

Potential Energy of a permanent dipole \vec{m} in an external field

$$U = -\vec{m} \cdot \vec{B}_{\text{ext}}$$

in an external field $\nabla \times \vec{B}_{\text{ext}} = 0$

Can not derive our result yet

Since ρ true dependence of \vec{J}

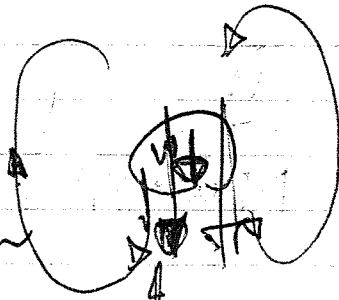
not considered.

Lowest Energy State

$$\vec{m} \parallel \vec{B}_{\text{ext}}$$

tend to align

$\vec{B} \sim \vec{B}$ due to \vec{m}



To Here end of 10th lecture

12 pages

Review

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} \quad \nabla \cdot \vec{B} = 0$$

Force on a current distribution

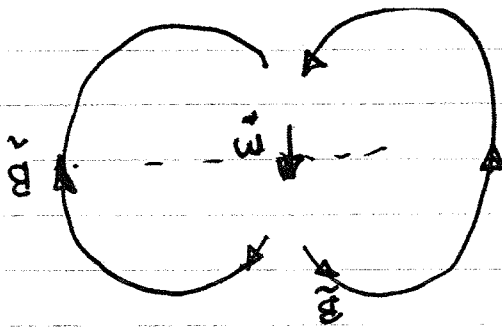
$$\vec{F} = \int d^3x \vec{j}(\vec{x}) \times \vec{B}(\vec{x})$$

Vector potential

$$\vec{B} = \nabla \times \vec{A}$$

Coulomb gauge $\nabla \cdot \vec{A} = 0$

$$\vec{A} = \frac{1}{c} \int d^3x' \frac{\vec{j}(\vec{x}')}{|\vec{x} - \vec{x}'|}$$



magnetic dipole field

$$\vec{B} = \nabla \times \vec{A} = \frac{3\vec{x} \vec{m} \cdot \vec{x} - \vec{m} |\vec{x}|^2}{|\vec{x}|^3} - \nabla \left(\frac{\vec{m} \cdot \vec{x}}{|\vec{x}|^3} \right)$$

outside region of current

magnetization

magnetic moment density

$$\vec{M}(\vec{x}')$$

$$\vec{m} = \frac{1}{2} \int d^3x' \vec{x}' \times \vec{J}(\vec{x}') / 2$$

magnetic moment

$$\vec{A} = -\frac{\mu_0 / 4\pi}{|\vec{x}|^3} \vec{x} \times \vec{m} = \frac{\vec{m} \times \vec{x}}{|\vec{x}|^3}$$

Vector Potential due to a localized current distribution $|\vec{x}'| \gg |\vec{x}|$

Relation between magnetic moment and angular momentum

$$\vec{m} = \frac{1}{2} \int \vec{x}' \times \vec{J}(\vec{x}') d^3x'$$

suppose $\vec{J}(\vec{x}')$ is composed

of an ensemble of charges

and mass m with charges q • each

charge has an instantaneous

position \vec{x}_i and velocity \vec{v}_i

delta function

$$\vec{J}(\vec{x}') = \sum_N q_i \vec{v}_i \delta(\vec{x}' - \vec{x}_i)$$

L density why:

Remember for continuous

$$\vec{J} = \rho \vec{v} \quad \rho \approx \sum_i \delta(\vec{x}' - \vec{x}_i)$$

$$\vec{m} = \frac{1}{2} \sum_i m_i \vec{x}_i \times \vec{v}_i \quad \vec{v}_i = \frac{d\vec{x}_i}{dt}$$

$\tilde{L}_i =$ angular momentum of i th charge.

$$\tilde{m} = \frac{e}{am} \sum_i \tilde{L}_i = \frac{e}{am} \tilde{L}$$

total angular momentum vector

Classical Result

Forces and Torques on localized current distribution

$$\vec{F} = \int d^3x' \underbrace{\vec{j}(\vec{x}') \times \vec{B}(\vec{x}')}_{\text{Force density}}$$

suppose $\vec{B}(\vec{x}')$ is localized

to a region of size R

and $\vec{B}(\vec{x}')$ varies smoothly

over this region.

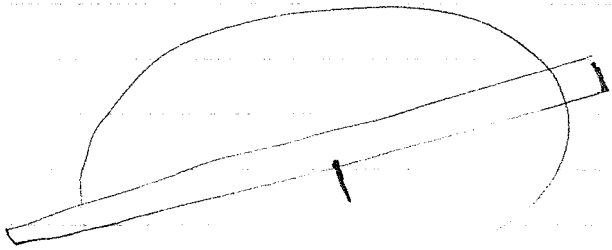
$l=1,2,3$
 $x_{1,2}$

~~Second term~~

$$\frac{c}{i} \int d^3x' \tilde{j}(\tilde{x}') \times (\tilde{x}' \cdot \nabla \tilde{B}(\tilde{x}))$$

$$= \sum_{l=1}^3 \int d^3x' \tilde{j}(\tilde{x}') \times \tilde{B}(\tilde{x}') \times \tilde{x}'$$

$\tilde{x}'=0$



$$(\nabla \cdot \tilde{j}) = 0$$

FIRST TERM VANISHES $\int d^3x' \tilde{j}(\tilde{x}') = 0$

$$+ \frac{c}{i} \int d^3x' \tilde{j}(\tilde{x}') \times (\tilde{x}' \cdot \nabla \tilde{B}(\tilde{x})) + \dots$$

$$\tilde{F} \approx \frac{c}{i} \int d^3x' \tilde{j}(\tilde{x}') \times \tilde{B}(0)$$

Taylor Expand \tilde{B}

$$\tilde{B}(\tilde{x}) = \tilde{B}(0) + \tilde{x} \cdot (\nabla \tilde{B}(\tilde{x})) + \dots$$

$\tilde{x}'=0$

$$\vec{F} = \nabla(\vec{m} \cdot \vec{B})$$

assumes \vec{m} is constant $\vec{m} \propto \vec{B}$

$$\vec{F} = (\vec{m} \times \nabla) \times \vec{B} = \nabla \vec{m} \cdot \vec{B} - \vec{m} \nabla \cdot \vec{B}$$

$\nabla \cdot \vec{B} = 0$

$$F_i = + \sum_{jk} \epsilon_{ijk} [(\vec{m} \times \nabla)_j B_k]$$

$$\frac{1}{c} \int d^3x' \vec{x} \cdot \vec{x}' \cdot \vec{J}'_j = - [\vec{x} \times \vec{m}]_j$$

in comp replace \vec{x} by ∇B_k

$$= - \vec{x} \times \vec{m}$$

$$\int d^3x' \vec{x} \cdot \vec{x}' \cdot \vec{J}' = - \frac{1}{c} \int d^3x' [\vec{x} \times (\vec{x}' \times \vec{J}')]]$$

We have previously derived

$$\epsilon_{ijk} = \begin{cases} +1 & \text{if } i,j,k \text{ are in order} \\ 0 & \text{otherwise} \end{cases}$$

$$\vec{F}_i = \sum_{jk} \epsilon_{ijk} \int d^3x' J'_j(x') \frac{1}{r} x'_k \left(\frac{\partial}{\partial x'_k} B_k \right)$$

Second Term $\vec{x}' \cdot \nabla B_k$

same as \vec{E} field due to dipole \vec{p}

$$\vec{B} = -\frac{\mu_0}{4\pi} \frac{\vec{m} \times \nabla}{r^3} = \frac{\mu_0}{4\pi} \frac{3\vec{m} \times \vec{r} - \vec{m} r^2}{r^5}$$

$$\tilde{N} = \int d^3x' \tilde{\mathbf{x}}' \cdot \tilde{\mathbf{B}} = \tilde{M} \times \tilde{\mathbf{B}}$$

$$= \int d^3x' \tilde{\mathbf{x}}' \cdot \tilde{\nabla}' \tilde{\mathbf{B}} = 0$$

$$\int d^3x' \tilde{\mathbf{x}}' \cdot \tilde{\nabla}' \tilde{\mathbf{B}} = 0$$

($\tilde{\nabla}' \cdot \tilde{\mathbf{B}} = 0$)

$$= \int d^3x' \tilde{\mathbf{x}}' \cdot [\tilde{\mathbf{B}} - \tilde{\mathbf{B}}(\infty)]$$

$$\tilde{N} = \int d^3x' \tilde{\mathbf{x}}' \cdot (\tilde{\mathbf{B}}(\infty) - \tilde{\mathbf{B}}(0))$$

Lowest order term

$$\tilde{\mathbf{B}} = \tilde{\mathbf{B}}(0) + \tilde{\nabla} \tilde{\mathbf{x}} \cdot \tilde{\mathbf{B}}$$

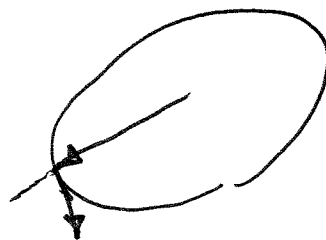
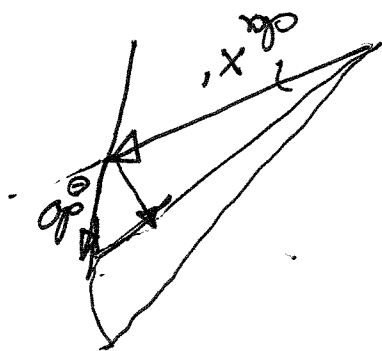
$$= \int d^3x' \tilde{\mathbf{x}}' \cdot \tilde{\nabla}' (\tilde{\mathbf{x}} \times \tilde{\mathbf{B}})$$

$$\tilde{N} = \int d^3x' \tilde{\mathbf{x}}' \cdot \tilde{\mathbf{F}}$$

The total torque

$$|\tilde{m}| = \int \phi da$$

$$da = x' dy_2 = \frac{z}{2}$$



derive \tilde{m} for a power loop

$$\tilde{N} = \int d^3\tilde{x}' (\tilde{x}' \times (\tilde{\nabla} \times \tilde{B}^{ext})) \times \tilde{B}^{ext}(0)$$

TORQUE

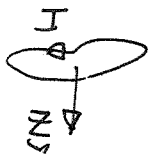
$$\tilde{F} = \tilde{\nabla} \tilde{m} \cdot \tilde{B}^{ext}(\tilde{x}) \Big|_{\tilde{x}=0}$$

$$\tilde{B}^{ext}(\tilde{x}) \approx \tilde{B}^{ext}(0) + \tilde{x} \cdot \tilde{\nabla} \tilde{B}^{ext} \Big|_{\tilde{x}=0}$$

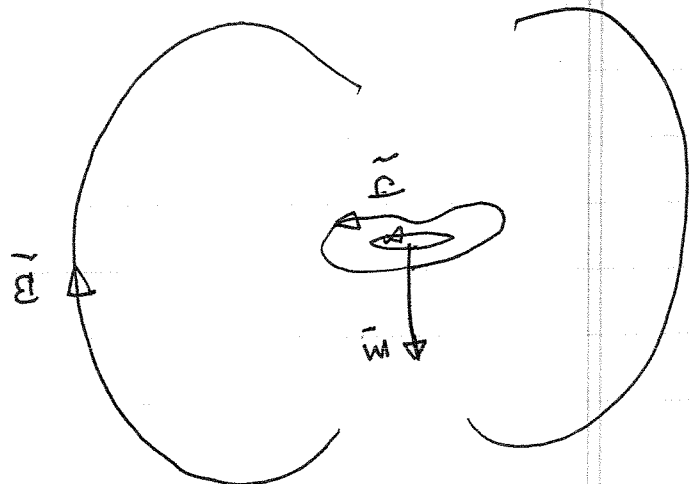
Force on a dipole in an external Field

$$\tilde{m} = \frac{1}{2} \int \tilde{x}' \times d\tilde{x}' I = \tilde{z} I A$$

Evaluate \tilde{m} for a current loop



Planar



IF $\tilde{x} \parallel \tilde{m}$ ~~$\tilde{B} \parallel \tilde{m}$~~ $\tilde{B} \cdot \tilde{m} > 0$
 if $\tilde{x} \perp \tilde{m}$ $\tilde{B} \cdot \tilde{m} < 0$

Review of Fields due to localized

currents

In general: $\vec{B} = \nabla \times \vec{A}$

$$\vec{A} = \frac{\mu_0}{4\pi} \int d^3x' \frac{\vec{J}(\vec{x}')}{|\vec{x} - \vec{x}'|}$$

For localized current density expand

$$\frac{1}{|\vec{x} - \vec{x}'|} \approx \frac{1}{|\vec{x}|} \left[1 - \frac{\vec{x}' \cdot \vec{x}}{|\vec{x}|^2} + \dots \right]$$

$$\vec{A} = -\frac{\mu_0}{4\pi} \frac{\vec{x} \times \vec{m}}{|\vec{x}|^3}$$

where $\vec{m} = \frac{1}{2} \int d^3x' \vec{x}' \times \vec{J}(\vec{x}')$

dipole moment

$$\vec{B} = \nabla \times \vec{A} = \frac{\mu_0}{4\pi} \left[\frac{3\vec{x} \times \vec{m} - x^2 \vec{m}}{|\vec{x}|^5} \right]$$

same as E field due to dipole p

Magnetic materials

$$\nabla \times \vec{B} = \mu_0 \vec{j} + \mu_0 \vec{j}_m$$

follow

procedure of dielectrics

free current density

induced current density

$$\vec{j} = \vec{j}_f + \vec{j}_i$$

note: $\nabla \cdot \vec{j}_f = 0$

$$\nabla \cdot \vec{j}_i = 0$$

$$\rho_i = -\nabla \cdot \vec{P}$$

$$\nabla \cdot \vec{E} = 4\pi(\rho_f + \rho_i)$$



answer

$$\vec{j}_i = \nabla \times \vec{M}$$

$\nabla \cdot \vec{j}_i = 0$ implies $\vec{j}_i = \nabla \times (\text{something})$

$$\vec{M} = \sum_i \langle \vec{m}_i \rangle n_i$$

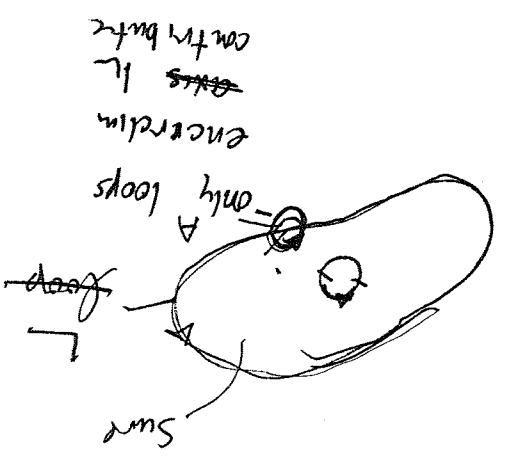
number density
average magnetic moment

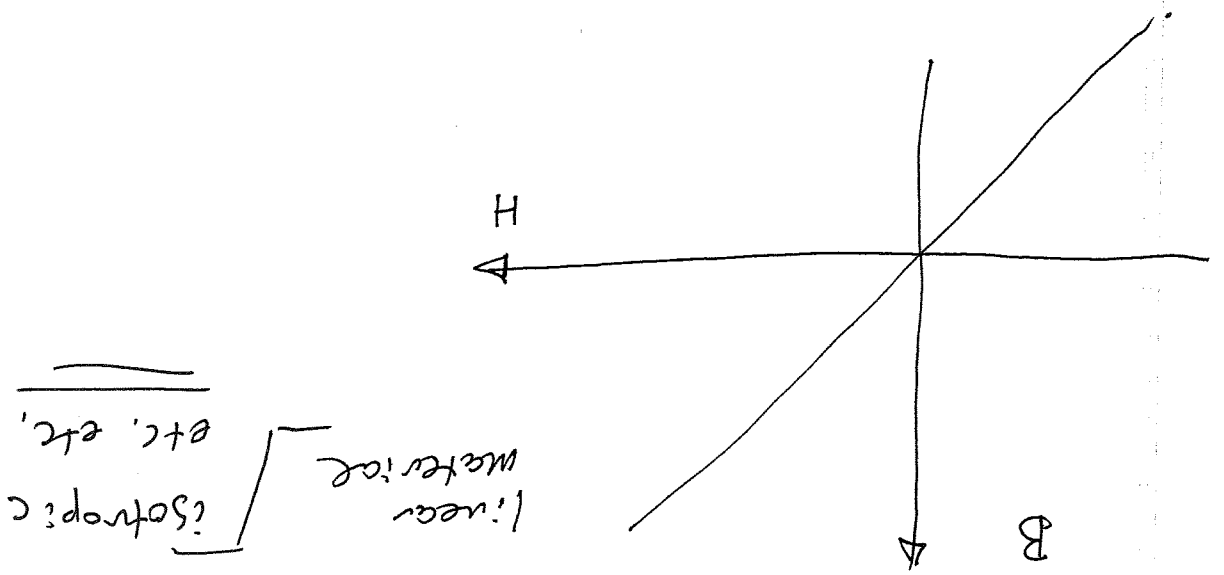
$\vec{M} =$ magnetization or

magnetic moment density

~~Alternate~~ approach

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{enclosed}$$





$$\bar{H} = \mu \bar{H} \quad \bar{H} \sim \mathbb{R} \cdot \mathbb{J}^{free}$$

$$\tilde{B} = \mu_0 (\tilde{H} + \tilde{M}) \quad \text{if } \tilde{M} \propto \tilde{H}, \tilde{B}$$

call this \bar{H}

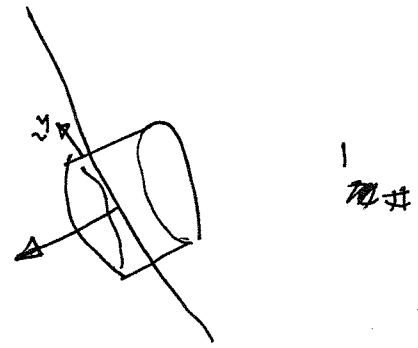
$$\nabla \times \begin{bmatrix} \tilde{B} \\ \mu_0 \tilde{M} \end{bmatrix} = \tilde{J}^{free}$$

$$= \mu_0 [\tilde{J}^{free} + \nabla \times \tilde{M}]$$

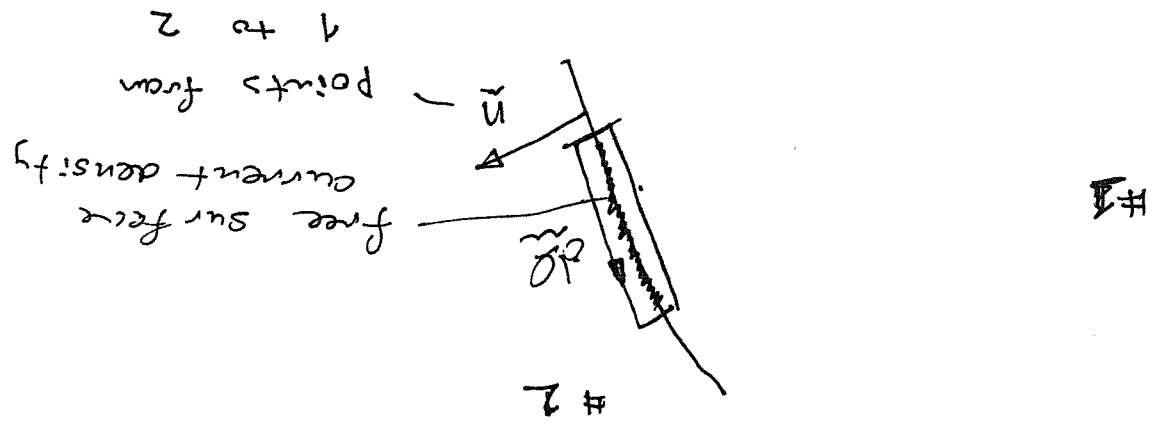
$$\nabla \times \tilde{B} = \mu_0 \tilde{J} = \mu_0 [\tilde{J}^{free} + \tilde{J}^{ind}]$$

Boundary Conditions

$$\nabla \cdot \tilde{\mathbf{B}} = 0 \quad \nabla \times \tilde{\mathbf{H}} = \tilde{\mathbf{j}}_{free}$$



$$\int_{\#2}^{\#1} \tilde{\mathbf{B}} \cdot \tilde{\mathbf{n}} dA = \Delta \quad \tilde{\mathbf{n}} \cdot (\tilde{\mathbf{B}}_2 - \tilde{\mathbf{B}}_1) = 0$$



$$\oint \tilde{\mathbf{H}} \cdot d\tilde{\mathbf{l}} = I_{enclsd}$$

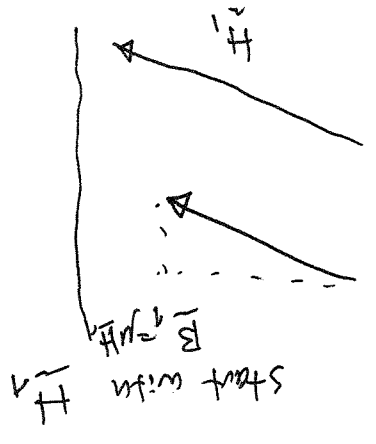
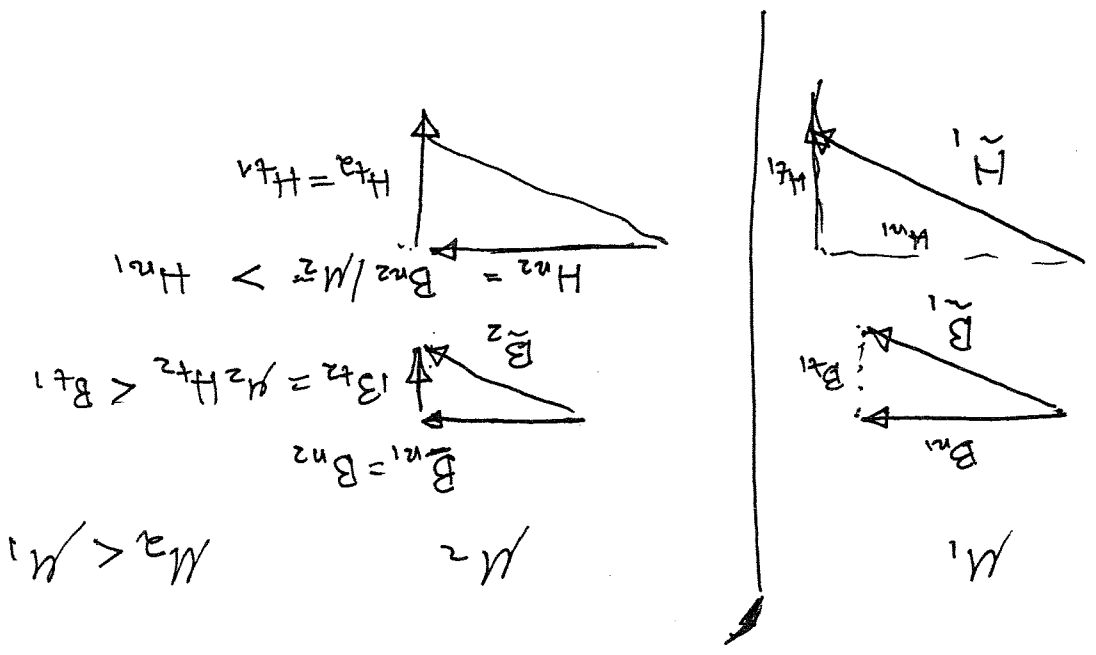
$$(\tilde{\mathbf{H}}_2 - \tilde{\mathbf{H}}_1) \cdot d\tilde{\mathbf{l}}_2 = \tilde{\mathbf{K}}_s \cdot (\tilde{\mathbf{n}} \times d\tilde{\mathbf{l}}_2)$$

$$= (\tilde{\mathbf{K}}_s \times \tilde{\mathbf{n}}) \cdot d\tilde{\mathbf{l}}_2$$

$$\tilde{\mathbf{H}}_2 - \tilde{\mathbf{H}}_1 = \tilde{\mathbf{K}}_s \times \tilde{\mathbf{n}}$$

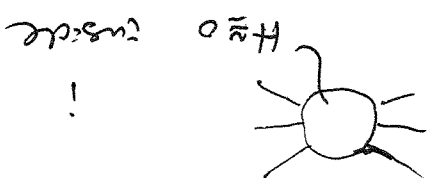
$$\tilde{\mathbf{K}}_s = \tilde{\mathbf{n}} \times (\tilde{\mathbf{H}}_2 - \tilde{\mathbf{H}}_1)$$

free



What happens when $\mu_2 > \mu_1$

! if B_{t2} is finite $H_{t2} \rightarrow 0$



$D = \epsilon \bar{E}$
 $\bar{B} = \mu \bar{H}$

solve with