



Faradays Law of Induction

Up to this point we have considered magnetic

fields produced and electric fields produced

by steady state currents and charge distributions

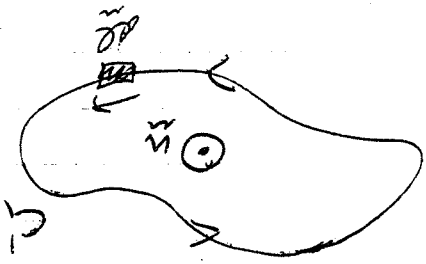
~~fields~~ we now want to begin to generalize

our results to allow for the situation.

Consider a circuit  $C$  as shown. The total

magnetic flux linking this

changing  $\vec{B} \rightarrow \vec{E}$



$$\Psi = \int_S \vec{B} \cdot \vec{n} \, dS$$

Observations are that an electric field is produced in the circuit if the magnetic flux  $\Phi$  changes in time.  $\nabla \times \vec{E} = -\dot{\vec{B}}$

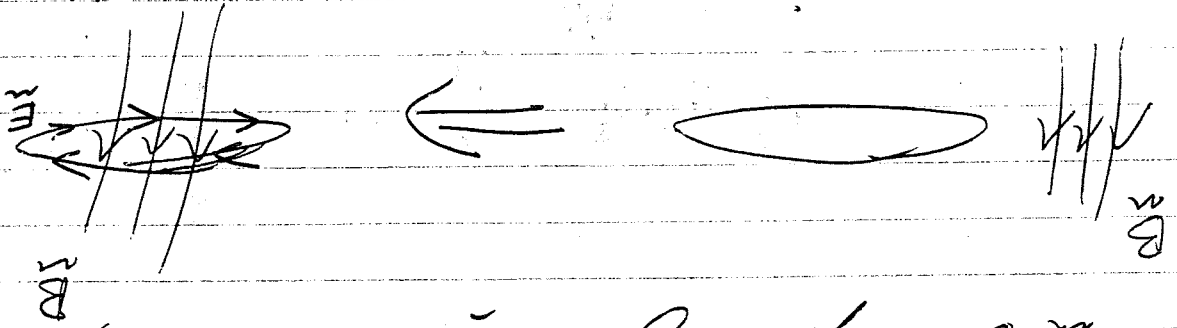
$$\vec{E} = \oint \vec{E} \cdot d\vec{l} = \text{electromotive force}$$

$$\boxed{\vec{E} = -\dot{\vec{A}} \vec{F}} \quad \text{Faradays}$$

one direction of the electric field is such that

the magnetic current would oppose the change

in the magnetic field (Lenz' law)



In generalise this result to a arbitrary closed loop in space.

$$\mathcal{E} = \oint_C \vec{E} \cdot d\vec{r} = \int_S \nabla \times \vec{E} \cdot d\vec{S}$$

$$\int_{\partial S} [\nabla \times \vec{E} + \frac{\partial \vec{B}}{\partial t}] \cdot d\vec{S} = 0$$

Since this is valid for any surface  $S$ , the result must be valid locally at any point

$$\nabla \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$$

Maxwell's 3rd eq

Faraday's law

and that for  $\partial B / \partial t = 0$  this reduces to

$$\nabla \times \vec{E} = -\vec{\nabla} \phi$$

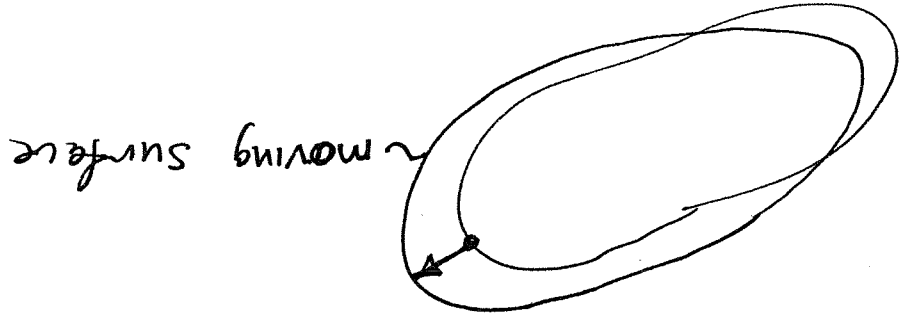
As a result, assigned potential

consider a magnetic field

$\tilde{B}(\vec{x}, t)$  and

associated  $\tilde{E}(\vec{x}, t)$

$$\frac{d}{dt} \int_{L(t)} \tilde{B} \cdot d\vec{l} = \frac{d}{dt} \int_{S(t)} (\tilde{B} \cdot \tilde{n}) da$$

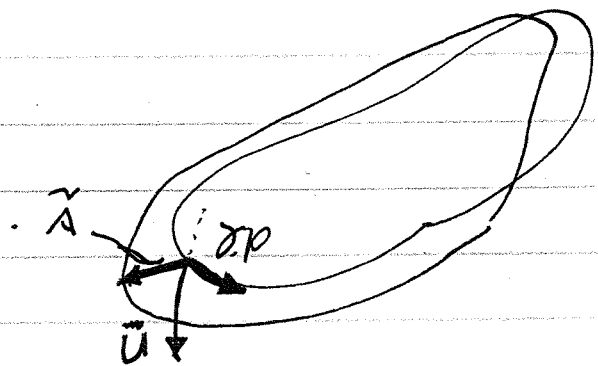


$$\frac{d}{dt} \int_{S(t)} (\tilde{B} \cdot \tilde{n}) da = \int_{S(t)} \frac{\partial}{\partial t} \tilde{B} \cdot \tilde{n} da +$$

$$\frac{1}{S(t)} \int_{S(t)} \tilde{B} \cdot \tilde{n} da$$

Let  $\vec{v}(\vec{x}, t)$  be the velocity of

the element of the surface



$$\vec{n} \, da = -d\vec{r} \times \vec{v} \, dt =$$

$$-c \, \nabla \times \vec{E}$$

$$\frac{d}{dt} \int_{S(t)} (\vec{B} \cdot \vec{n}) \, da = \int_{S(t)} da \, \vec{n} \cdot \frac{\partial \vec{B}}{\partial t}$$

$$- \int_C \vec{B} \cdot d\vec{\ell} \times \vec{v}$$

$$- \int_C d\vec{\ell} \cdot \vec{v} \times \vec{B}$$

"electric field felt  
by charge moving with  
velocity  $\vec{v}$ "

$$\vec{E}' = \vec{E} + \vec{v} \times \vec{B}$$

electromotive force

$$\text{EMF} = \int \vec{v} \cdot \vec{E}' = - \frac{d\Phi}{dt}$$

$$\vec{E}' = \vec{E} + \vec{v} \times \vec{B}$$

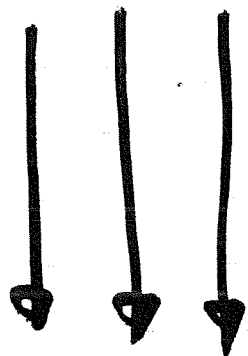
$$= - \int \vec{v} \cdot \vec{E}' =$$

$$\frac{d}{dt} = - \int \vec{v} \cdot \left[ \vec{E} + \vec{v} \times \vec{B} \right]$$

For moving loop

# Homopolar generator

get examples



$$\dot{\mathbf{B}} = \mathbf{B}_0 \mathbf{e}_z$$

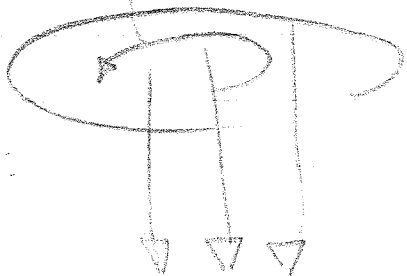
constant

$$\dot{\mathbf{r}} = 0$$

Now add spinning metal disc

spinning disc there is

Force  $\propto \dot{\mathbf{r}} \times \mathbf{B} = \dot{\mathbf{E}}$



$$\dot{\mathbf{r}} = v \mathbf{e}_\theta = \Omega r \mathbf{e}_\theta$$

For a good conductor  $\dot{\mathbf{E}} = 0$

$$\mathbf{E}_r + v \mathbf{e}_\theta \cdot \mathbf{B}_0 = 0 \quad \mathbf{E}_r = -\Omega r B_0$$

electrostatic field

$$\phi = \frac{\Omega r^2 B_0}{2}$$

$$= \int_0^a \frac{dr}{c} \cdot \Omega a^2 B_0 = \frac{\Omega a^2 B_0}{c}$$

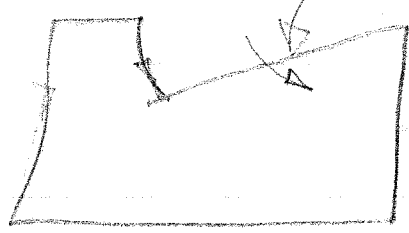
$$\int d\vec{l} \cdot (\vec{E} + \vec{v} \times \vec{B})$$

if the element is embedded in the moving disc

$$\vec{E} = -\nabla\phi$$

loop = 0

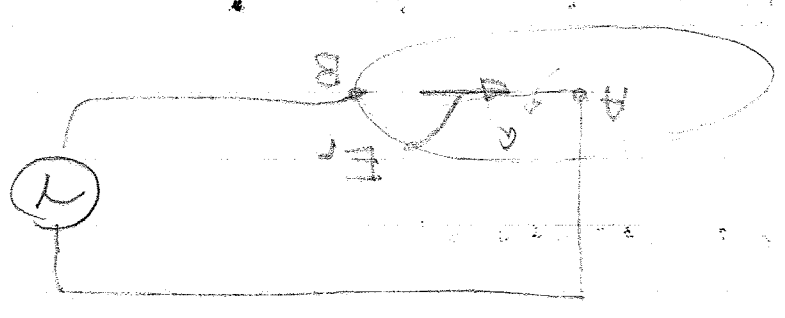
$$\int d\vec{l} \cdot (\vec{E} + \vec{v} \times \vec{B})$$



$$\mathcal{E}_{mf} = \oint d\vec{l} \cdot (\vec{E} + \vec{v} \times \vec{B})$$

$$\phi_B - \phi_A = \frac{\Omega a^2 B_0}{c}$$

Voltmeter



# Dynamics

resistivity =  $\frac{1}{\sigma}$

$$\bar{\mathbf{E}} + \tilde{\mathbf{v}} \times \tilde{\mathbf{B}} = n \tilde{\mathbf{j}}$$

$$\bar{\mathbf{E}} = n \tilde{\mathbf{j}} - (\tilde{\mathbf{v}} \times \tilde{\mathbf{B}})$$

$$+ \frac{\partial \tilde{\mathbf{B}}}{\partial t} = -\nabla \times \bar{\mathbf{E}}$$

$$\tilde{\mathbf{v}} = \frac{1}{\mu_0} \nabla \times \tilde{\mathbf{B}}$$

$$= -\nabla \times n \frac{1}{\mu_0} (\nabla \times \tilde{\mathbf{B}}) + \nabla \times (n \tilde{\mathbf{v}} \times \tilde{\mathbf{B}})$$

$$\frac{\partial \tilde{\mathbf{B}}}{\partial t} = -\frac{1}{\mu_0} \nabla (\nabla \cdot \tilde{\mathbf{B}}) - \nabla^2 \tilde{\mathbf{B}} + \tilde{\mathbf{B}} \cdot \nabla \tilde{\mathbf{v}} + \nabla \cdot \tilde{\mathbf{B}} \tilde{\mathbf{v}}$$

$$-\tilde{\mathbf{v}} \cdot \nabla \tilde{\mathbf{B}} - \tilde{\mathbf{B}} \cdot \nabla \tilde{\mathbf{v}}$$

NOW LET'S ~~ASSUME~~  $(\nabla \cdot \tilde{\mathbf{B}}) = 0$

ASSUME

$$\nabla \cdot \tilde{\mathbf{v}} = 0$$

incompressible flow

$\tilde{\mathbf{v}}$ -given

$$\frac{\partial \tilde{\mathbf{B}}}{\partial t} + \tilde{\mathbf{v}} \cdot \nabla \tilde{\mathbf{B}} = \tilde{\mathbf{B}} \cdot \nabla \tilde{\mathbf{v}} + \left(\frac{1}{\mu_0}\right) \nabla^2 \tilde{\mathbf{B}}$$

time derivative following flow

flow

APPLICATION

to inhomogeneous flow

DIFFUSION OF

$\tilde{\mathbf{B}}$

due to resistivity

(can be very small in a

good conductor)

what is Behaviour of  $\tilde{B}(\vec{x}, t)$  ?

Suppose  $\tilde{V} = \text{const}$

$$\frac{\partial \tilde{B}}{\partial t} + \tilde{V} \cdot \nabla \tilde{B} = \left(\frac{\mu_0}{m}\right) \nabla^2 \tilde{B}$$

all coefficients independent of  $\vec{x}, t$

Look for  $\tilde{B} = \text{Re} \left\{ \tilde{B} e^{i\vec{k} \cdot \vec{x} - \omega t} \right\}$

$$(\omega + i\vec{k} \cdot \tilde{V}) \tilde{B} = -\frac{\mu_0}{m} k^2 \tilde{B}$$

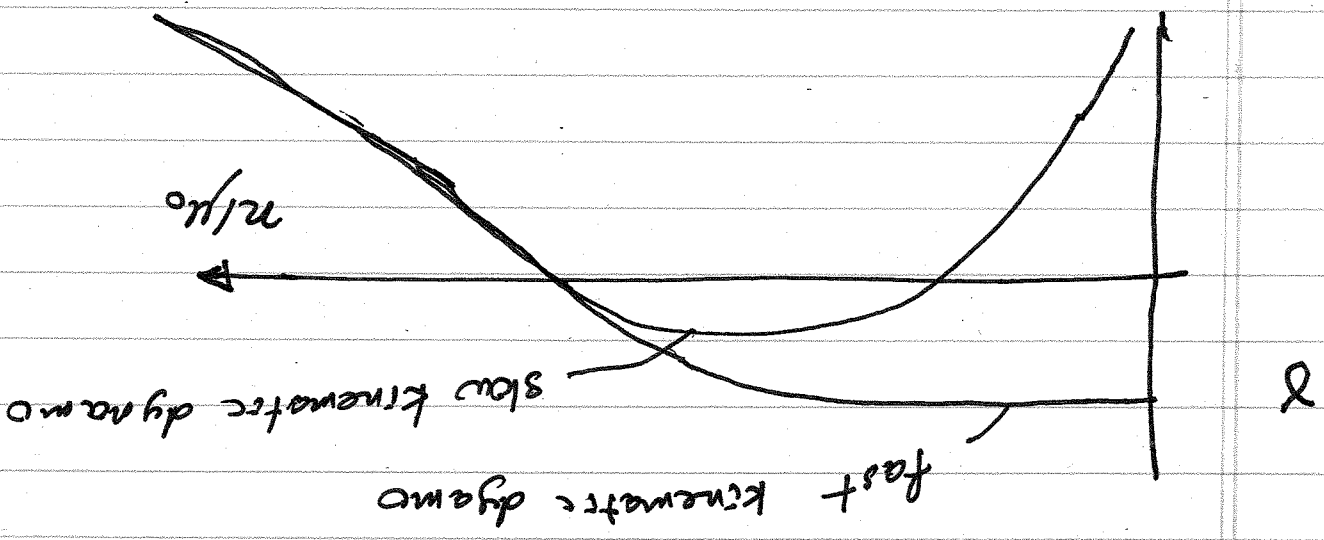
$$\omega = i\vec{k} \cdot \tilde{V} - \frac{\mu_0}{m} k^2$$

$\text{Re} \{ \omega \} < 0$

field decays in time!

3

What if  $\tilde{V}$  is inhomogeneous  
Then the following is possible



If  $r > 0$  small seed  $\tilde{B}$  is amplified  
in time.

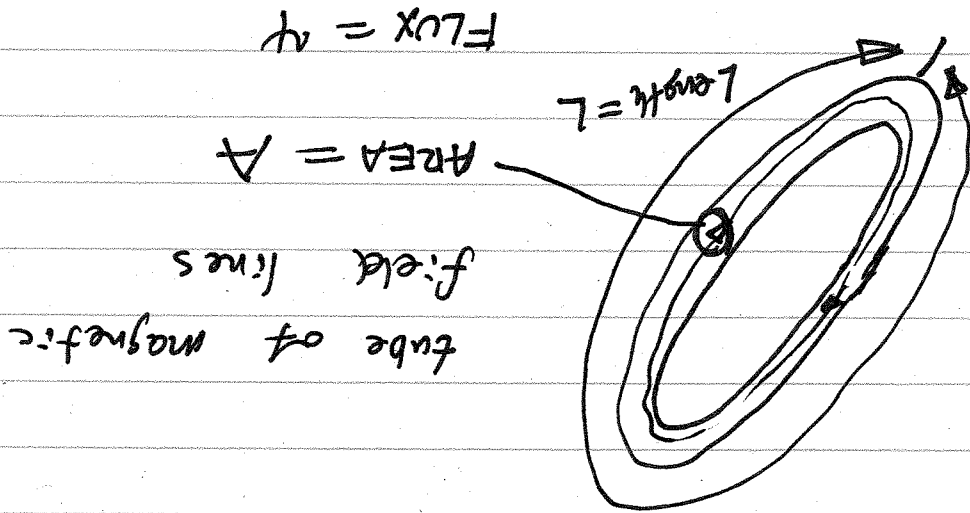
ORIGIN OF: TERRESTRIAL, SOLAR,  
GALACTIC MAGNETIC  
FIELDS

4

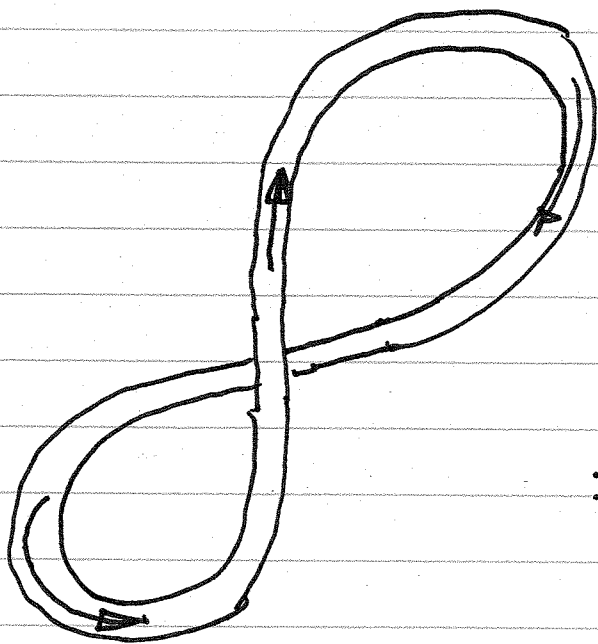
WHAT KIND OF  $\tilde{V}(\vec{x}, t)$  IS NEEDED?

ANS: stretch - twist - fold

CONSIDER CARTON:



The diagram shows a small surface element  $dA$  on the surface of a magnetic field tube. A normal vector  $\vec{n}$  is shown pointing outwards from the surface. The magnetic field vector  $\vec{B}$  is shown as a line passing through the surface element. The text  $\Psi = \int \vec{B} \cdot \vec{n} dA$  is written next to the diagram.

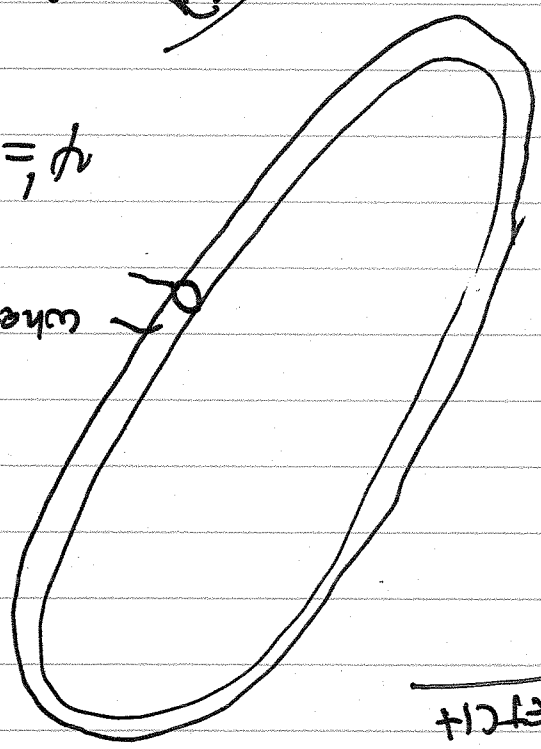


TWIST

$$\oint \text{curl}(\vec{E} + \vec{v} \times \vec{B}) = -\frac{d\Phi}{dt} = 0$$

$$\psi = \psi'$$

what is  $\psi'$ ?



STRETCH

Imagine flow stretches tube to  $2L$

~~$\psi = \text{const}$~~

$$\vec{V} = A'\vec{L}' = AL$$

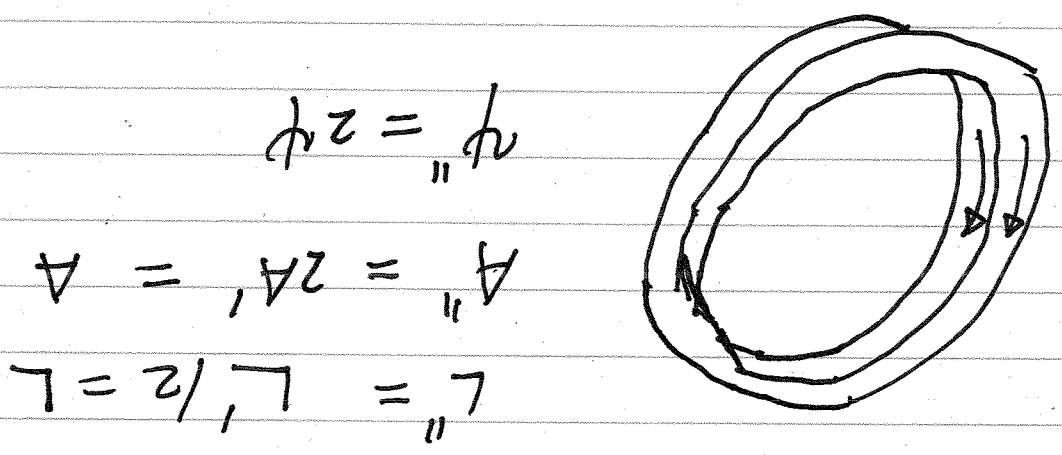
$$\nabla \cdot \vec{v}$$

$$L' = 2L$$

$$A' = A/2$$

FOLD

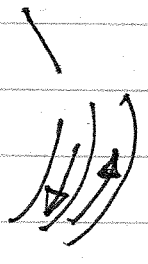
New loop



FIELD HAS ~~BEEN~~ DOUBLED

Picture is actually more complicated

Sources stretching and twisting results in this situation



resistor (flux diffusion cases cancellation)