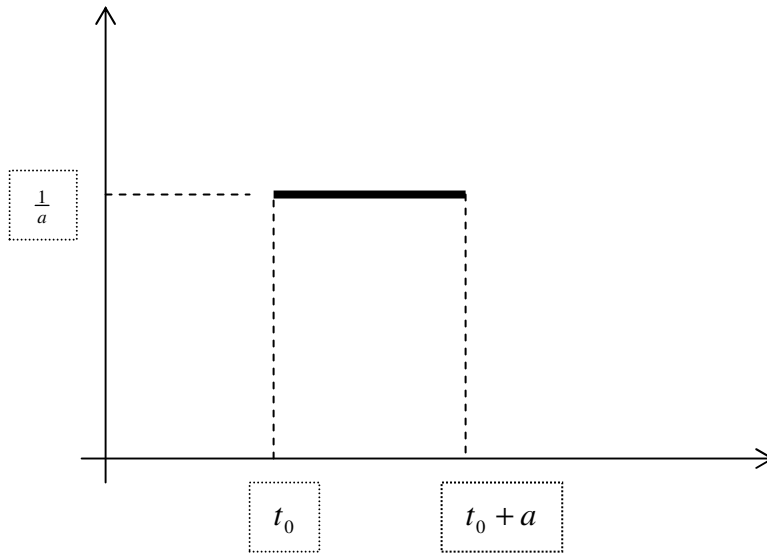


The Unit Impulse Function¹

Consider square wave of width a and height $1/a$, anchored at t_0 , i.e.:

$$\Delta_a(t_0) = \frac{1}{a}(u(t-t_0) - u(t-(t_0+a))) = \frac{u(t-t_0) - u(t-t_0-a)}{a}$$



Then by definition: $\delta(x-x_0) = \lim_{a \rightarrow 0} \Delta_a(x_0) = \lim_{a \rightarrow 0} \frac{(u(x-x_0) - u(x-x_0-a))}{a}$

Lemma 1: $\frac{d}{dt} u(t-t_0) = \delta(t-t_0)$

Proof: $\frac{d}{dt} u(t-t_0) = \lim_{h \rightarrow 0} \frac{u(t-t_0+h) - u(t-t_0)}{h}$

Define: $h = -a$. Then:

$$\frac{d}{dt} u(t-t_0) = \lim_{a \rightarrow 0} \frac{u(t-t_0-a) - u(t-t_0)}{-a} = \lim_{a \rightarrow 0} \frac{u(t-t_0) - u(t-t_0-a)}{a} = \delta(t-t_0)$$

¹ Also known as the 'Dirac delta-function,' in homage to physicist Paul Dirac.

Lemma 2: $L\{\delta(t - t_0)\} = e^{-st_0}$

Proof:

$$\begin{aligned} L\{\delta(t - t_0)\} &= L\left\{\lim_{a \rightarrow 0} \frac{u(t - t_0) - u(t - t_0 - a)}{a}\right\} = \lim_{a \rightarrow 0} \frac{1}{a} \{L\{u(t - t_0)\} - L\{u(t - t_0 - a)\}\} \\ &= \lim_{a \rightarrow 0} \frac{1}{a} \left\{ \frac{e^{-t_0 s}}{s} - L\{u(t - (t_0 + a))\} \right\} = \lim_{a \rightarrow 0} \frac{1}{a} \left\{ \frac{e^{-t_0 s}}{s} - \frac{e^{-s(t_0 + a)}}{s} \right\} = e^{-t_0 s} \lim_{a \rightarrow 0} \frac{(1 - e^{-as})}{as} \\ &= e^{-t_0 s} \lim_{a \rightarrow 0} \frac{\frac{d}{da}(1 - e^{-as})}{\frac{d}{da}(as)} = e^{-t_0 s} \lim_{a \rightarrow 0} \frac{se^{-as}}{s} = e^{-t_0 s} \lim_{a \rightarrow 0} e^{-as} = e^{-t_0 s} \end{aligned}$$

(Note how L'Hopital's Rule must be invoked to evaluate this above limit)

So the LT for the unit impulse function is a decaying exponential. Hence the inverse LT of a decaying exponential is the unit impulse function. Note: We can also verify the result in Lemma2 using THM4:

Since: $\frac{d}{dt}u(t - t_0) = \delta(t - t_0)$ then: $u(t - t_0) = \int_{t_0}^t \delta(\omega - t_0) d\omega$

Hence according to Thm4: $L\left\{\int_{t_0}^t \delta(\omega - t_0) d\omega\right\} = \frac{1}{s} L\{\delta(t - t_0)\}$

But we also know that: $L\left\{\int_{t_0}^t \delta(\omega - t_0) d\omega\right\} = L\{u(t - t_0)\} = \frac{e^{-t_0 s}}{s}$

Hence: $L\left\{\int_{t_0}^t \delta(\omega - t_0) d\omega\right\} = L\{u(t - t_0)\} = \frac{e^{-t_0 s}}{s} = \frac{1}{s} L\{\delta(t - t_0)\} \Rightarrow L\{\delta(t - t_0)\} = e^{-t_0 s}$

In particular, if $t_0 = 0$, then: $L\{\delta(t - 0)\} = L\{\delta(t)\} = e^{-s \cdot 0} = 1$

Or: $L^{-1}\{k\} = kL^{-1}\{1\} = k\delta(t)$

(I.e., the inverse LT of any constant function k (a horizontal line in s -space) is the unit impulse function with amplitude k , i.e. an infinite vertical ray in t -space!)

Example: Use the above result to find²: $L^{-1}\left[\frac{s^2}{s^2 - a^2}\right]$

² This was on part of the original formulation of Problem IIb, Assignment I. But I decided to make it simpler and remove the squared term.

Solution: According to Thm2:

$$F(s) = \frac{s^2}{s^2 - a^2} = s \frac{s}{s^2 - a^2} = sL\{\cosh at\} \Rightarrow sL\{\cosh at\} = L\left\{\frac{d}{dt} \cosh at\right\} + \cosh(a \cdot 0)$$

Hence:
$$F(s) = \frac{s^2}{s^2 - a^2} = sL\{\cosh at\} = L\{a \sinh at\} + 1$$

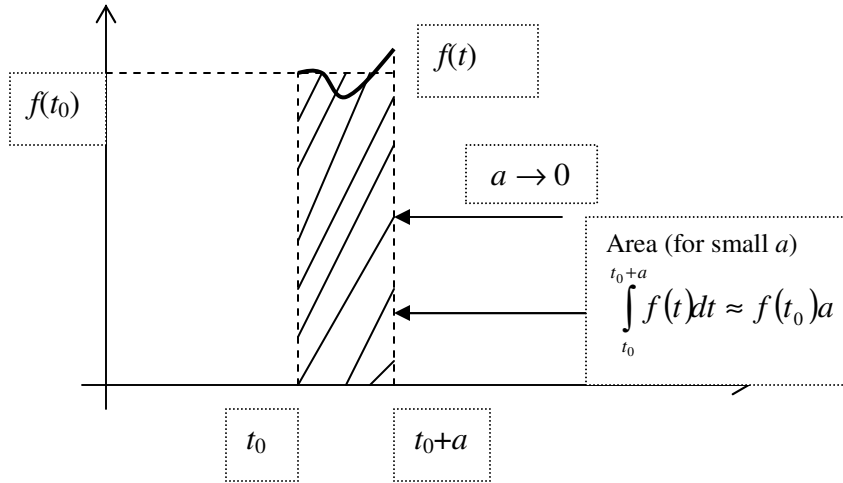
So:
$$L^{-1}\left[\frac{s^2}{s^2 - a^2}\right] = L^{-1}\{L\{a \sinh at\} + 1\} = a \sinh at + L^{-1}\{1\} = a \sinh at + \delta(t)$$

Lemma 3:
$$\int_{-\infty}^{\infty} f(t)\delta(t - t_0) = f(t_0)$$

Proof:

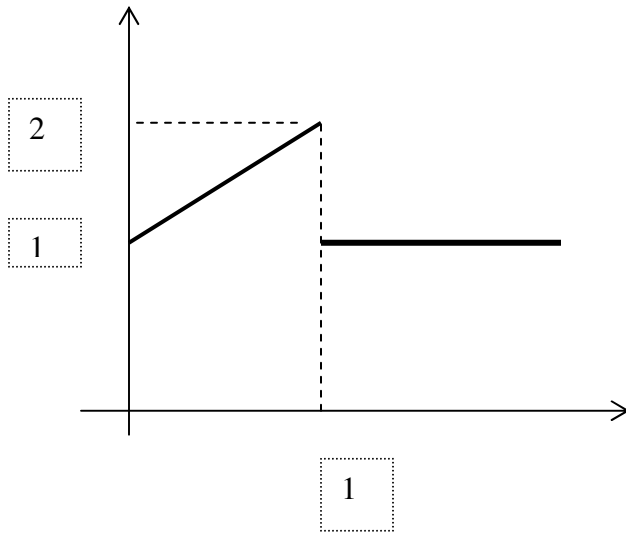
$$\begin{aligned} \int_{-\infty}^{\infty} f(t)\delta(t - t_0) &= \int_{-\infty}^{\infty} f(t) \lim_{a \rightarrow 0} \left[\frac{u(t - t_0) - u(t - t_0 - a)}{a} \right] \\ &= \lim_{a \rightarrow 0} \frac{1}{a} \left\{ \int_{-\infty}^{\infty} f(t)u(t - t_0)dt - \int_{-\infty}^{\infty} f(t)u(t - t_0 - a)dt \right\} = \lim_{a \rightarrow 0} \frac{1}{a} \left\{ \int_{t_0}^{\infty} f(t)dt - \int_{t_0+a}^{\infty} f(t)dt \right\} \\ &= \lim_{a \rightarrow 0} \frac{1}{a} \int_{t_0}^{t_0+a} f(t)dt = \lim_{a \rightarrow 0} \frac{1}{a} [f(t_0)a] = f(t_0) \end{aligned}$$

The last step in the proof involved approximating the integral $\int_{t_0}^{t_0+a} f(t)dt$ with the area of a rectangle of height $f(t_0)$ and width a . Note that in the limit $a \rightarrow 0$, the approximation becomes exact. (See figure next page) :



• **Example (Modified 2d) Shen, p53)**

Find $f(t)$, $f'(t)$ and the associated LTs of $f(t)$, $f'(t)$:



a.)

$$f(t) = (t+1)[u(t) - u(t-1)] + 1 \cdot u(t-1) = tu(t) + u(t) - tu(t-1) - u(t-1) + u(t-1)$$

$$= (t+1)u(t) - tu(t-1)$$

b.)

$$\begin{aligned} f'(t) &= (t+1)\frac{d}{dt}u(t) + u(t) - t\frac{d}{dt}u(t-1) - u(t-1) \\ &= (t+1)\delta(t) + u(t) - t\delta(t-1) - u(t-1) = t\delta(t) + \delta(t) + u(t) - t\delta(t-1) - u(t-1) \\ &= \delta(t) - t\delta(t-1) + u(t) - u(t-1) \end{aligned}$$

(Recall remark, p51, Sheng: $t\delta(t) = 0$)

c.)

$$\begin{aligned} L\{f(t)\} &= L\{(t+1)u(t) - tu(t-1)\} = L\{tu(t)\} + L\{u(t)\} - L\{tu(t-1)\} \\ &= (-1)\frac{d}{ds}L\{u(t)\} + L\{u(t)\} - (-1)\frac{d}{ds}L\{u(t-1)\} \\ &= -\frac{d}{ds}\left(\frac{1}{s}\right) + \frac{1}{s} + \frac{d}{ds}\left(\frac{e^{-s}}{s}\right) = \frac{1}{s^2} + \frac{1}{s} + \frac{-se^{-s} - e^{-s}}{s^2} = \frac{1}{s^2}(1 + s - se^{-s} - e^{-s}) \end{aligned}$$

d.)

$$\begin{aligned} L\{f'(t)\} &= L\{\delta(t) - t\delta(t-1) + u(t) - u(t-1)\} = L\{\delta(t)\} - L\{t\delta(t-1)\} + L\{u(t)\} - L\{u(t-1)\} \\ &= 1 - (-1)\frac{d}{ds}L\{\delta(t-1)\} + L\{u(t)\} - L\{u(t-1)\} = 1 + \frac{d}{ds}(e^{-s}) + \frac{1}{s} - \frac{e^{-s}}{s} = 1 - e^{-s} + \frac{1}{s}(1 - e^{-s}) \\ &= \frac{1}{s}(s - se^{-s} + 1 - e^{-s}) = \frac{1}{s}(1 + s - se^{-s} - e^{-s}) \end{aligned}$$

Note the obvious connection between the answer in d) and c). This can be substantiated using THM2:

$$L\{f'(t)\} = sL\{f(t)\} - f(0)$$

Here: $f(0) = (0+1)u(0) - 0u(0-1) = 0$ (since $u(t) = 0$ for all $t \leq 0$)

And as results c) and d) clearly show:

$$L\{f'(t)\} = sL\{f(t)\}$$

$$\begin{aligned} \text{Since: } L\{f'(t)\} &= \frac{1}{s}(1 + s - se^{-s} - e^{-s}) \\ L\{f(t)\} &= \frac{1}{s^2}(1 + s - se^{-s} - e^{-s}) \end{aligned}$$