right] \quad j = -i = \sqrt{-1}$$

$$E_x = \operatorname{Re} \left[ E_0 \exp[j(\omega t - kz - \theta)] \right]$$

Griffiths uses the Physics convention

$$\begin{aligned}E_x &= \operatorname{Re} \left[ E_0 \exp \left[ i(kz - \omega t + \theta) \right] \right] \\&= \operatorname{Re} \left[ (E_0 e^{i\theta}) \exp \left[ i(kz - \omega t) \right] \right]\end{aligned}$$

$$\hat{E}_x = (E_0 e^{i\theta}) \text{ complex phasor amplitude}$$

# Quick Review of Phasors

Phasors - a way of representing amplitude and phase

Imaginary number

$$i = \sqrt{-1}$$

Engineers use  $j$   
Physicists and mathematicians  
use  $i$

Complex number

$$Z = X + iY$$

X is the real part  
Y is the imaginary part

Complex numbers follow the same rules of algebra as regular numbers

$$Z_1 = X_1 + iY_1 \quad Z_2 = X_2 + iY_2$$

Addition:

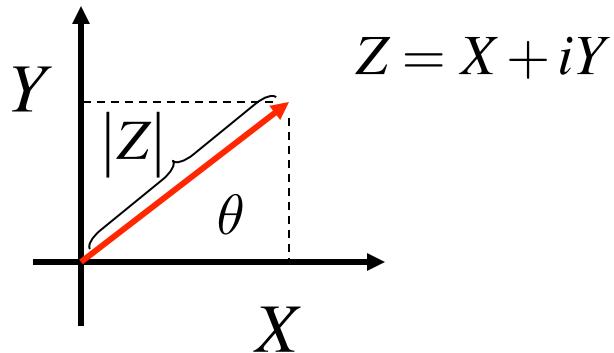
$$Z_1 + Z_2 = (X_1 + X_2) + i(Y_1 + Y_2)$$

Multiplication:

$$Z_1 Z_2 = (X_1 + iY_1)(X_2 + iY_2) = X_1 X_2 + i^2 Y_1 Y_2 + i(X_1 Y_2 + X_2 Y_1)$$

$-1$

**A complex number is specified by two real numbers**



Instead of real and imaginary parts can give magnitude and phase

$$|Z| = \sqrt{X^2 + Y^2}$$

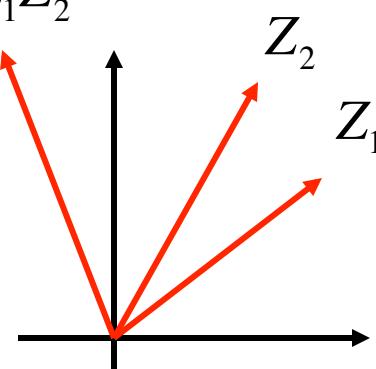
$$\tan \theta = Y / X$$

Multiplying complex numbers - part 2

Magnitudes multiply

Phases add

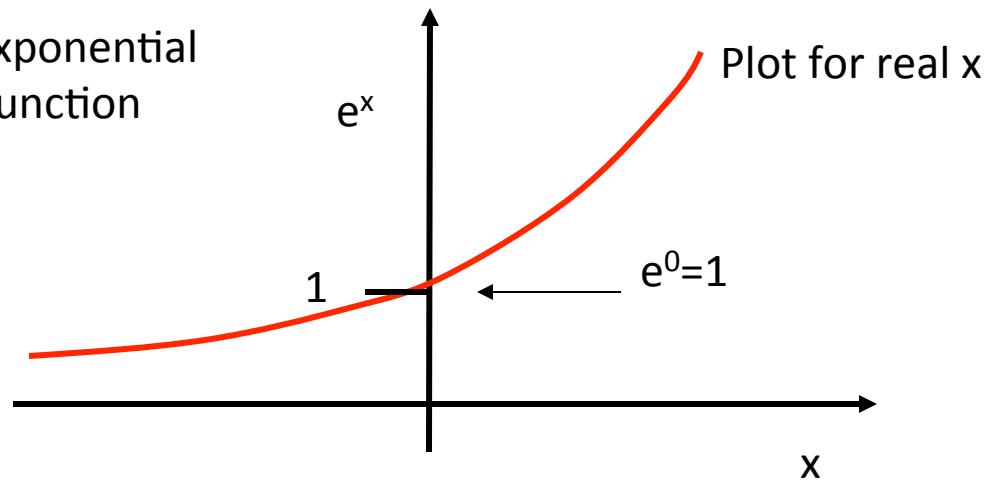
$$Z_3 = Z_1 Z_2$$



$$|Z_3| = |Z_1| |Z_2|$$

$$\theta_3 = \theta_1 + \theta_2$$

Exponential  
function



Plot for real  $x$

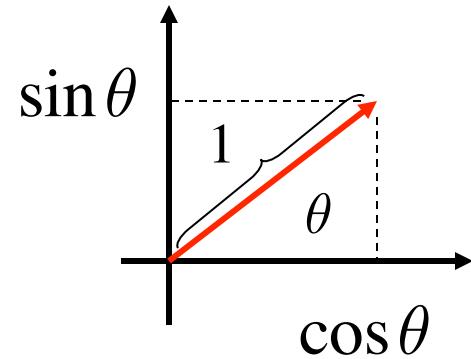
But, what if  $x$  is imaginary?

Let  $X = \cos \theta, Y = \sin \theta$   
 $Z = \cos \theta + i \sin \theta$

Then you can show:

$$\frac{dZ}{d\theta} = iZ$$

$$e^{i\theta} = \cos \theta + i \sin \theta$$



So:  $Z(\theta) = Z(0)e^{i\theta} = e^{i\theta}$

$$Z(0) = 1 + i0 = 1$$

## Phasors

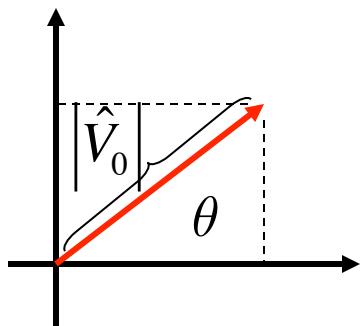
Suppose I have an oscillating voltage

$$\begin{aligned}V(t) &= V_0 \cos[\omega t - \theta] \\&= V_0 \cos[\theta - \omega t]\end{aligned}$$

I can write this as the real part of a complex number.

$$V(t) = \operatorname{Re} \left[ (V_0 e^{i\theta}) e^{-i\omega t} \right]$$

Call this  $\hat{V}_0$  a complex amplitude or “phasor”



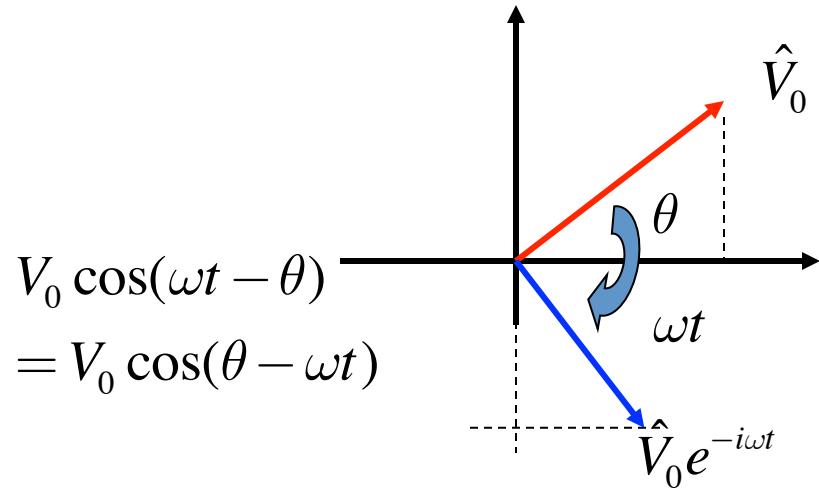
$$\hat{V}_0 = V_0 e^{i\theta}$$

$$|\hat{V}_0| = V_0$$

In this class, hat means a complex #

Magnitude of phasor gives peak amplitude of signal.  
Angle give phase of signal.

$$V(t) = \operatorname{Re}[\hat{V}_0 e^{-i\omega t}]$$



Multiplying  $\hat{V}_0$  by  $e^{-i\omega t}$  rotates the angle of the product by  $-\omega t$

Remember:  $|Z_3| = |Z_1||Z_2|$

$$\theta_3 = \theta_1 + \theta_2$$

$$\frac{d}{dt} I(t) = \frac{d}{dt} \operatorname{Re}[\hat{I}_0 e^{-i\omega t}] = \operatorname{Re}\left[\frac{d}{dt} \hat{I}_0 e^{-i\omega t}\right] = \operatorname{Re}[-i\omega \hat{I}_0 e^{-i\omega t}]$$

Now suppose

$$V(t) = L \frac{d}{dt} I(t)$$

$$\hat{V}_0 = -i\omega L \hat{I}_0$$

$$\operatorname{Re}[\hat{V}_0 e^{-i\omega t}] = \operatorname{Re}[-i\omega L \hat{I}_0 e^{-i\omega t}]$$

$$\frac{d}{dt} \Rightarrow -i\omega$$

Real parts equal

$$\hat{V}_0 e^{-i\omega t} = -i\omega L \hat{I}_0 e^{-i\omega t}$$