

## H-theorem

Recall that the rate of entropy density production in the distribution for species  $j$  is

$$\frac{\partial S_j}{\partial t} = - \int d^3v \ln f_j \sum_k C_{jk} \quad (1)$$

Let  $\dot{S}_{jj}$  denote the contribution to this entropy production from the *nonlinear* Fokker-Planck operator for self-collisions:

$$\dot{S}_{jj} = - \int d^3v \ln f_j C_{jj} \quad (2)$$

(It would be silly to use a *linearized* operator before we knew we had a Maxwellian.)

**Claim:**  $\dot{S}_{jj}$  vanishes only if  $f_j$  is a drifting Maxwellian.

### Proof:

For the rest of this proof I'll drop the subscript on  $f_j$ . Plugging in the Landau form of the collision operator,

$$\begin{aligned} \dot{S}_{jj} &= \gamma \int d^3v \ln f \nabla_v \cdot \int d^3v' \nabla_g \nabla_g g \cdot [f' \nabla_v f - f \nabla_{v'} f'] \\ &= -\gamma \int d^3v \int d^3v' f' f [\nabla_v \ln f] \cdot \nabla_g \nabla_g g \cdot [\nabla_v \ln f - \nabla_{v'} \ln f'] \\ &= -\gamma \int d^3v \int d^3v' f' f [\nabla_{v'} \ln f'] \cdot \nabla_g \nabla_g g \cdot [\nabla_{v'} \ln f' - \nabla_v \ln f] \end{aligned} \quad (3)$$

where in the last line we have switched primed and unprimed variables. Averaging the last two lines, we get

$$\dot{S}_{jj} = -\frac{\gamma}{2} \int d^3v \int d^3v' f' f [\nabla_v \ln f - \nabla_{v'} \ln f'] \cdot \nabla_g \nabla_g g \cdot [\nabla_v \ln f - \nabla_{v'} \ln f'] \quad (4)$$

Letting  $\mathbf{H} = \nabla_v \ln f - \nabla_{v'} \ln f'$ , then

$$\dot{S}_{jj} = -\frac{\gamma}{2} \int d^3v \int d^3v' \frac{f' f}{g^3} \left\{ g^2 H^2 - [\mathbf{g} \cdot \mathbf{H}]^2 \right\} \quad (5)$$

If  $\dot{S}_{jj}$  vanishes, then the quantity in curly braces must vanish for all  $\mathbf{v}$  and for all  $\mathbf{v}'$ . Using the Schwartz inequality for  $\mathbb{R}^3$ , it must then be that

$$\mathbf{g} \parallel \mathbf{H} \text{ for all } \mathbf{v} \text{ and for all } \mathbf{v}'. \quad (6)$$

At this point it is helpful to cast the problem into the notation of Sussman and Wisdom. As part of this notation, throughout this document I will use parentheses only to denote the input parameters to a function, not to denote the order of operations; for the latter I will use only square and curly braces. The statement (6) means that there exists a function  $\lambda : \mathbb{R}^6 \rightarrow \mathbb{R}$  such that the following three equations are true for all  $a, b, c, d, e$ , and  $h$ :

$$[a - d] \lambda(a, b, c, d, e, h) = F_0(a, b, c) - F_0(d, e, h), \quad (7)$$

$$[b - e] \lambda(a, b, c, d, e, h) = F_1(a, b, c) - F_1(d, e, h), \quad (8)$$

$$[c - f] \lambda(a, b, c, d, e, h) = F_2(a, b, c) - F_2(d, e, h) \quad (9)$$

where  $F_i = \partial_i \ln f$ . (Here I'm just using  $a$ ,  $b$ , and  $c$  to represent the components of  $\mathbf{v}$ , and using  $d$ ,  $e$ , and  $h$  to represent the components of  $\mathbf{v}'$ .) Considering (7) for the particular case that  $a = d$ , then

$$F_0(a, b, c) = F_0(a, e, h). \quad (10)$$

Therefore  $F_0$  must be independent of its latter two inputs, so we can write

$$F_0(a, b, c) = x(a) \quad (11)$$

where  $x$  is some  $\mathbb{R} \rightarrow \mathbb{R}$  function. Repeating this procedure for (8) and (9), we find there must exist functions  $y$  and  $z$  such that

$$F_1(a, b, c) = y(b) \quad \text{and} \quad (12)$$

$$F_2(a, b, c) = z(c). \quad (13)$$

Using these results we can rewrite (7)-(9) as

$$[a - d] \lambda(a, b, c, d, e, h) = x(a) - x(d), \quad (14)$$

$$[b - e] \lambda(a, b, c, d, e, h) = y(b) - y(e), \quad (15)$$

$$[c - h] \lambda(a, b, c, d, e, h) = z(c) - z(h). \quad (16)$$

Note that if  $a \neq d$ , then (14) can be written such that the right-hand side is independent of  $b, c, e$ , and  $h$ :

$$\lambda(a, b, c, d, e, h) = [x(a) - x(d)] / [a - d] \quad (17)$$

Thus,  $\lambda$  must be independent of its 2<sup>nd</sup>, 3<sup>rd</sup>, 5<sup>th</sup>, and 6<sup>th</sup> inputs, at least outside of the hyperplane in its input space defined by  $a = d$ . If we require that  $f$  be smooth, then since  $x(a) = [\partial_0 \ln f](a, b, c)$  is a derivative of the distribution function,  $x$  will be differentiable. In this case, in the limit  $d \rightarrow a$ , the right-hand side of (17) converges to  $[Dx](a)$ , and so we can write

$$\lambda(a, b, c, a, e, h) = [Dx](a). \quad (18)$$

The right-hand side is independent of  $b, c, e$ , and  $h$ , and so  $\lambda$  must be independent of these inputs on the hyperplane  $a = d$  as well.

The argument above can be repeated using (15) to show that  $\lambda$  is also independent of its 1<sup>st</sup> and 4<sup>th</sup> inputs. Therefore  $\lambda$  is a constant.

We can now write (14) as

$$x(a) - a\lambda = x(d) - d\lambda. \quad (19)$$

Since this equation is true for all  $a$  and  $d$ , each side of this equation is a constant (i.e. independent of  $a, b, c, d, e$ , and  $h$ .) We name this constant  $x_0$ . Recalling the definition of  $x$  in terms of the distribution function, then

$$[\partial_0 \ln f](a, b, c) = a\lambda + x_0. \quad (20)$$

Integrating in  $a$ , then

$$\ln f(a, b, c) = \frac{\lambda}{2} a^2 + x_0 a + A(b, c) \quad (21)$$

where  $A(b, c)$  is the “constant of integration.” Repeating the last few steps using (15) and (16) instead of (14), we find that there must exist constants  $y_0$  and  $z_0$  and functions  $B(a, c)$  and  $C(a, b)$  such that

$$\ln f(a, b, c) = \frac{\lambda}{2} b^2 + y_0 b + B(a, c) \quad (22)$$

and

$$\ln f(a, b, c) = \frac{\lambda}{2} c^2 + z_0 c + C(a, b). \quad (23)$$

The only way that (21), (22), and (23) can simultaneously be true is if

$$\ln f(a, b, c) = \frac{\lambda}{2} [a^2 + b^2 + c^2] + x_0 a + y_0 b + z_0 c + w \quad (24)$$

where  $w$  is another constant. Now recognizing  $\lambda \rightarrow -M / T$ ,  $x_0 = M V_x / T$ , etc., we see that  $f$  is a Maxwellian.