

Subsonic
Fluid response : Reduced eqns for $\partial_t \ll \sqrt{g} \rightarrow 0$ ^{or}
for $g \neq 0$

$$\partial_t n + \vec{\nabla} \cdot (n \vec{u}) = 0 \quad (1)$$

$$nM d\vec{u}/dt = -\vec{\nabla} p + nM \vec{g} \quad (2)$$

$$dp/dt + \gamma p \vec{\nabla} \cdot \vec{u} = 0 \quad (3)$$

Expect $\partial_t \ll \sqrt{g/L} \ll kc_s$

Expect convection $\frac{u}{L} \ll \partial_t$

Size of terms :

$$(1) \quad \partial_t n : \frac{nu}{L} \Rightarrow 1:1$$

$$(3) \quad \partial_t : u/L \Rightarrow 1:1$$

$$(2) \quad \cancel{nM} \partial_t : \frac{p}{L n M} : nM g$$

$$\ll \frac{u}{L} \partial_t : \frac{c_s^2}{L^2} : \frac{g}{L}$$

$$\ll \partial_t^2 : (c_s/L)^2 : (g/L)$$

$$\epsilon^2 : 1 : \epsilon^2$$

lowest order

$$\vec{\nabla} p = 0 \Rightarrow p = p(t) \rightarrow p_0$$

$$\Rightarrow \boxed{\begin{aligned} \partial_t n + \vec{\nabla} \cdot (n \vec{u}) &= 0 & (1') \\ \vec{\nabla} \cdot \vec{u} &= 0 & (3') \end{aligned}} \leftarrow \text{incompressible}$$

$$nM \frac{d\vec{u}}{dt} = -\vec{\nabla} p + nM \vec{g}$$

annihilate

$$\Rightarrow \vec{\nabla} \times (nM \frac{d\vec{u}}{dt}) = \vec{\nabla} n \times M \vec{g}$$

$$\Rightarrow \boxed{\vec{\nabla} \times (n \frac{d\vec{u}}{dt}) = \vec{\nabla} n \times \vec{g}} \quad (2')$$

(1'), (2'), (3') - New set for $\{n, \vec{u}\}$ _{4 vars}

Reduced eqns for low frequency.

p is eliminated.

There are sound waves here.

Linearized Red^d Eqs and RT mode ^{R3}

Equil $\vec{\nabla} n \times \vec{g} = 0 \Rightarrow n(x)$

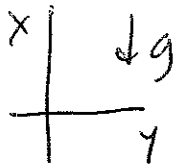
$n(x)$ arbitrary

Perturb

$$\partial_t \tilde{n} + \vec{\nabla} \cdot (n \tilde{\vec{u}}) = 0$$

$$\vec{\nabla} \cdot \tilde{\vec{u}} = 0$$

$$\vec{\nabla} \times (n \partial_t \tilde{\vec{u}}) = \vec{\nabla} \tilde{n} \times \vec{g}$$



$\tilde{n} \rightarrow \tilde{n}(x) e^{iky - i\omega t}$; $k_z = 0$
is ok

Note $\hat{y} \cdot \vec{\nabla} \times n \tilde{\vec{u}} = -\partial_x (n \tilde{u}_z) = 0$

$$\Rightarrow \tilde{u}_z = 0$$

only $\{\tilde{n}, \tilde{u}_x, \tilde{u}_y\}$

Now $\vec{\nabla} \cdot \tilde{\vec{u}} = 0 \Rightarrow \tilde{\vec{u}} = \vec{z} \times \vec{\nabla} \tilde{\varphi}$

$\tilde{\varphi}$ = stream function

$$\Rightarrow \vec{\nabla} \times (n \vec{z} \times \vec{\nabla} \tilde{\varphi}) = \vec{z} \cdot \vec{\nabla} (n \vec{\nabla} \tilde{\varphi})$$



$$\Rightarrow \begin{cases} \partial_t \tilde{\eta} + \hat{z} \times \vec{\nabla} \tilde{\varphi} \cdot \vec{\nabla} \eta = 0 \\ \partial_t \vec{\nabla} \cdot (n \vec{\nabla} \tilde{\varphi}) = g \partial_y \tilde{\eta} \end{cases}$$

24

$$\Rightarrow -i\omega \tilde{\eta} - ik n' \tilde{\varphi} = 0$$

$$-i\omega \vec{\nabla} \cdot (n \vec{\nabla} \tilde{\varphi}) = +ik g \tilde{\eta}$$

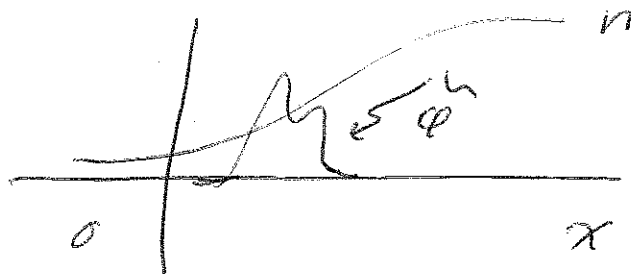
Eliminate $\tilde{\eta} \Rightarrow$

$$\vec{\nabla} \cdot [n \vec{\nabla} \tilde{\varphi}] = \frac{gk^2}{\omega^2} \frac{n'}{n} \tilde{\varphi}$$

eigenvalue eqn

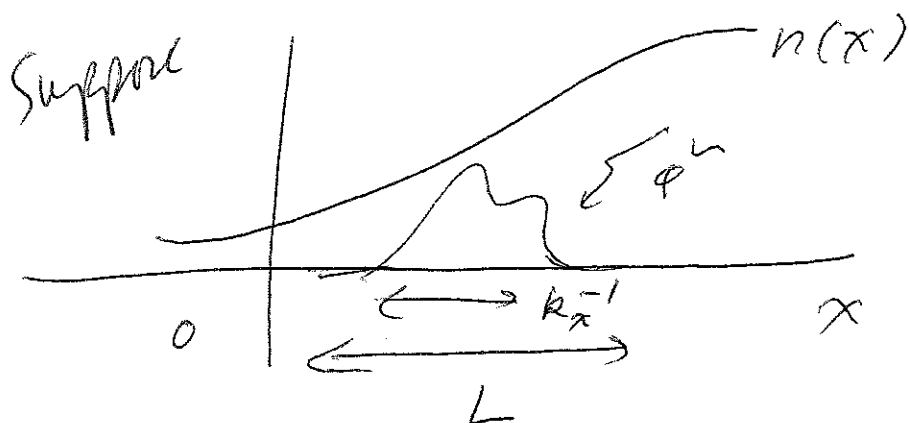
$$\Rightarrow \tilde{\varphi}'' - k^2 \tilde{\varphi} = \left(\frac{gk^2}{\omega^2} \right) \frac{n'}{n} \tilde{\varphi}$$

$$\tilde{\varphi}(|x| \rightarrow \infty) \rightarrow 0$$



Diffuse Profile Analysis (RT)

RS



let $n = n(x)$

Note: $k_y \rightarrow \infty \Rightarrow \infty \Rightarrow \omega^2 = -\frac{g k^2}{\omega^2 n} \frac{n'}{n}$

$\Rightarrow \omega^2 \approx -\frac{g n'}{n} (x)$ can't be

but expect max instability where

$\frac{n'}{n} (x) \rightarrow \max \Rightarrow \left(\frac{n'}{n}\right) \approx \left(\frac{n'}{n}\right)_{\max} - \left(\frac{n'}{n}\right)_{\max} \frac{(x-x_m)^2}{2}$

$$\frac{\omega''}{k^2} \omega \approx \frac{g \left(\frac{n'}{n}\right)_m}{\omega^2 (n)_m} \left[1 - \frac{s^2}{L_n^2}\right] \omega \quad s \equiv x - x_m$$

self cons $\Leftrightarrow k_x L_n \gg 1$

$$\Rightarrow \omega_{ss} = k^2 \left\{ 1 + \left(\frac{g n'}{\omega n}\right)_m \left[1 - \frac{s^2}{L_n^2}\right] \right\} \omega$$

$\omega'(|x| \rightarrow \infty) \rightarrow 0$

Parabolic Cylinder Egn

eigenvalue problem

⇒ H.O. in QM, Hermite Fns

Try lowest: $\psi \rightarrow \psi e^{-s^2/2\Delta^2}$

$$s \equiv (x - x_m)$$

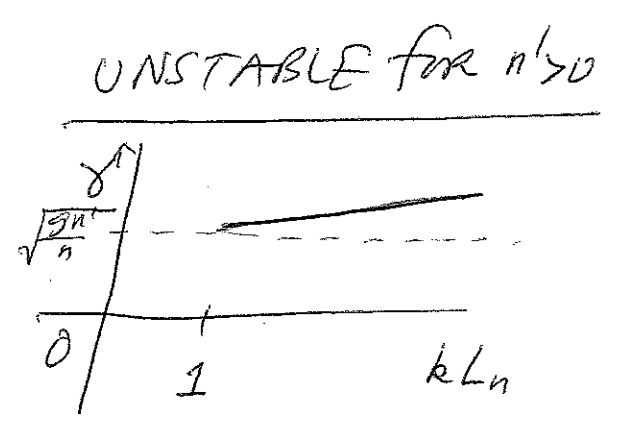
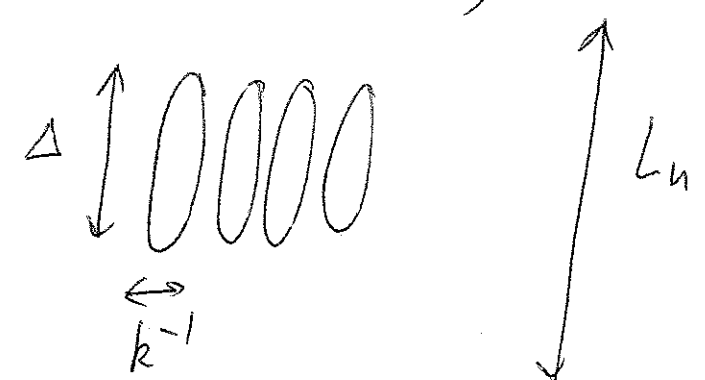
$$\psi_s = -\frac{s\psi}{\Delta^2}, \quad \psi_{ss} = -\frac{\psi}{\Delta^2} + \frac{s^2\psi}{\Delta^4}$$

$$\Rightarrow \boxed{\omega^2 = \frac{-g(n'/n)_{max}}{(1 + 1/k^2\Delta^2)}} \quad \text{and} \quad \frac{k^2 L_n^2}{k^4 \Delta^4} = 1 + \frac{1}{k^2 \Delta^2}$$

But demand $\Delta \ll L_n$

$$\Rightarrow k\Delta \approx \sqrt{kL_n} \Rightarrow \Delta \approx \sqrt{L_n/k}$$

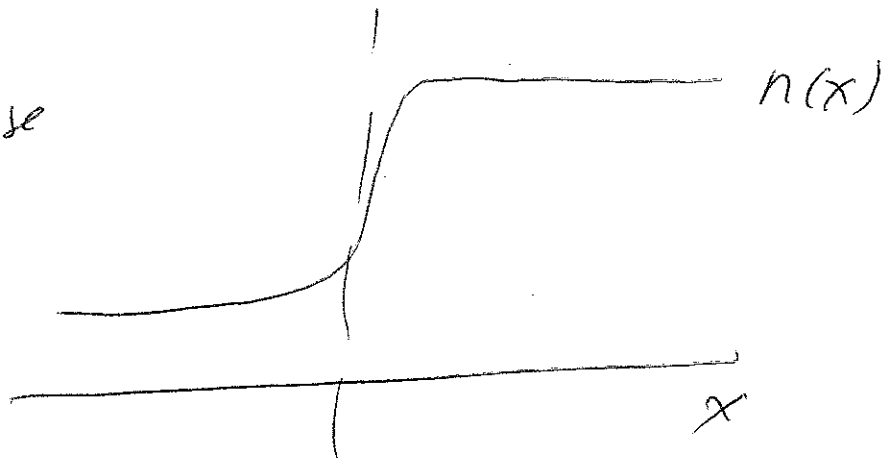
with $kL_n \gg 1, k\Delta \gg 1$



Sharp Boundary Analysis (RT)

RT

Suppose



i.e. suppose $L_n \rightarrow 0$ (sharp change)

\Rightarrow let $n(x) \rightarrow$ step function @ $x=0$

$kL_n \ll 1$ and $L_n \ll \Delta$

$$\boxed{\psi'' - k^2 \psi = \frac{g h^2}{\omega^2} \frac{n'}{n} \psi}$$

$k \neq 0 \Rightarrow n' = 0 \Rightarrow \psi'' - k^2 \psi = 0$

$$\Rightarrow \psi = \begin{cases} \hat{\psi}_+ e^{-kx} & , \quad 0 < x \\ \hat{\psi}_- e^{\pm kx} & , \quad x < 0 \end{cases}$$

Now $\frac{d}{dx} \left[n \frac{d\psi}{dx} \right] = \frac{g h^2}{\omega^2} \frac{dn}{dx} \psi$

$$\int_{-E}^E dx \Rightarrow \left[n \frac{d\psi}{dx} \right]_{-E}^E = \frac{g h^2}{\omega^2} \left[[n\psi]_{-E}^E - \int_{-E}^E dx n \psi' \right]$$

$$\Rightarrow [n\tilde{\phi}]_+^+ = \frac{g\hbar^2}{\omega^2} [n\tilde{\phi}]_+^+$$

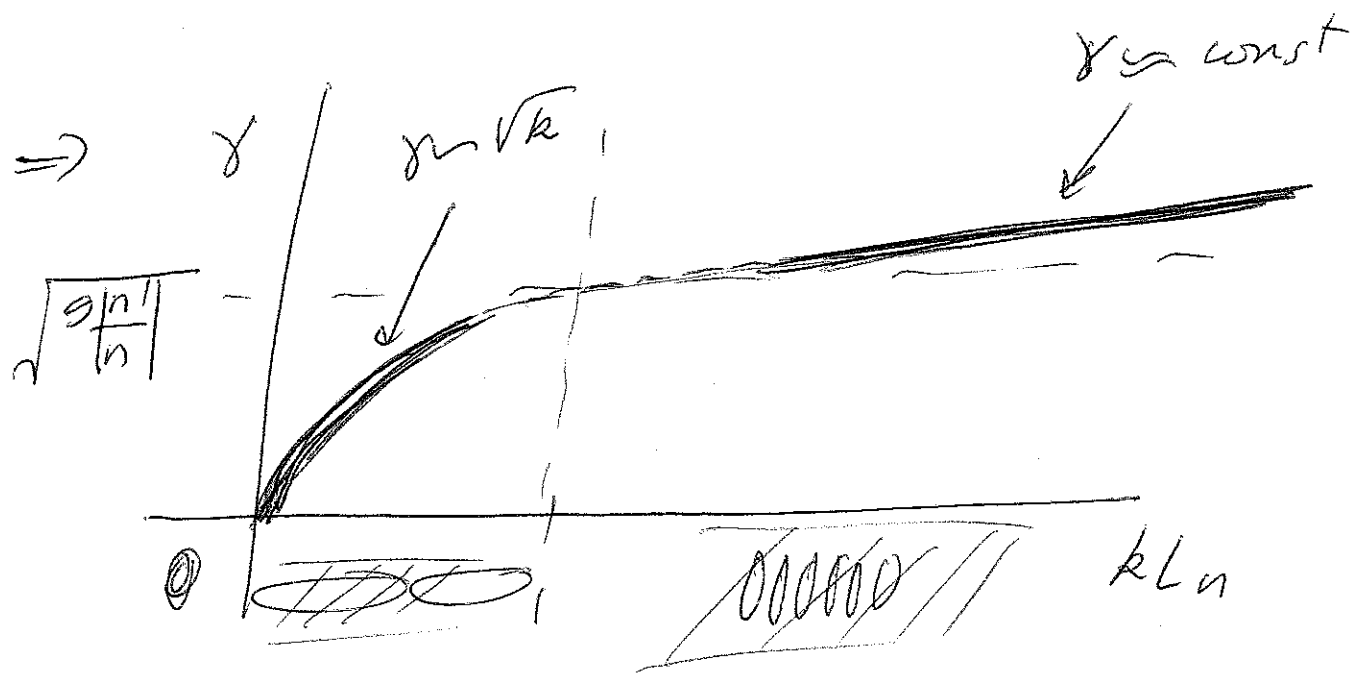
18

also can show $[\tilde{\phi}]_+^+ = 0$, by a 2nd \int

$$\Rightarrow -n_+k - n_-k = \frac{g\hbar^2}{\omega^2} [n_+ - n_-]$$

$$\Rightarrow \boxed{\omega^2 = -gk \frac{(n_+ - n_-)}{(n_+ + n_-)}}$$

Need $k \Delta_n \ll 1$ and $\Delta_n \gg L_n$, $\Delta k \ll 1$
 $\Rightarrow kL_n \ll 1$



γ peaks weakly as $k \rightarrow \infty$